



Exploring Higgs Boson Properties

From Precision Measurements to Interpretation

Chiara Arcangeletti

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The Standard Model and the Higgs boson

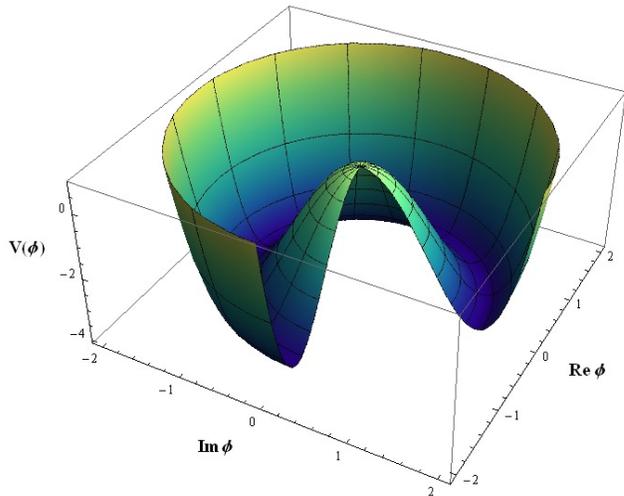
- Electromagnetic interaction
- Strong interaction
- Weak interaction

Gauge bosons, fermions, interactions, but...

...no masses...

Spontaneous symmetry breaking:

Higgs Mechanism

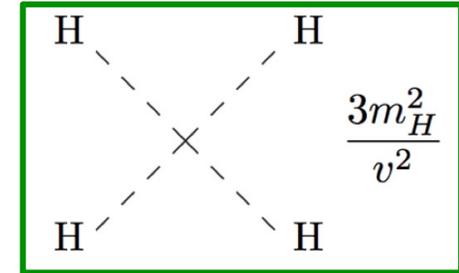


$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\Psi} \not{D} \Psi$$

$$+ \bar{\Psi}_i y_{ij} \Psi_j \phi + \text{h.c.}$$

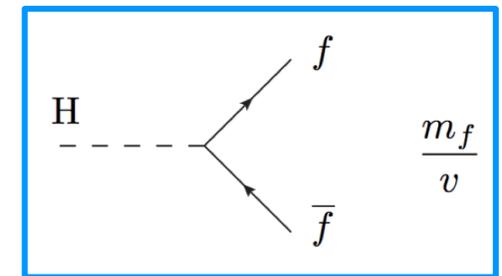
$$+ \frac{1}{2} D_\mu \phi^\dagger D^\mu \phi - V(\phi)$$

Self - interactions

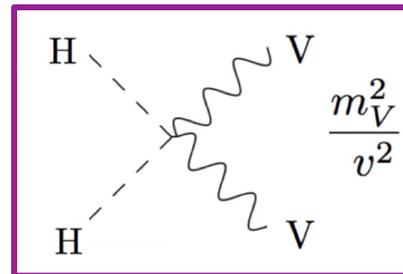


Yukawa interactions

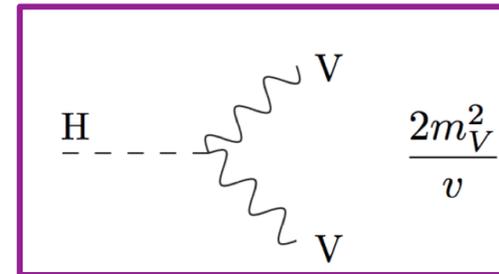
Fermion masses



Gauge interactions



Gauge boson masses

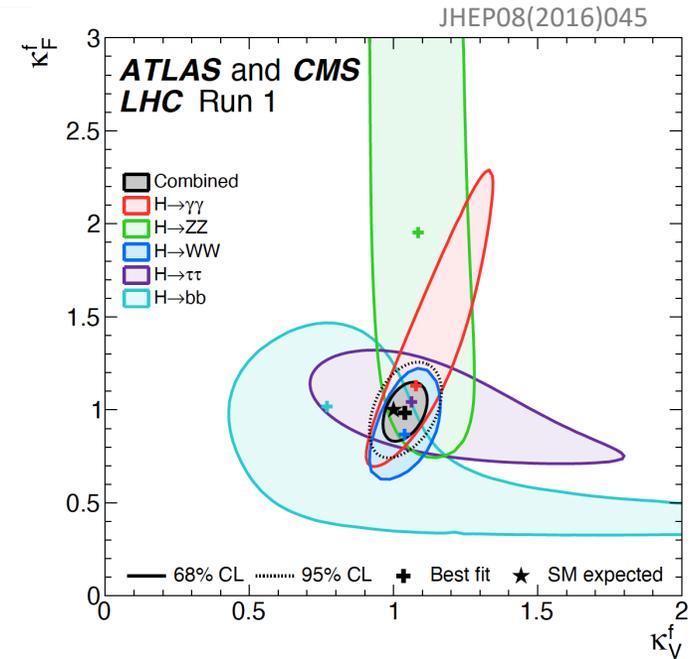
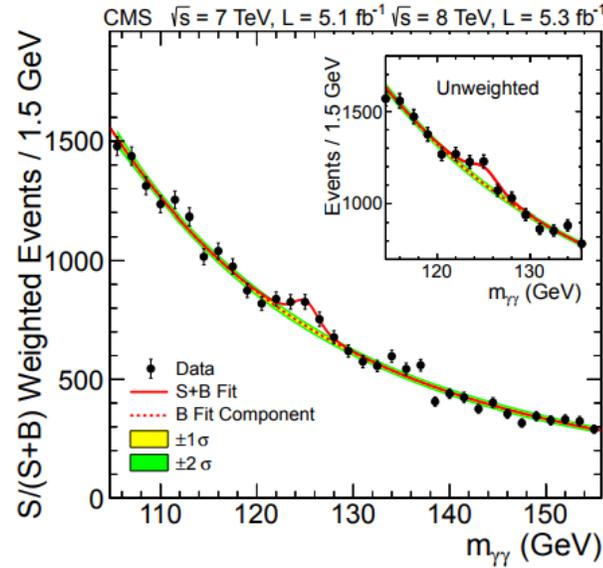
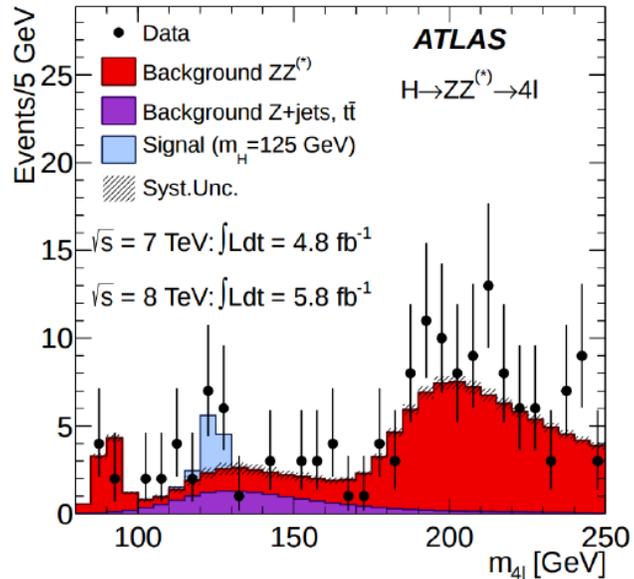


Couplings of the Higgs boson with the other particles as **fingerprints** of the Standard Model

The evolution of the Higgs boson

Run 1

Discovery of the Higgs boson... and first measurements of its properties!



Beginning of *Precision Era*...

More precise measurements of the Higgs mass, width, couplings and cross sections

More stringent constraints on anomalous Higgs boson couplings with other SM particles
Interpretation of the results in different theoretical framework

Run 2

Run 3

...walking towards higher energy and higher luminosity

looking to the Higgs boson production at 13.6 TeV!

The Large Hadron Collider

Proton – Proton collider of 27 km with 4 interaction points

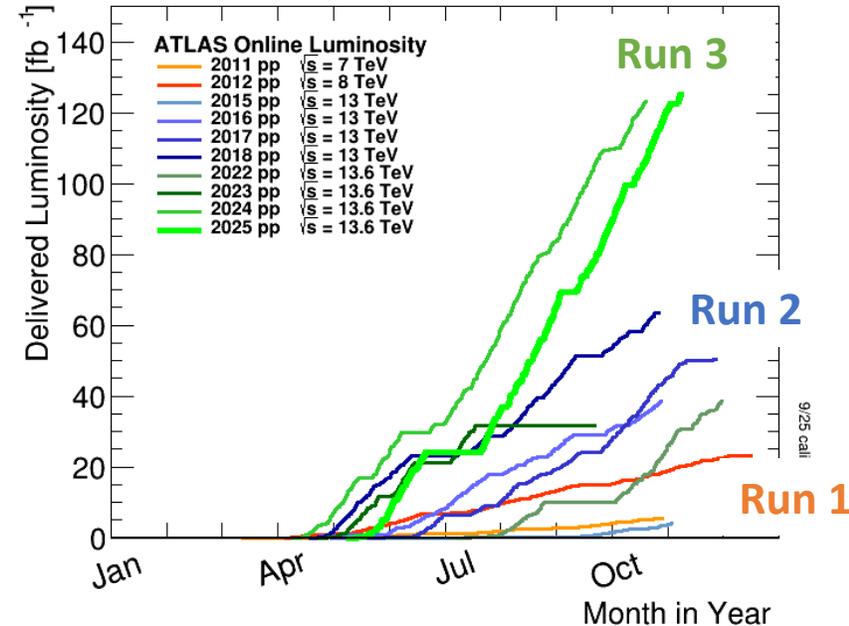
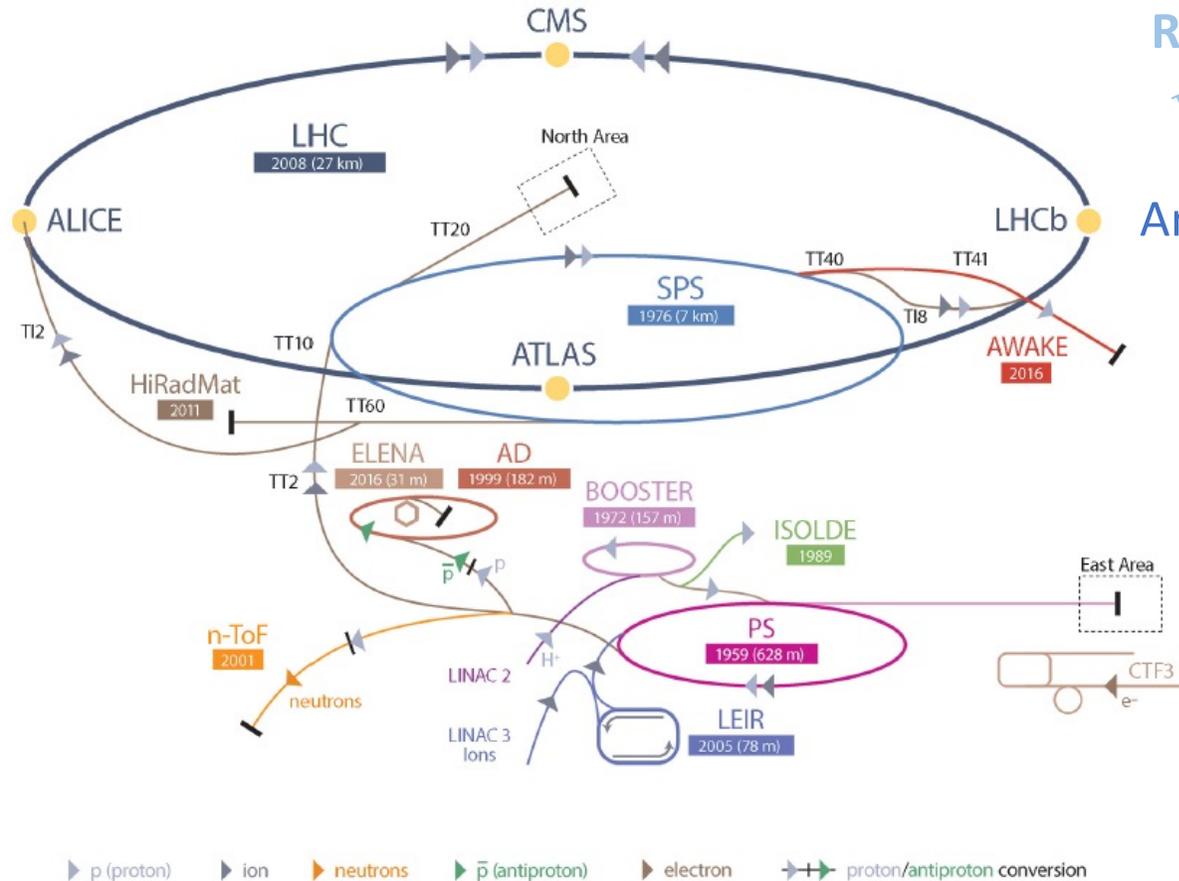
Why protons? Protons consist of quarks and gluons
 → their interactions produce other particles

Higher Energy collisions with respect to electron-positron colliders
 → hadron colliders are for “Discovery”

| | | |
|--------------------------------|-----------------------------|-------------------------------|
| Run 1 (2011-2012) | Run 2 (2015-2018) | Run 3 (2022-now) |
| $\sqrt{s} = 7 - 8 \text{ TeV}$ | $\sqrt{s} = 13 \text{ TeV}$ | $\sqrt{s} = 13.6 \text{ TeV}$ |

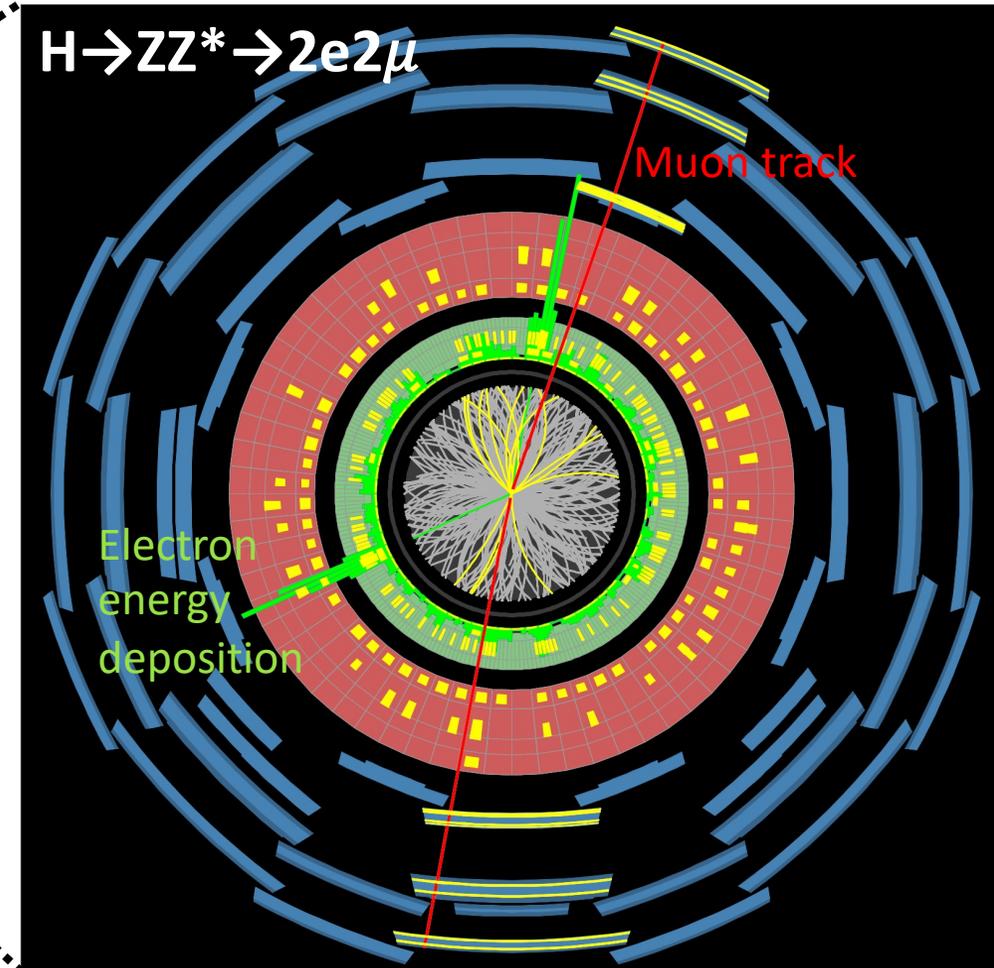
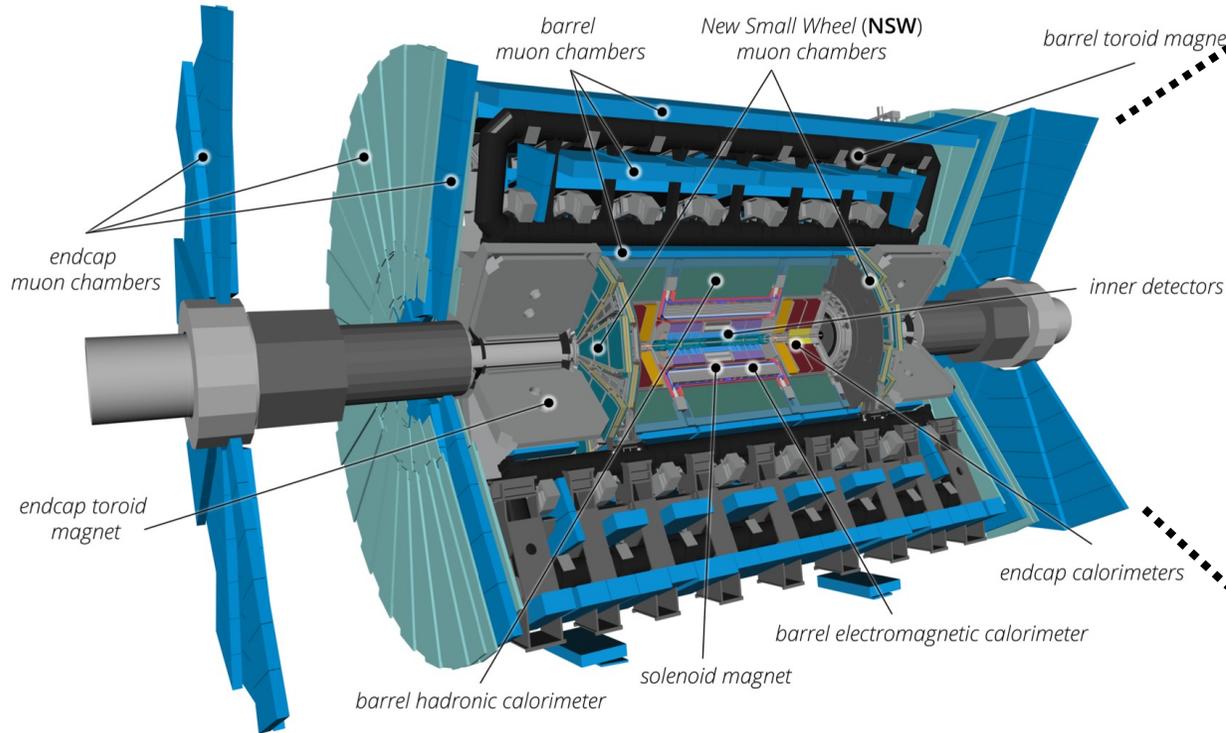


Amount of data collected in the collisions is called **Luminosity**



LHC collected **more than 450 fb⁻¹** of integrated luminosity until today since 2011

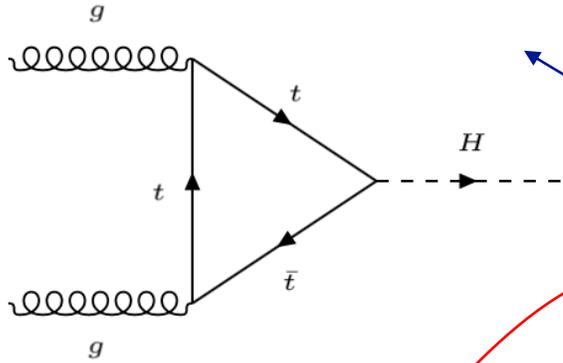
The ATLAS Detector



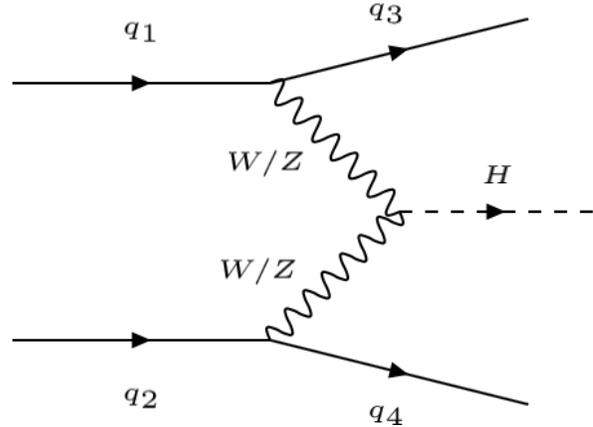
- Inner Detector
 - Vertex detector: pixel detector
 - Trackers: SCT, TRT
- Calorimeters
 - Electromagnetic (ECAL) : sampling Lar + lead
 - Hadronic (HCAL) : TileCal + LAr (endcap)
- Magnet System
 - Solenoid (2 T) + Toroid (0.5 – 1 T)
- Muon Spectrometer (MS)
 - Trackers: MDT, Micromegas
 - Trigger: RPC, TGC, sTGC

The Higgs boson production @ LHC

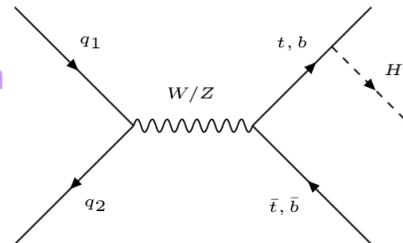
gluon-gluon fusion (ggF)



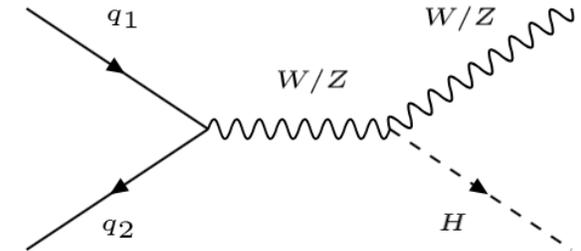
Vector boson fusion (VBF)



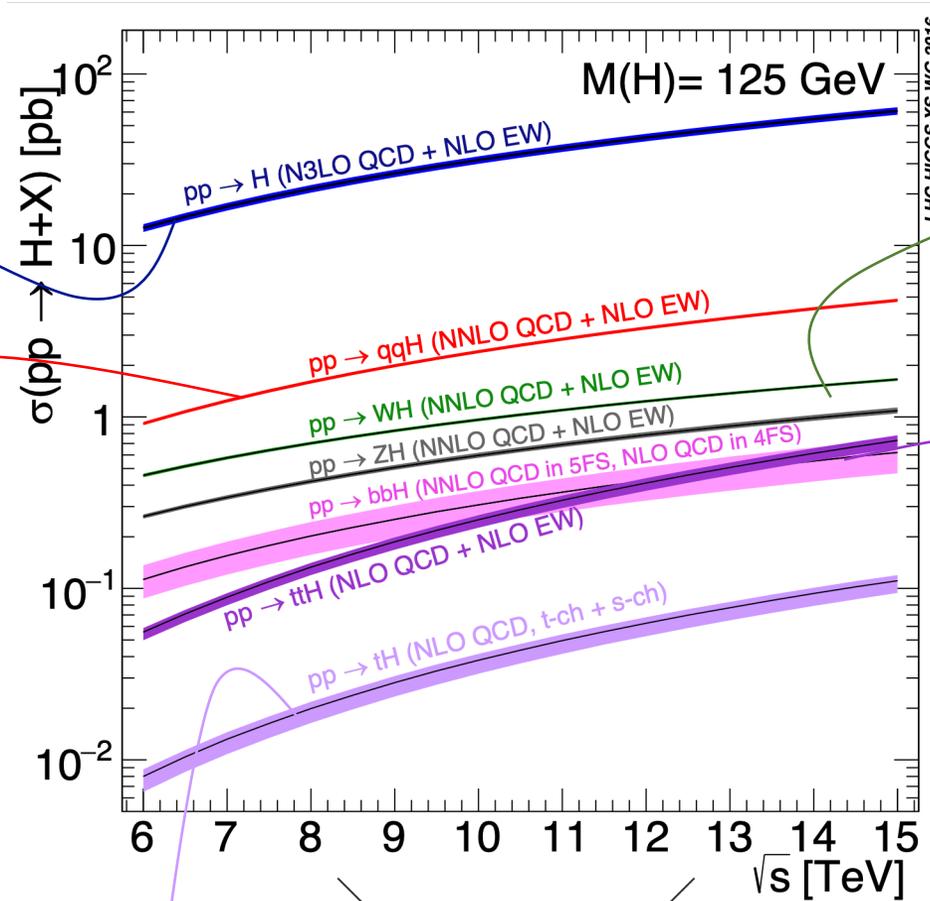
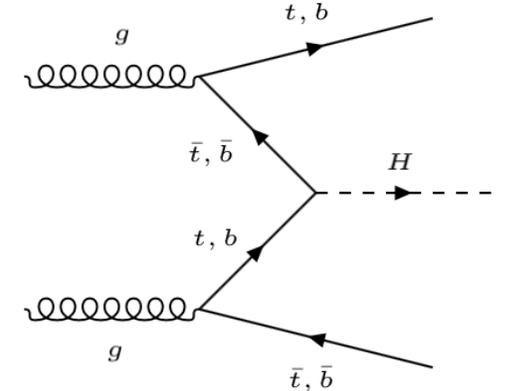
Associated production with single top (tH)



Associated production with a vector boson (VH)



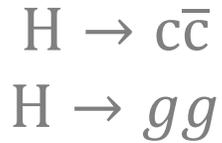
Associated production with top/bottom quark pair (ttH/bbH)



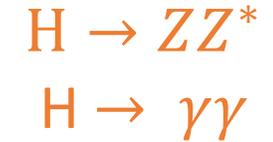
LHC HIGGS XS WG 2016

The Higgs boson decay modes

Hard to detect decays

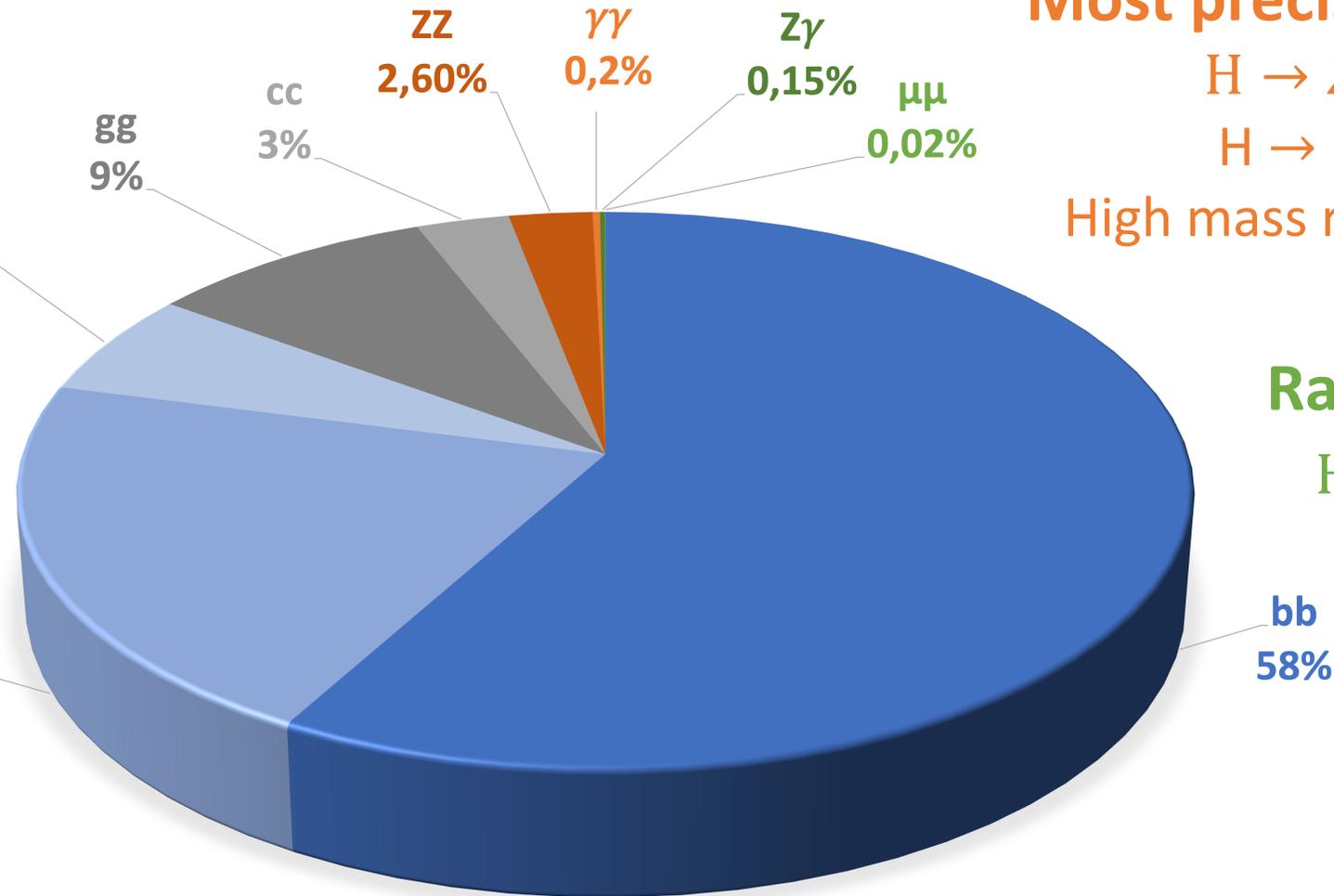
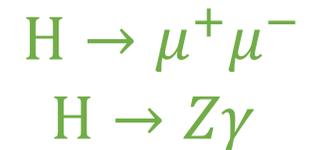


Most precise decays

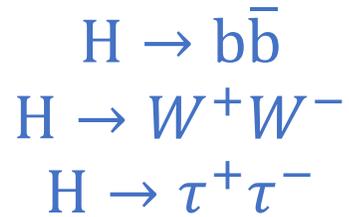


High mass resolution

Rare decays

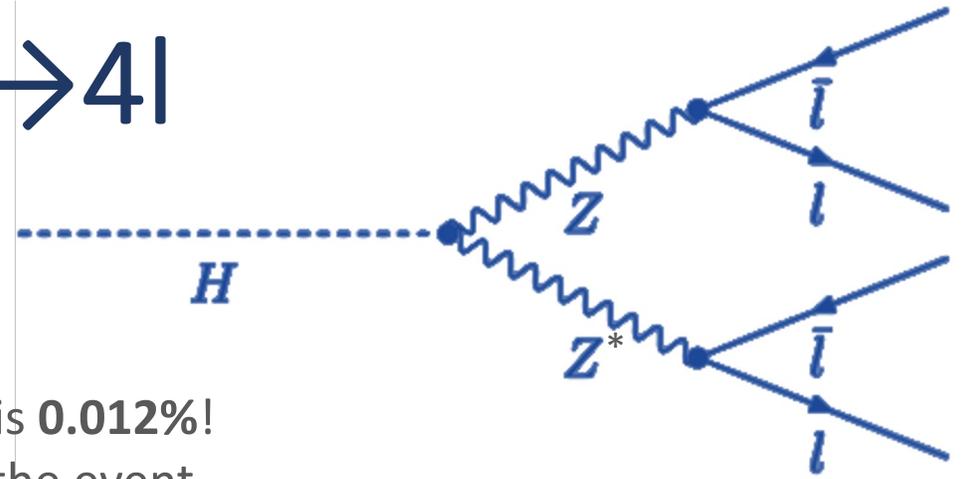


High probable decays



high backgrounds and
poor mass resolution

The Golden Channel: $H \rightarrow ZZ^* \rightarrow 4l$



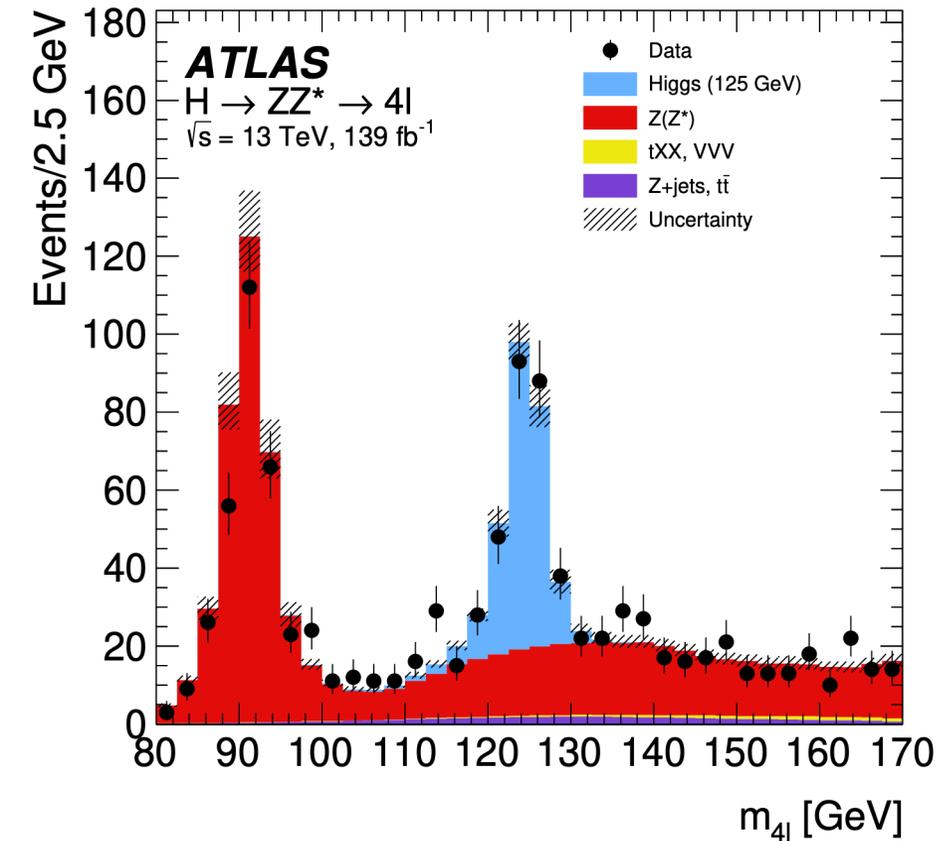
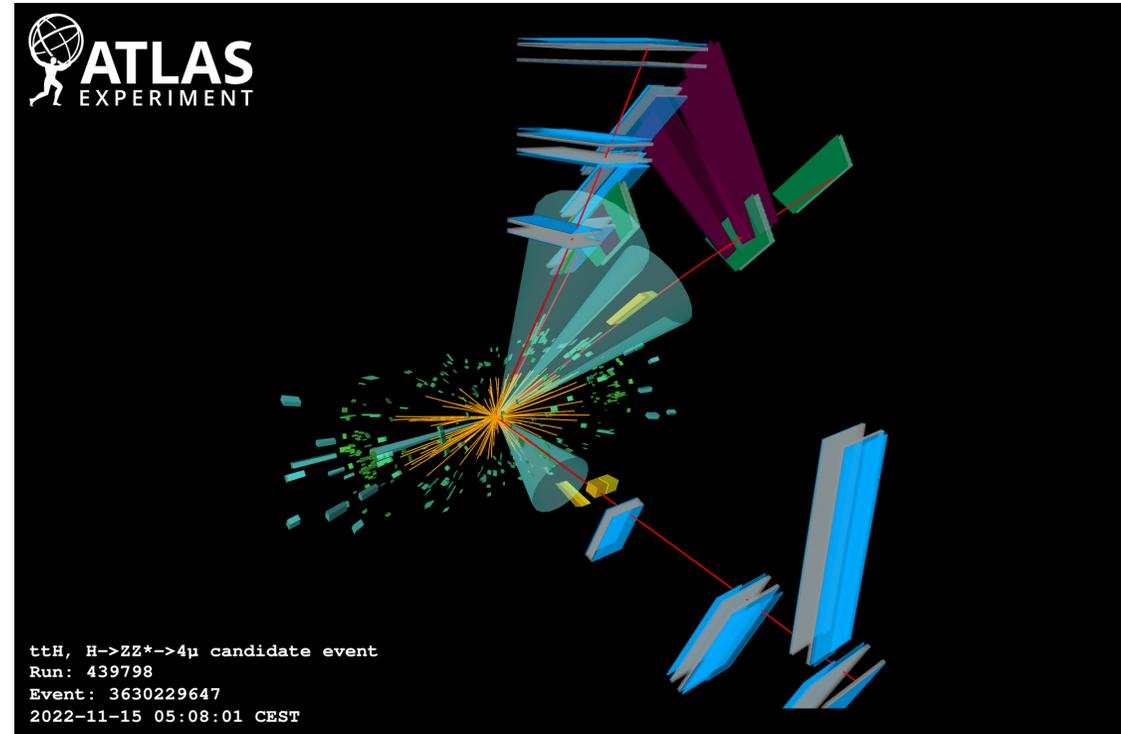
Final States: 4μ , $2\mu 2e$, $2e 2\mu$, $4e$

Fully leptonic final state

→ probability to detect it is **0.012%**!

→ very **clear signature** of the event

Very good mass resolution thanks to the excellent **lepton reconstruction** performance



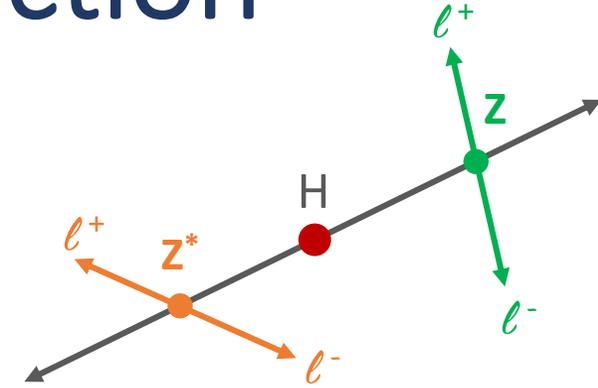
Very good **Signal/Background ratio** ~ 2

→ clearly identifiable over the background

$H \rightarrow ZZ^* \rightarrow 4l$: Higgs candidate selection

Two Same-Favour and Opposite-Sign (SFOS) lepton pairs

Leptons are required to be well separated and with $p_T > 20, 15, 10$ GeV and coming from a **common vertex**



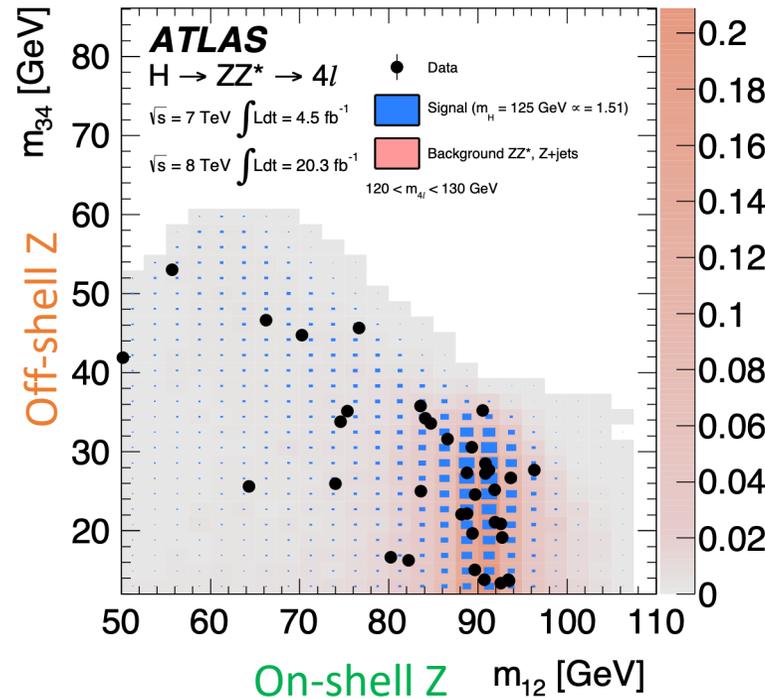
Mass requirements

On-shell Z

$$50 < m_{12} < 106 \text{ GeV}$$

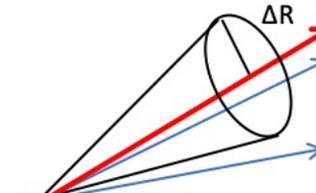
Off-shell Z

$$m_{\text{thr}}^* < m_{34} < 115 \text{ GeV}$$

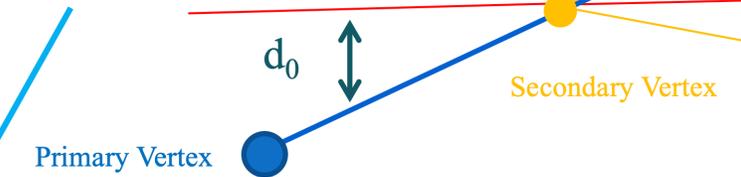


Lepton Isolation

Track- and calorimeter- based isolation requirements applied



Impact parameter d_0 significance cut



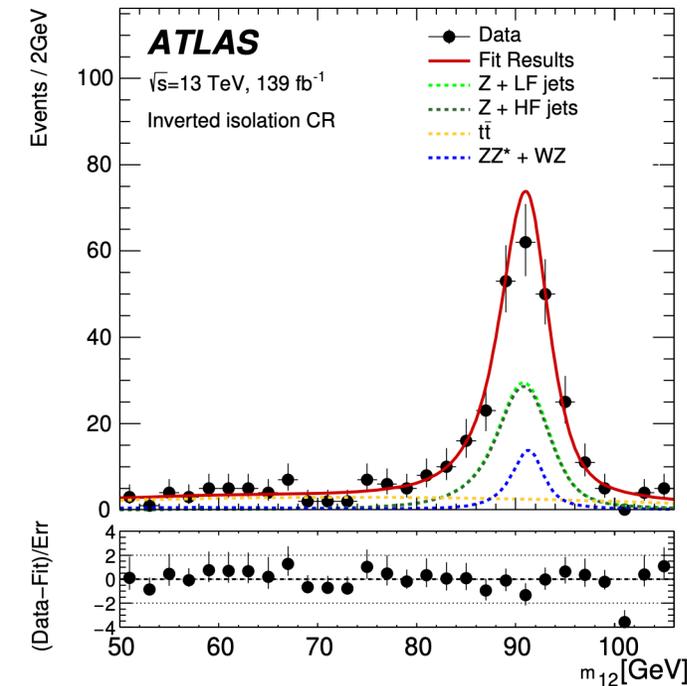
Suppress $t\bar{t}$ and Z+jets background

* $m_{\text{thr}} = 12$ GeV if $m_{4l} < 140$ GeV and rises linearly to 50 GeV for $m_{4l} = 190$ GeV.

H → ZZ* → 4l: background estimation

Irreducible Background (with 4 prompt leptons)

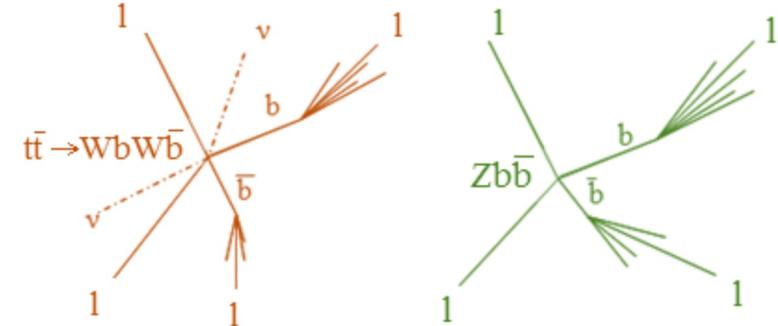
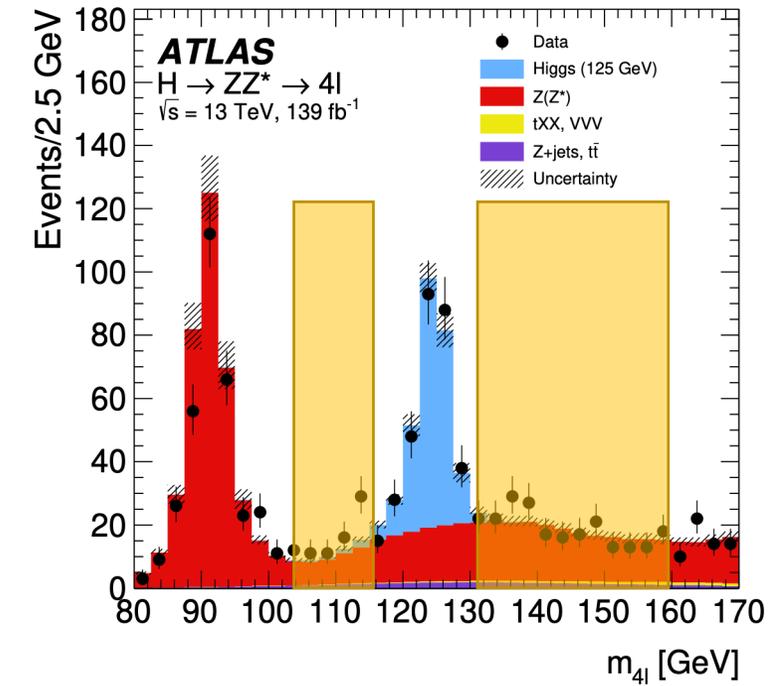
- ZZ* non-resonant production: 4 prompt leptons → estimated **from data** in the **mass sidebands [105-115] + [130-160] GeV**
- VVV and tXX from simulation



Reducible Background (tt̄, Z+jets, WZ)

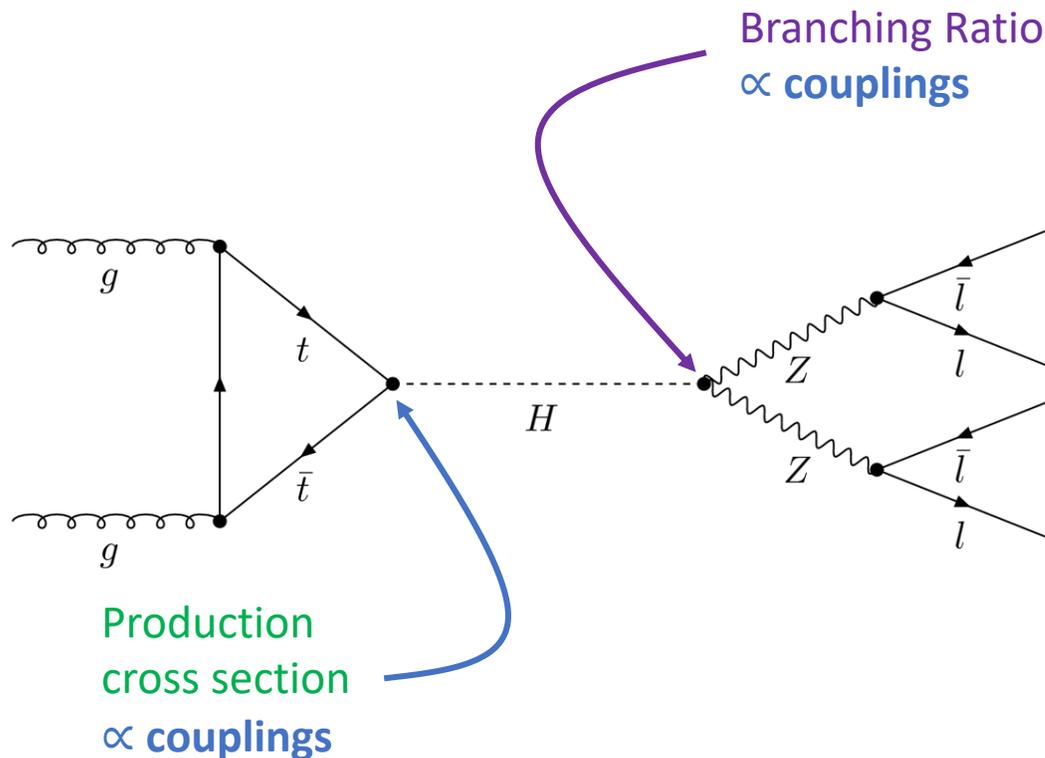
2 prompt leptons (Z) + 2 leptons from semi-leptonic decays (b- or c- quark)

- **Control Regions (CR)** defined inverting or relaxing Higgs candidate requirements
- **Data-driven estimation** in the CRs of the single components → **extrapolated into the Signal Region**



Higgs Cross Sections as Precision Observables

Cross section measurements are the core of the Higgs boson precision physics



σ : cross sections as probability to produce Higgs bosons

BR : Branching Ratios as probability that the Higgs decay in a specific channel

What do we measure? $\sigma \times BR$

How do we measure it?

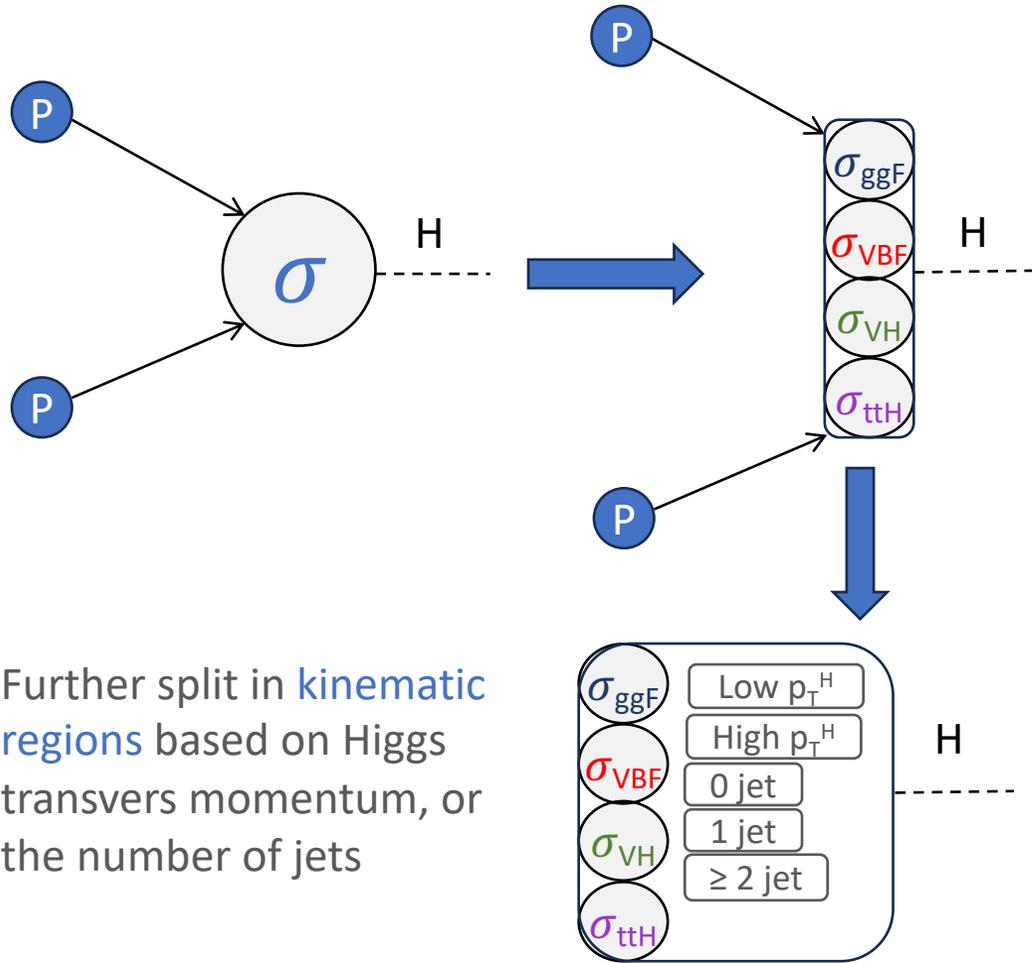
$$N_{obs} = \sigma \times BR \times L \times \epsilon$$

\propto selection efficiency (pointing to ϵ)
 \propto Luminosity (pointing to L)

Experimentally, we can only measure number of events selected as Higgs candidate

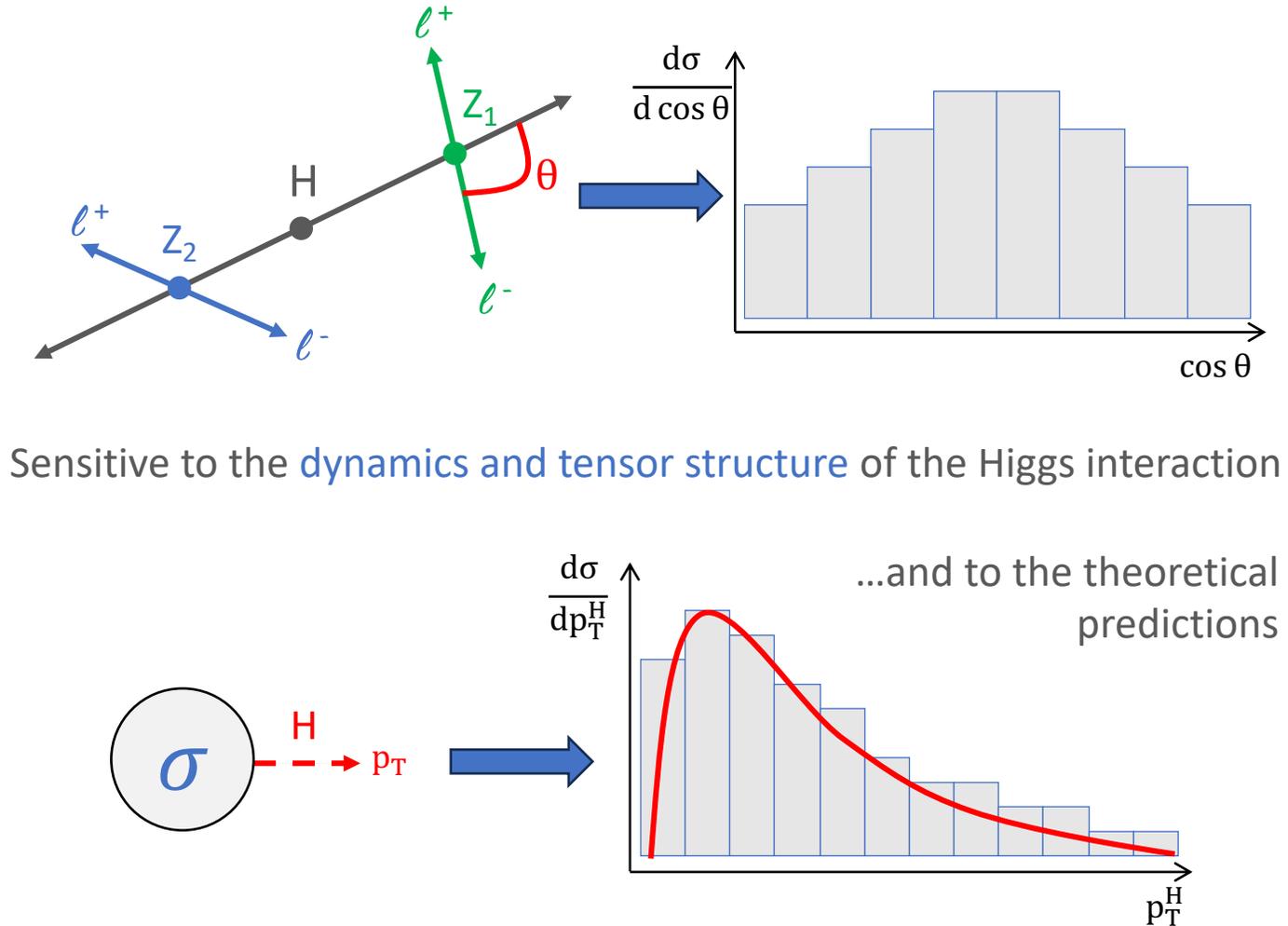
From inclusive to differential measurements

Measure $\sigma \times \text{BR}$ for different **Production Modes**



Further split in **kinematic regions** based on Higgs transverse momentum, or the number of jets

Measure $\sigma \times \text{BR}$ in bins of **Final-State Observables**



Sensitive to the **dynamics and tensor structure** of the Higgs interaction

...and to the theoretical predictions

Interpretations frameworks

Kappa-framework

Rescale Standard Model couplings

$$\sigma_i \times \text{BR}^f = \frac{\sigma_i(\vec{\kappa})\Gamma^f(\vec{\kappa})}{\Gamma_H}$$

$$\kappa_j^2 = \frac{\sigma_j}{(\sigma_j)_{SM}} \quad \kappa_j^2 = \frac{\Gamma^j}{(\Gamma^j)_{SM}}$$

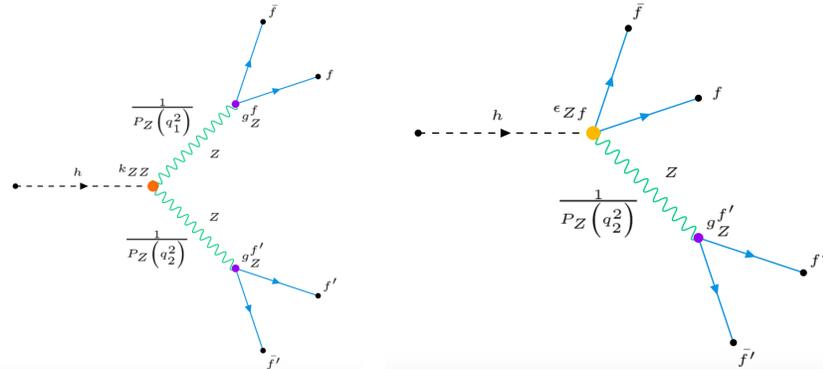
- Directly links rates to coupling strengths
- Simple, intuitive, widely used
- Sensitive only to effects on the rate

Pseudo - Observables

Parametrize momentum expansion of the on-shell Higgs decay amplitude

$$\frac{d\Gamma}{dm_{12}dm_{34}} \propto \sum_{f,f'} |F^{ff'}(m_{12}^2, m_{34}^2)|^2$$

$$F_1^{ff'}(q_1^2, q_2^2) = \kappa_{ZZ} \frac{g_Z^f g_Z^{f'}}{P_Z(q_1^2) P_Z(q_2^2)} + \frac{\epsilon_{Zf}}{m_Z^2} \frac{g_Z^{f'}}{P_Z(q_2^2)} + \frac{\epsilon_{Zf'}}{m_Z^2} \frac{g_Z^f}{P_Z(q_1^2)}$$



- Test kinematic structure of the interaction
- Easy interpretation of the effective coupling POs

Effective Field Theory

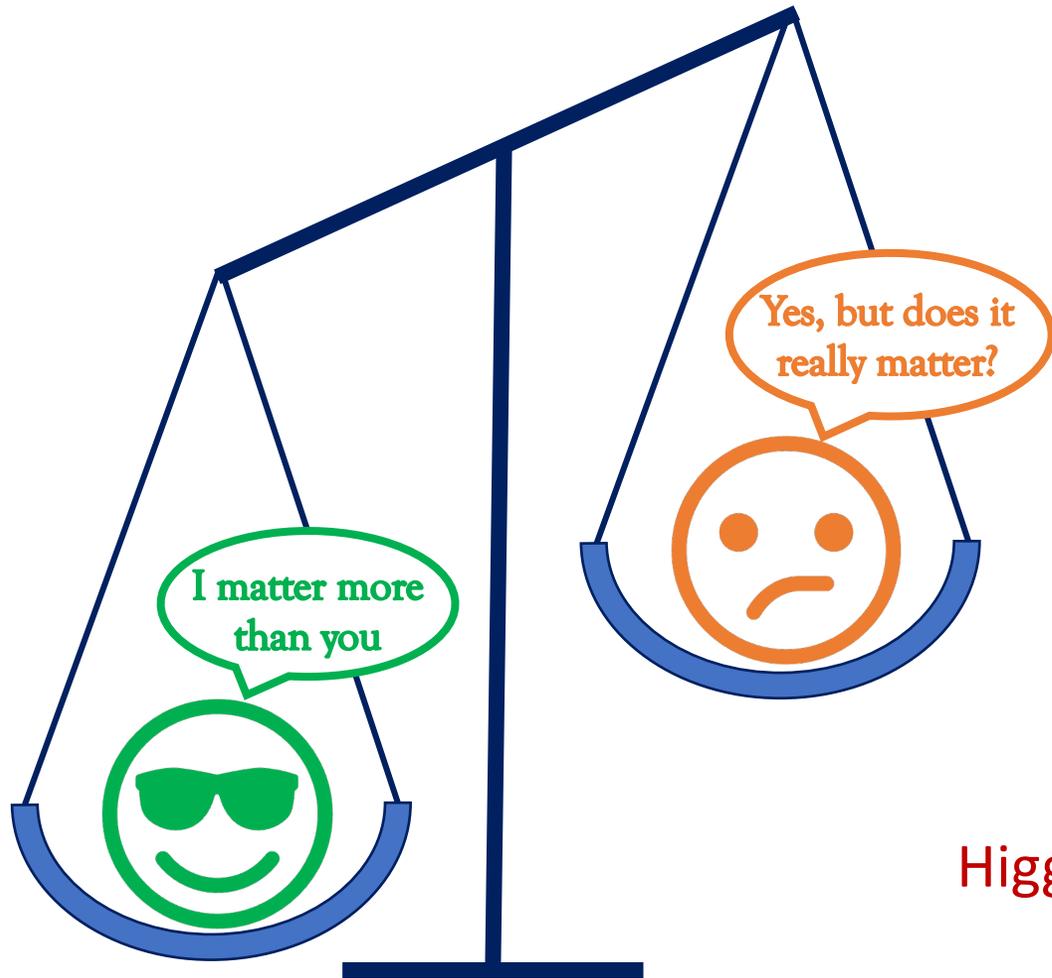
Add higher-dimensional operators to SM Lagrangian

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i^{(d)}}{\Lambda^{(d-4)}} \mathcal{O}_i^{(d)} \quad \text{for } d > 4$$

- Directly access possible anomalous couplings (*Wilson coefficients*) and new heavy physical states
- Interpret both shapes and rates in a field-theory framework

Different frameworks extract *different kinds* of information about Higgs couplings.

CP violation: anomalous Higgs couplings



The asymmetry between the **matter** and **antimatter** content in the universe is one of the **unsolved problems in Physics**



It implies the **violation of the charge conjugation and parity symmetry**

Standard Model can explain just a small level of the CP-violation required

→ Other sources of CP-violation must exist **Beyond the Standard Model**

Higgs is a good candidate to look at!

SM: Spin 0 and CP-even particle

We can look at the tensor structure of the Higgs interactions to identify possible mixed CP-even and CP-odd state

From Precision Measurements...

Fiducial and Production Cross Sections

Fiducial Cross Sections measurement

Provide cross section measurements in the most model independent way

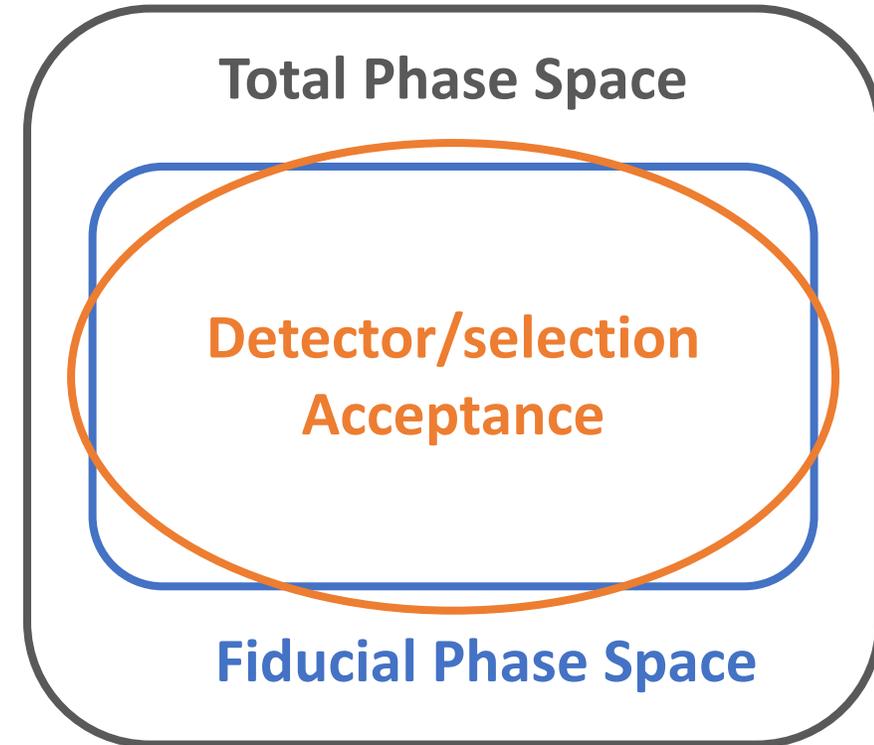
- **Fiducial phase space at particle-level** definition based on **detector and analysis selection acceptance** to minimize the extrapolation effects. The **fiducial cross section $\sigma_{fid} \cdot BR$** is defined as:

$$\sigma_{fid} \cdot BR = \sigma_{tot} \cdot BR \cdot A \quad A = \frac{N_{fiducial}}{N_{total}}$$

where:

BR = Branching ratio

A = acceptance



Fiducial Cross Sections measurement

Provide cross section measurements in the most model independent way

- **Fiducial phase space at particle-level** definition based on **detector and analysis selection acceptance** to minimize the extrapolation effects. The **fiducial cross section $\sigma_{fid} \cdot BR$** is defined as:

$$\sigma_{fid} \cdot BR = \sigma_{tot} \cdot BR \cdot A \quad A = \frac{N_{fiducial}}{N_{total}}$$

where:

BR = Branching ratio

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- **Correct for detector level effects**, efficiencies and resolution, defining the correction factor entering in the fiducial cross section extraction:

$$\sigma_{fid} \cdot BR = \frac{N_{signal}}{C_F \cdot L_{int}}$$

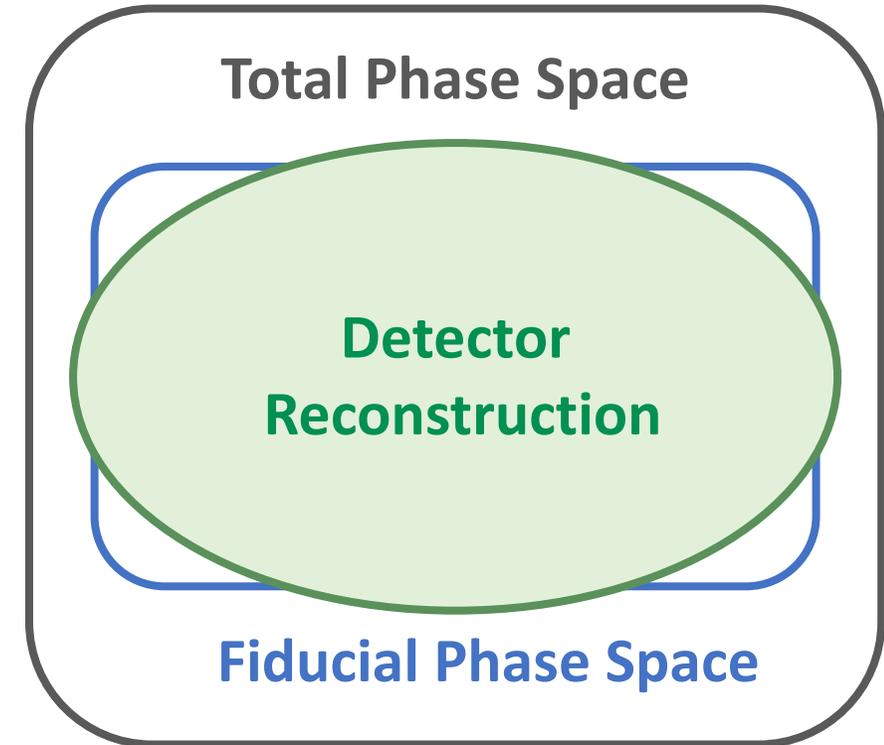
where:

L_{int} = integrated luminosity

C_F = correction factor

N_{signal} = number of signal events extracted fitting the observable able to discriminate signal vs background.

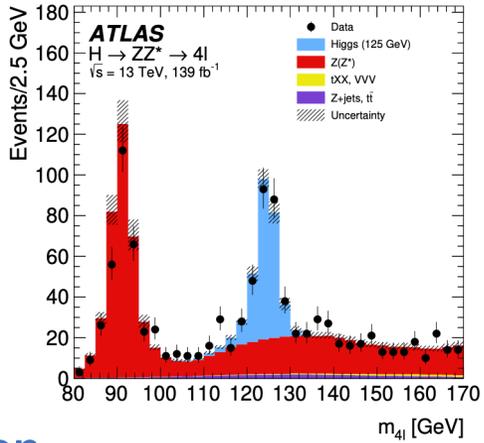
$$C_F = \frac{N_{reconstructed}}{N_{fiducial}} \longrightarrow$$



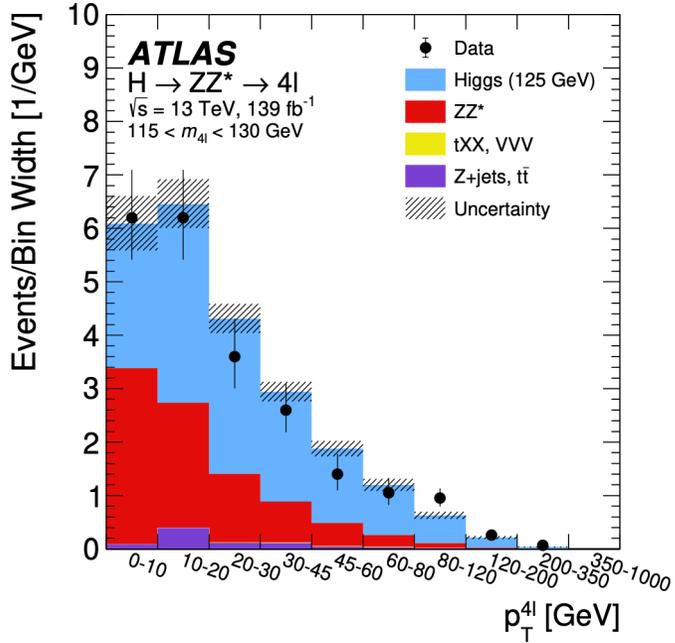
We **unfold** the reconstructed distribution of a given observable to estimate the **particle-level spectrum**

Fiducial Differential Cross Sections measurement

Discriminant
Signal vs
Background
Observable is m_{4l}
template fit in each
variable bin

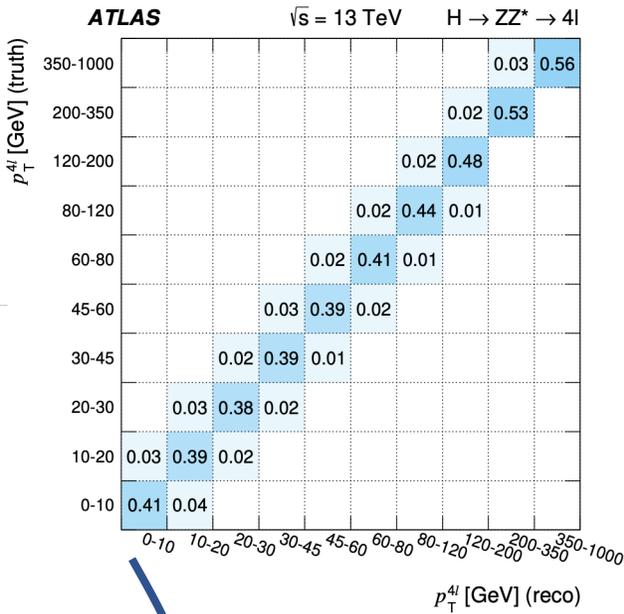
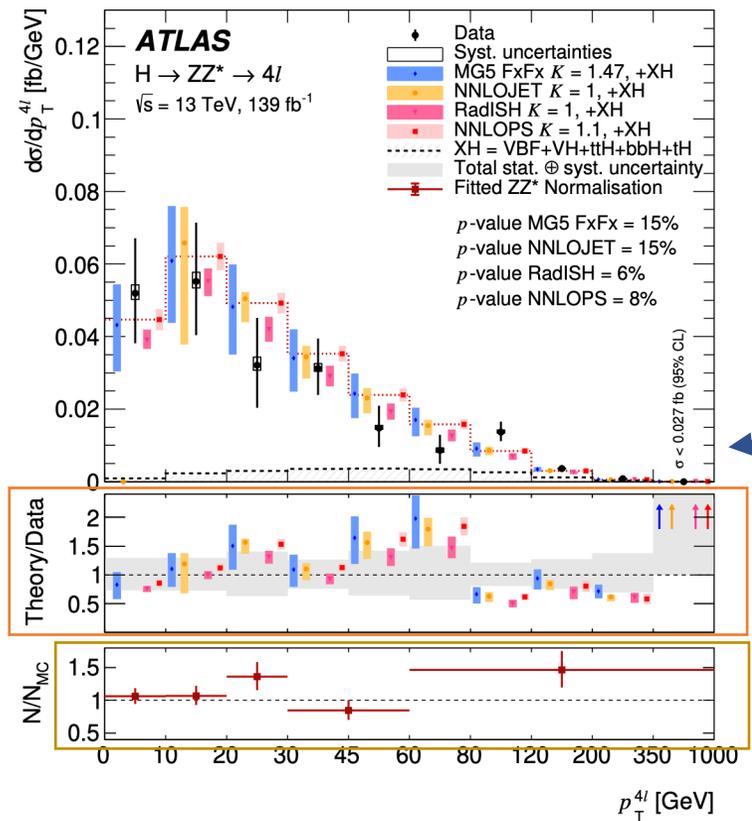


Reconstructed distribution



Compatibility
between data and
different theoretical
predictions

Particle – level distribution



Unfolding correcting
for detector effects:
acceptance and fiducial
efficiency

ZZ* background
estimation from data

Fiducial Differential Cross Sections measurement

Observables studied are sensitive to different Higgs boson properties or theoretical predictions

Higgs kinematic variables (some)

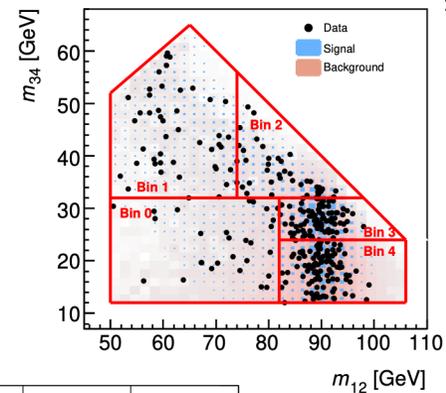
- p_T^{4l} : perturbative QCD, light quark coupling
- y^{4l} : parton density function
- $m_{12}, m_{34}, \cos\theta_1, \cos\theta_2, \phi, \phi_1, \cos\theta^*$: Sensitive to spin/parity properties in the decay vertex

Jet – related variables (some)

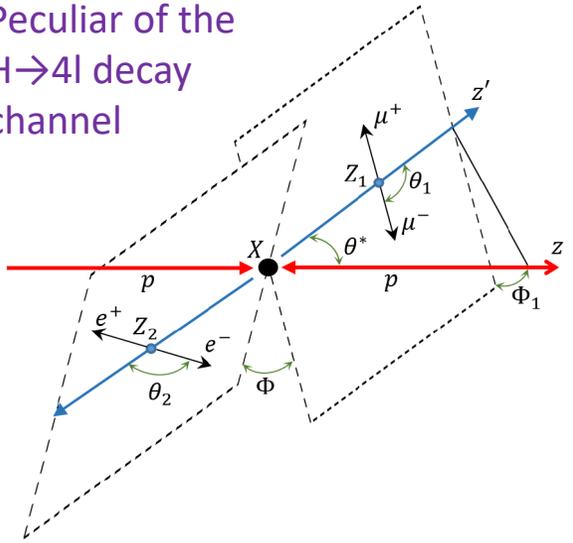
- $N_{jet}, m_{jj}, \Delta\eta_{jj}$: different production mode
- p_T^{j1}, p_T^{j2} : quark gluon radiation
- $\Delta\phi_{jj}$: Sensitive to spin/parity properties in the VBF production vertex

2D differential distributions (some)

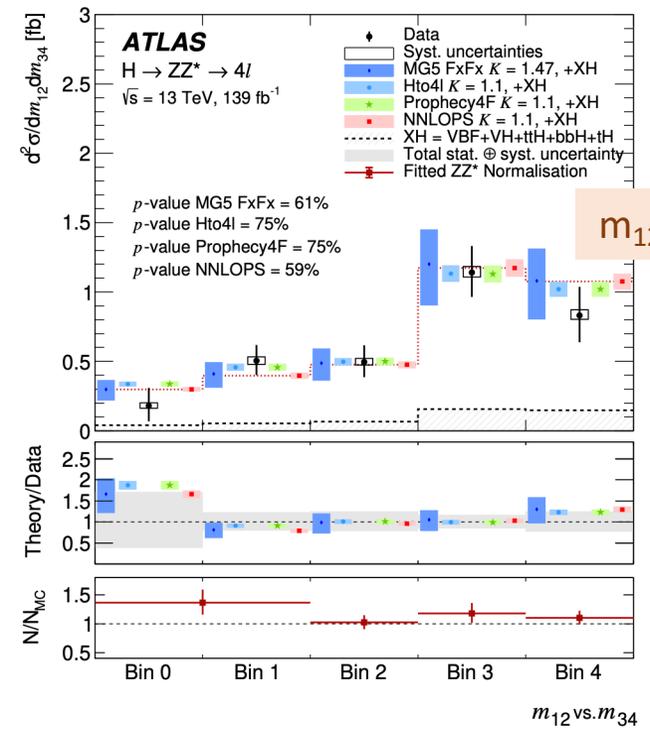
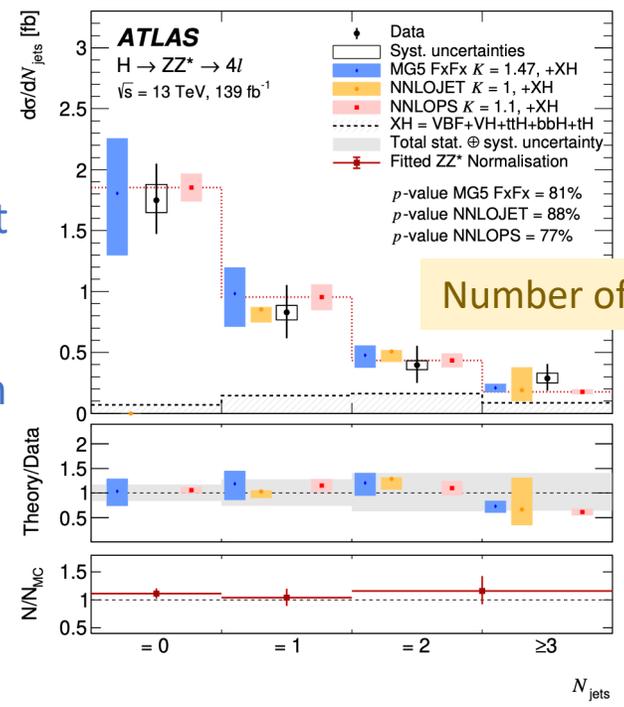
- $m_{12}m_{34}$: modification
- HZZ vertex



Peculiar of the $H \rightarrow 4l$ decay channel



Run 2 provided very large dataset of 140 fb^{-1} increasing the sensitivity in each differential bin



Selection of results from Run 2 paper

Fiducial Differential Cross Sections measurement

Observables studied are sensitive to different Higgs boson properties or theoretical predictions

Higgs kinematic variables (some)

$m_{12}, m_{34}, \cos\theta_1, \cos\theta_2, \phi, \phi_1, \cos\theta^*$:
Sensitive to spin/parity properties in the decay vertex

Jet – related variables (some)

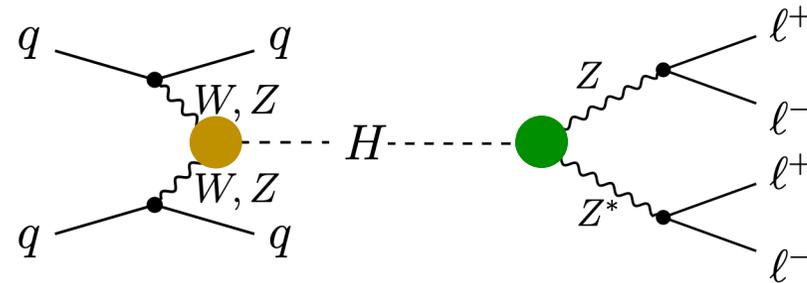
$\Delta\phi_{jj}$: Sensitive to spin/parity properties in the VBF production vertex

Lepton and jets kinematic information can be combined in a much more powerful observable sensitive to CP effects → **Optimal Observable**

The **matrix element** of the process is the sum of a SM CP-even contribution plus a BSM CP-odd one

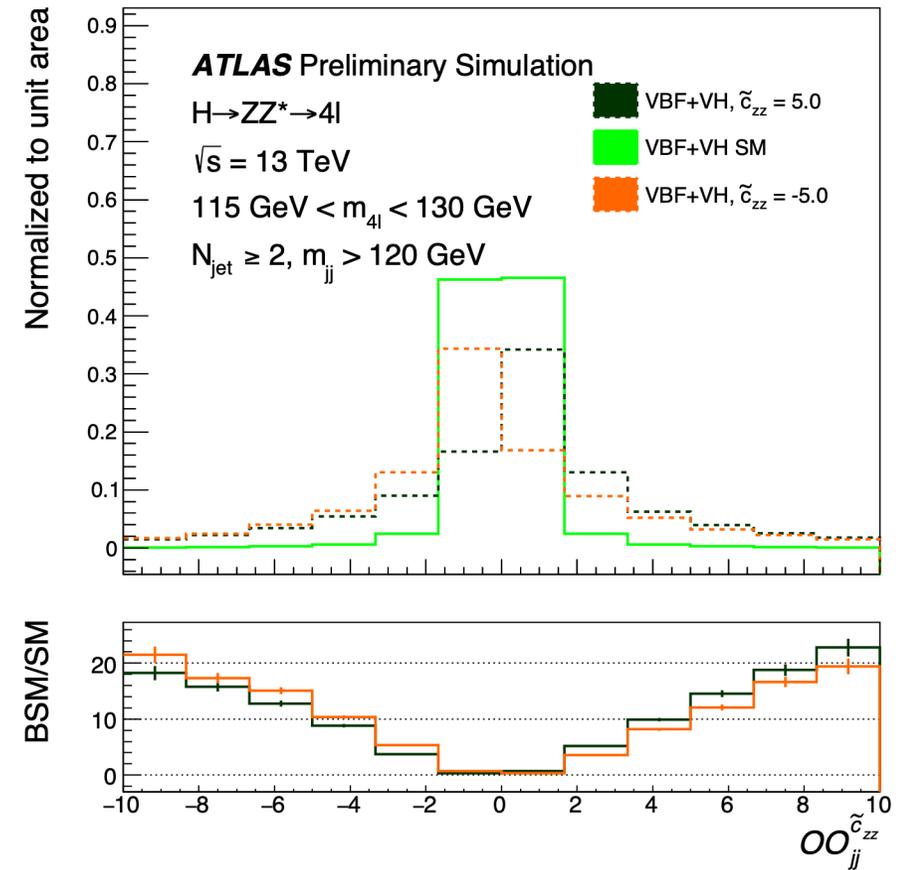
$$|\mathcal{M}|^2 = \left| \mathcal{M}_{SM} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{M}_{BSM,i} \right|^2$$

Interference term become our Optimal Observable



$$OO = \frac{2\Re(\mathcal{M}_{SM}^* \mathcal{M}_{BSM})}{|\mathcal{M}_{SM}|^2}$$

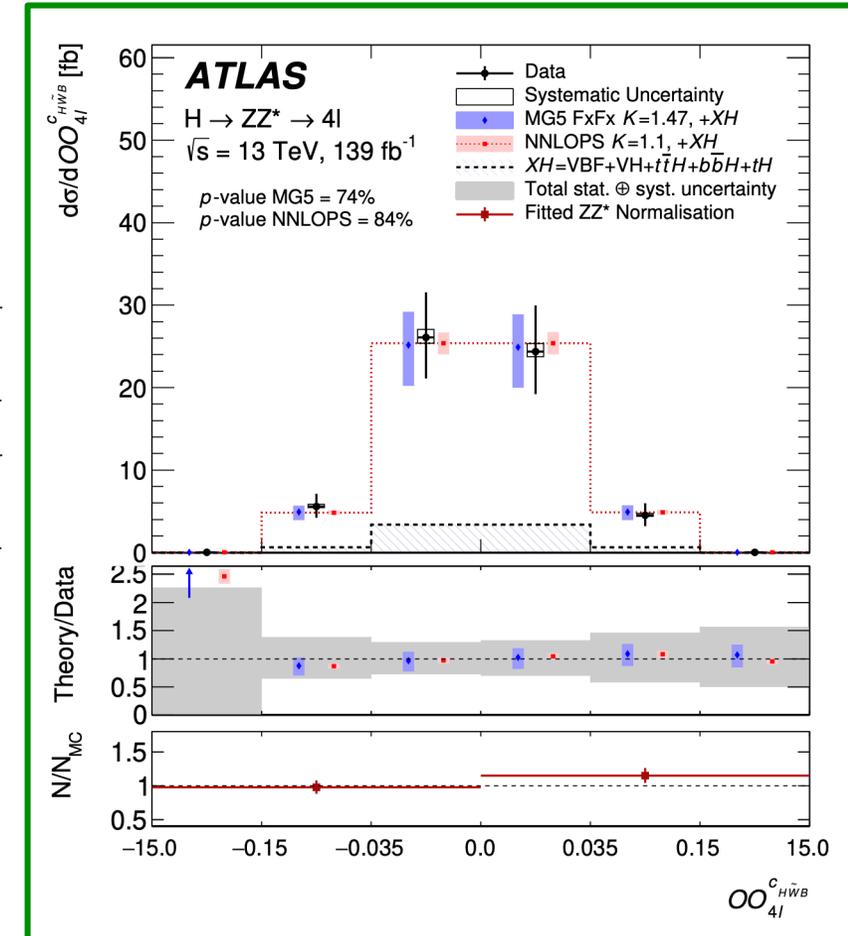
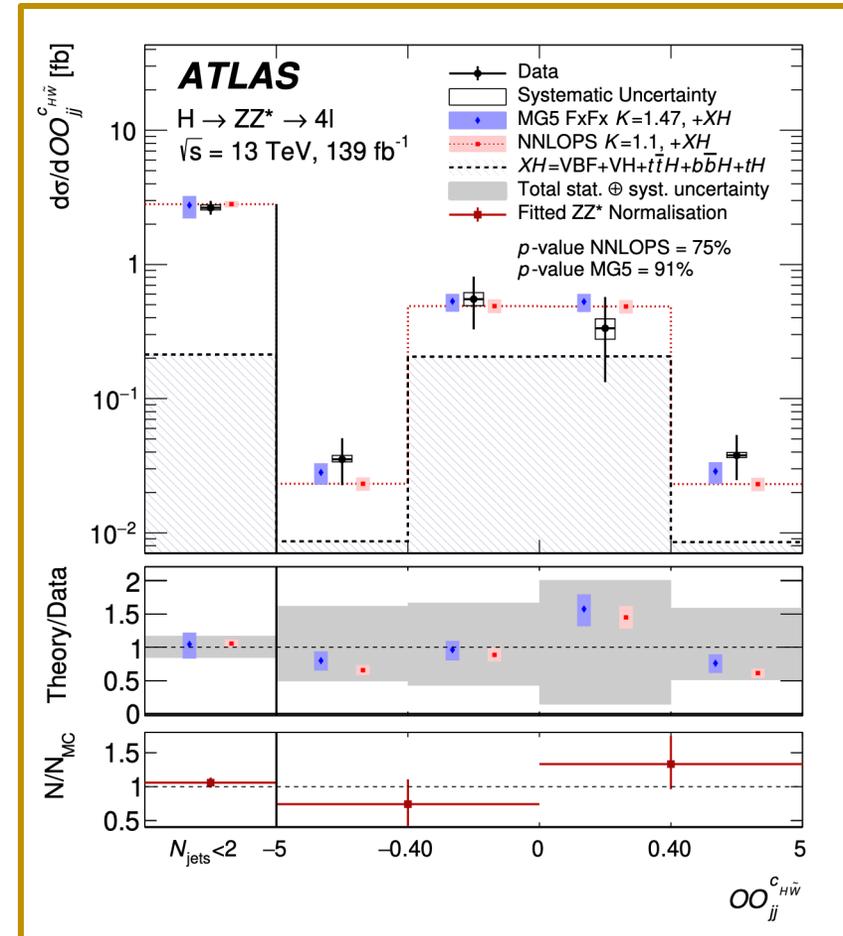
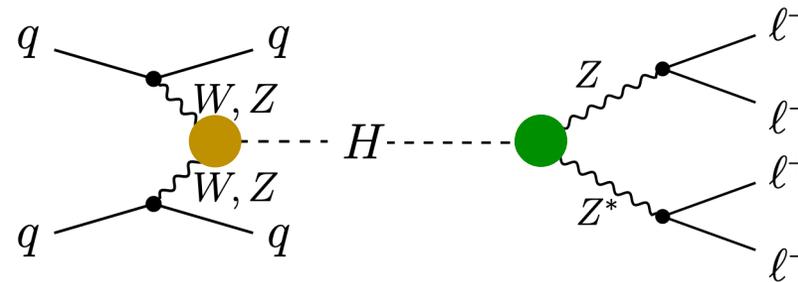
The interference become asymmetric in presence of CP-odd contribution in the coupling → sensitive to the sign



Fiducial Differential Cross Sections measurement

Observables studied are sensitive to different Higgs boson properties or theoretical predictions

Different Optimal Observables can be studied looking for anomalous CP-odd coupling at the **production** and/or **decay vertex**

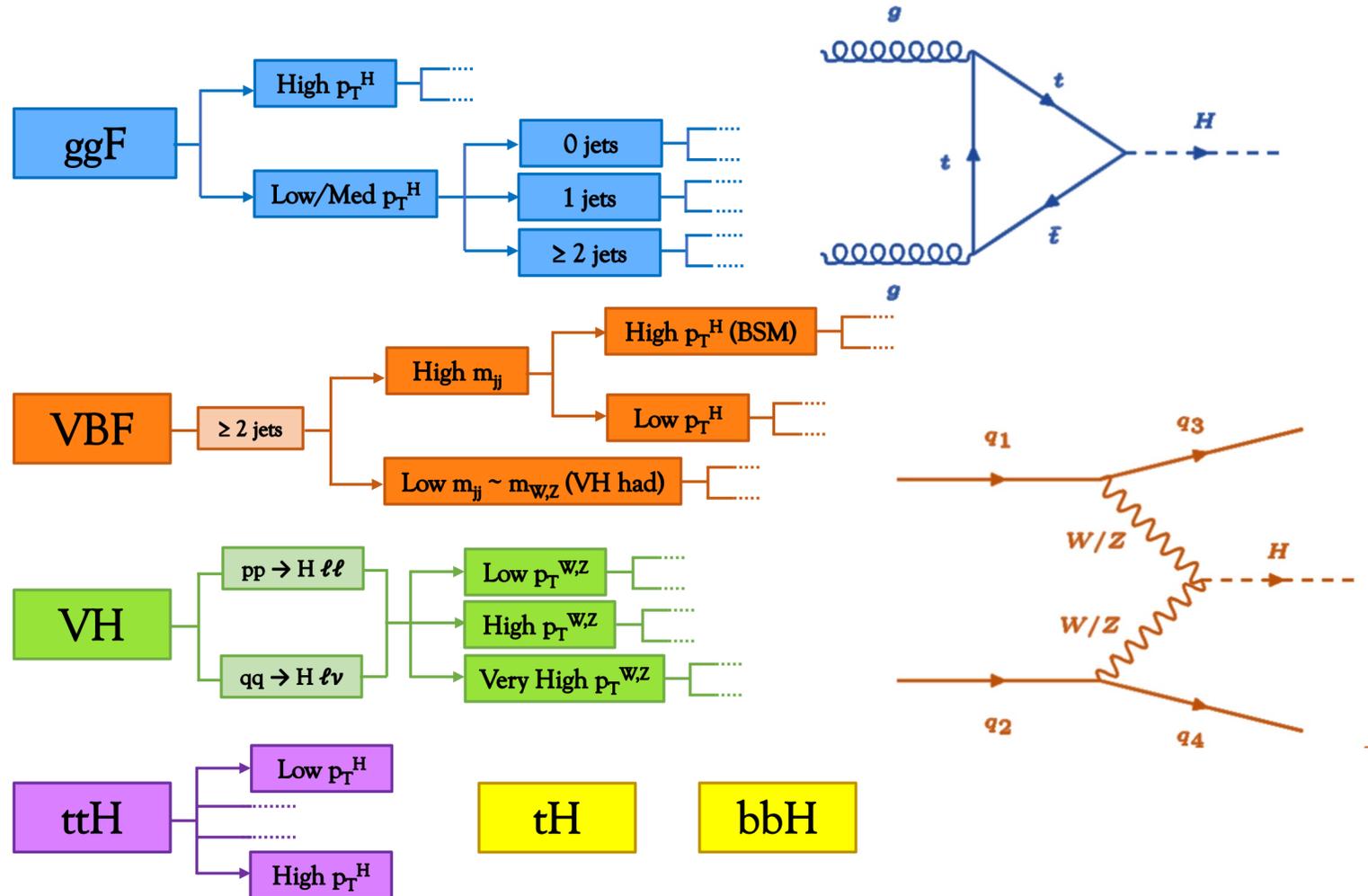
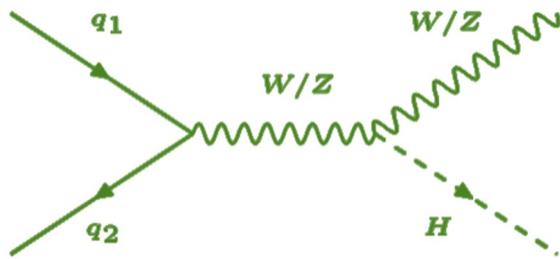


Production Cross Section measurements: STXS

Simplified Template Cross Section Framework

Exclusive regions in the Higgs phase space of the Higgs production processes, based on the kinematics of the Higgs and of the particles/jets produced in association

Maximize the experimental sensitivity to possible BSM effects

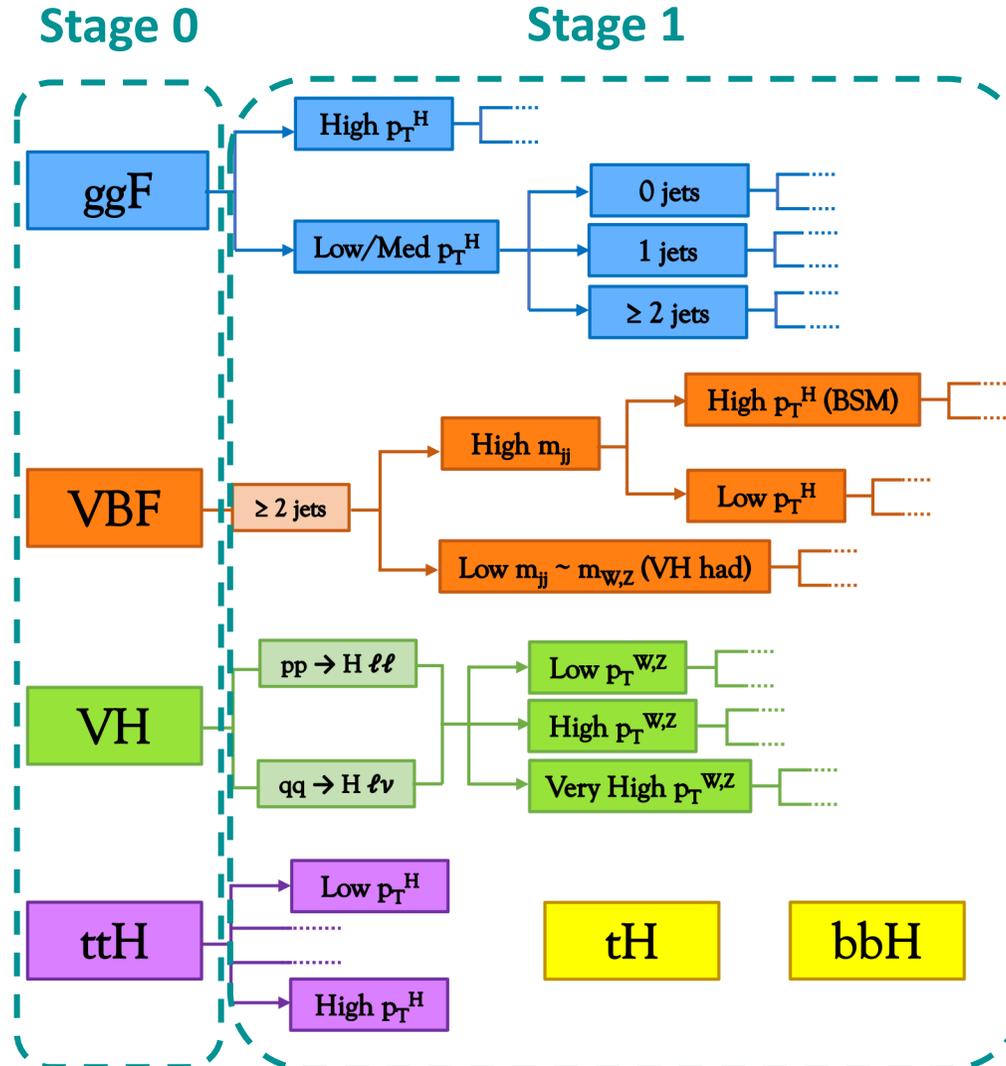


Production Cross Section measurements: STXS

Exclusive regions in the Higgs phase space of the Higgs production processes, based on the kinematics of the Higgs and of the particles/jets produced in association

Maximize the experimental sensitivity to possible BSM effects

Different STXS Stages definition at particle-level, increasingly fine granularity



Production Cross Section measurements: STXS

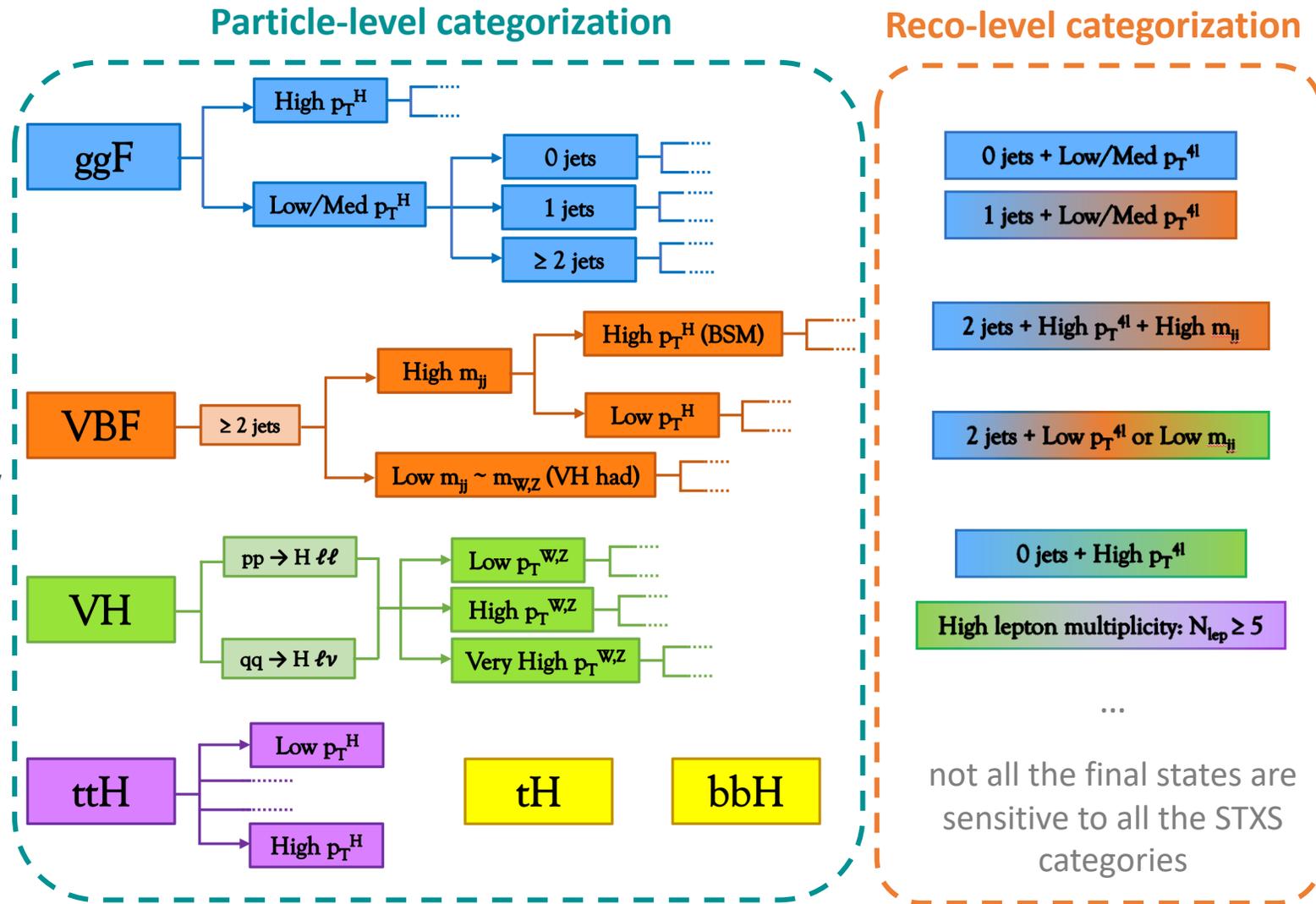
Exclusive regions in the Higgs phase space of the Higgs production processes, based on the kinematics of the Higgs and of the particles/jets produced in association

Maximize the experimental sensitivity to possible BSM effects

Different STXS Stages definition at particle-level, increasingly fine granularity

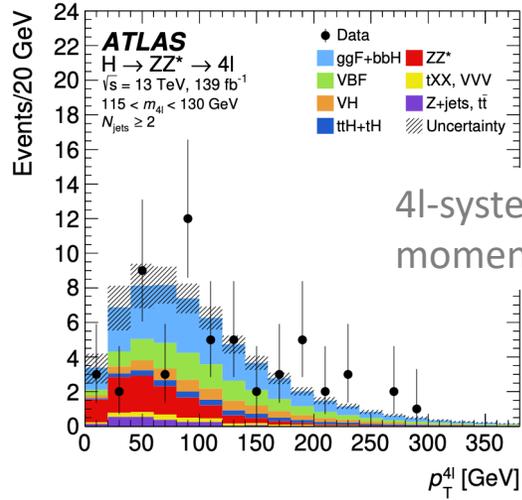
Reconstructed event categorization as close as possible to this definition → minimize model-dependent extrapolation

Different production modes can populate a reco-level categories → discriminant observables to extract production cross sections

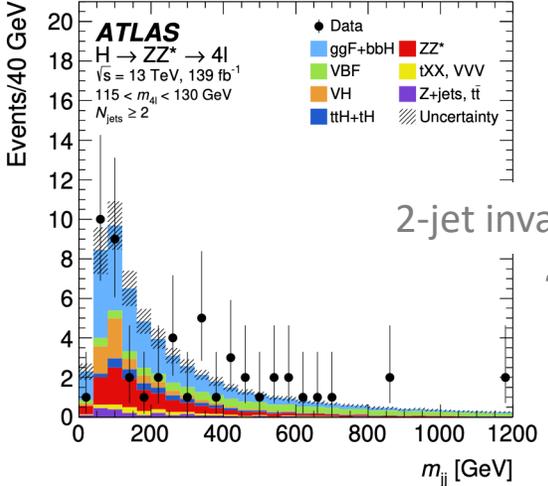


Production Cross Section measurements: STXS

Discriminant Observable between different production modes: **Neural Networks**

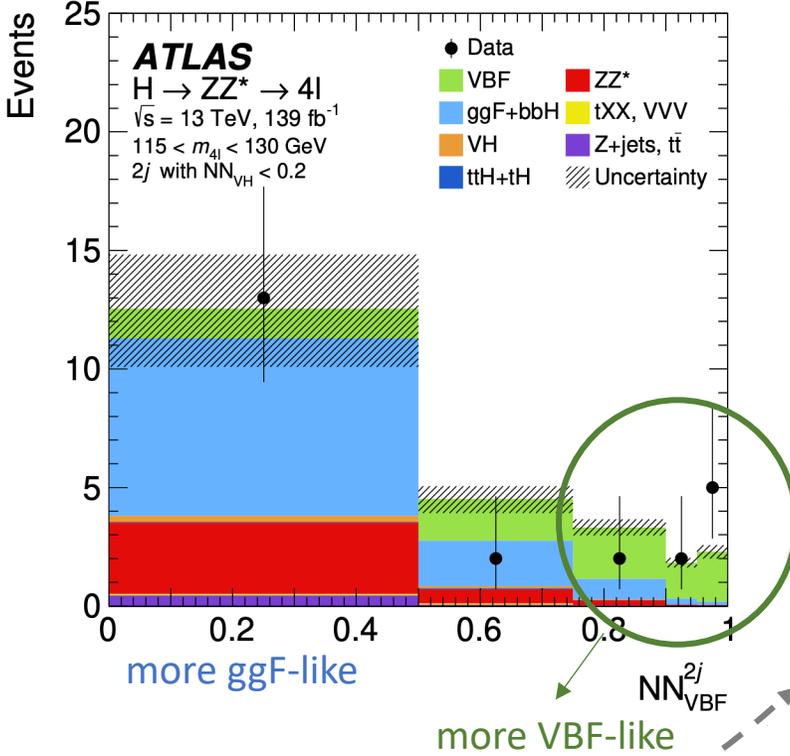


4l-system transverse momentum in for $N_{\text{jets}} \geq 2$



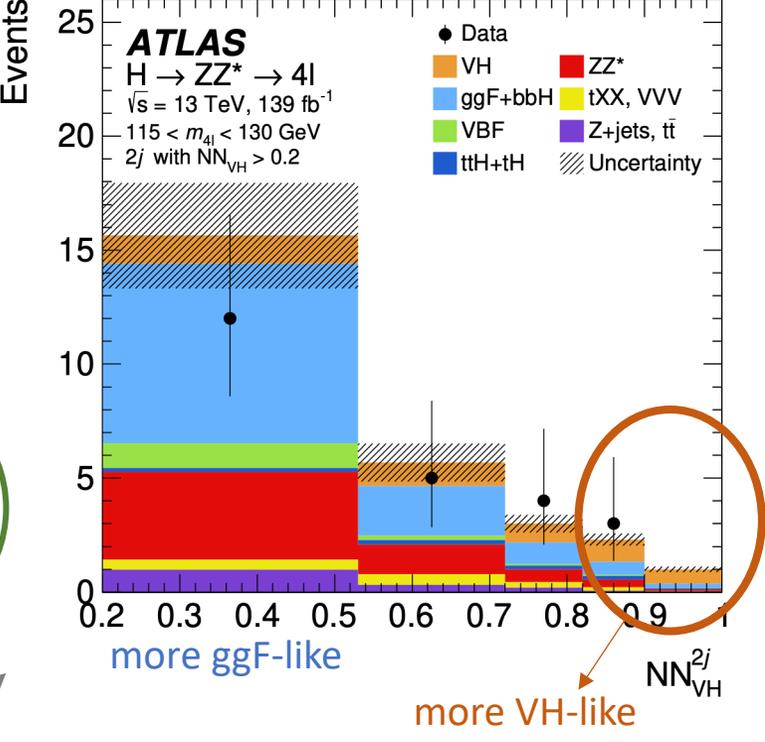
2-jet invariant mass

Only kinematic observables not powerful enough as discriminant... but can be used as input of the NN



more ggF-like

more VBF-like



more ggF-like

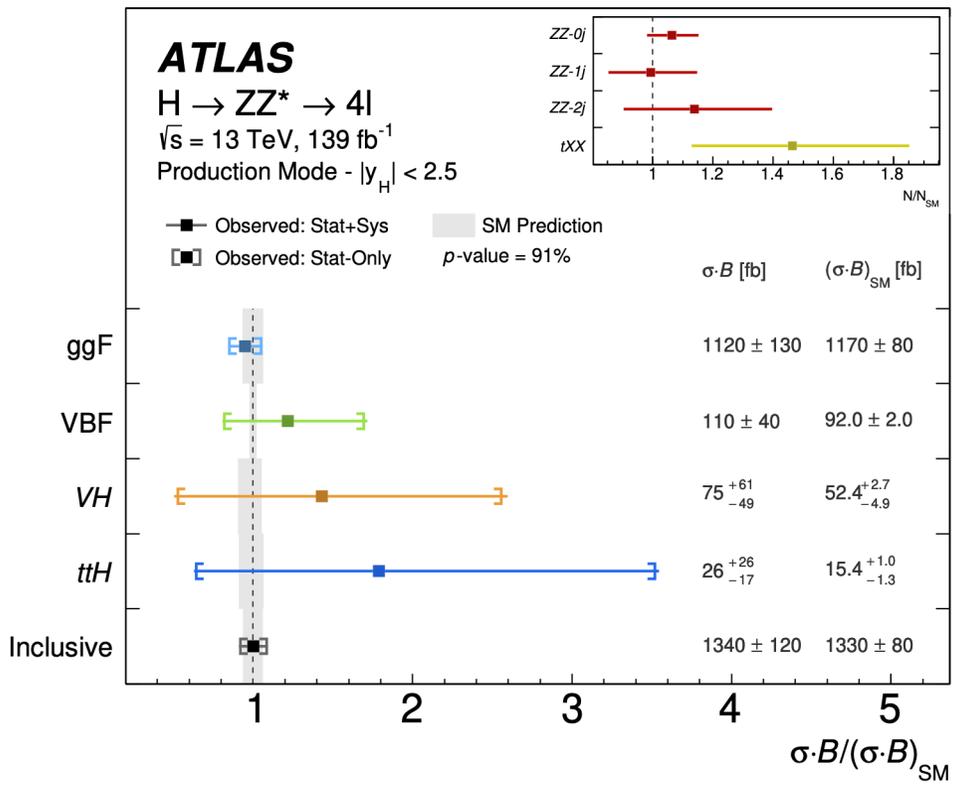
more VH-like

Production Cross Section measurements: STXS

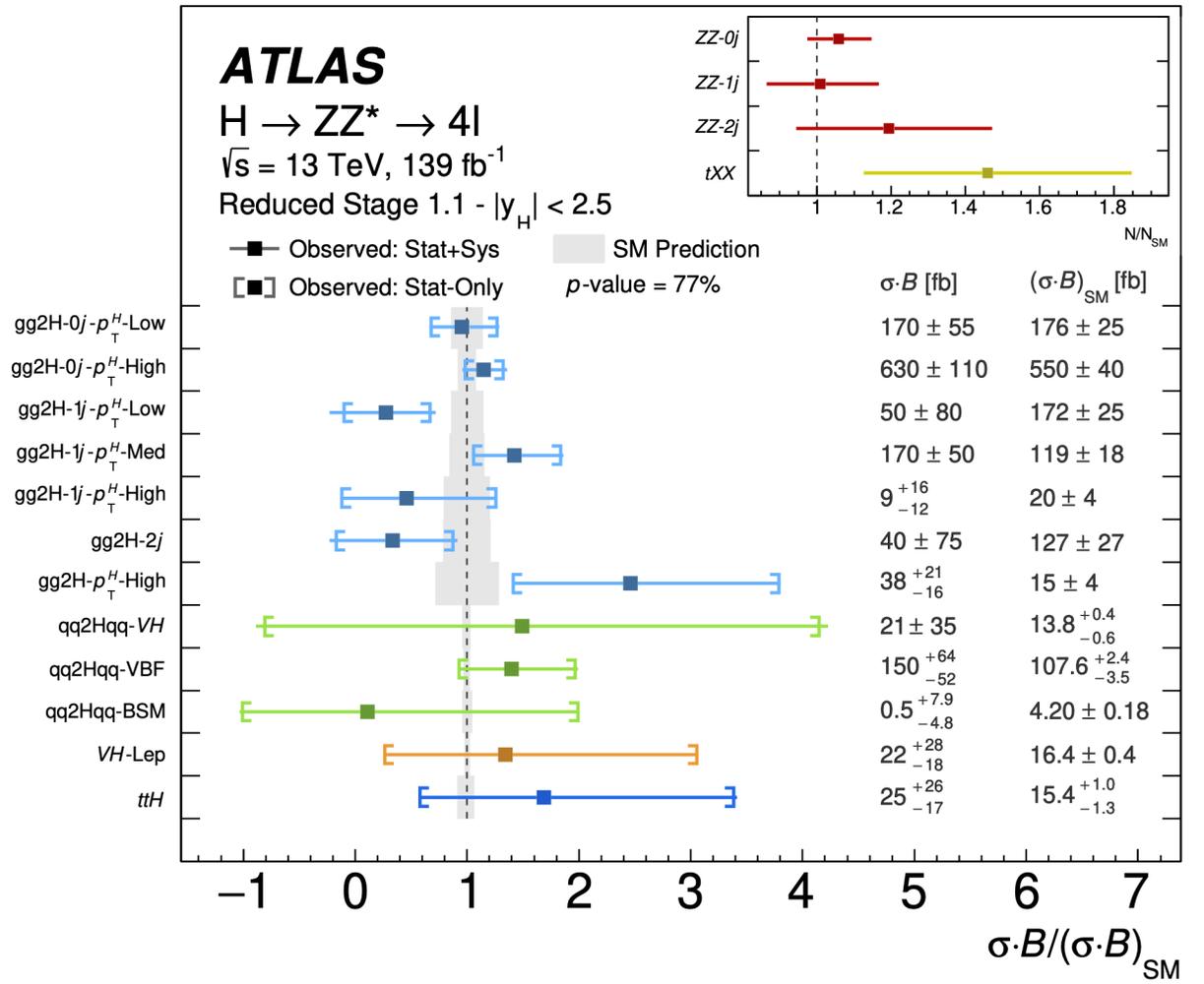
Results from Run 2

Sensitivity is statistically limited, but large Run 2 dataset allowed to probe more STXS regions than previous analyses

Production Mode Stage (Stage 0)



(Reduced) Stage 1.1

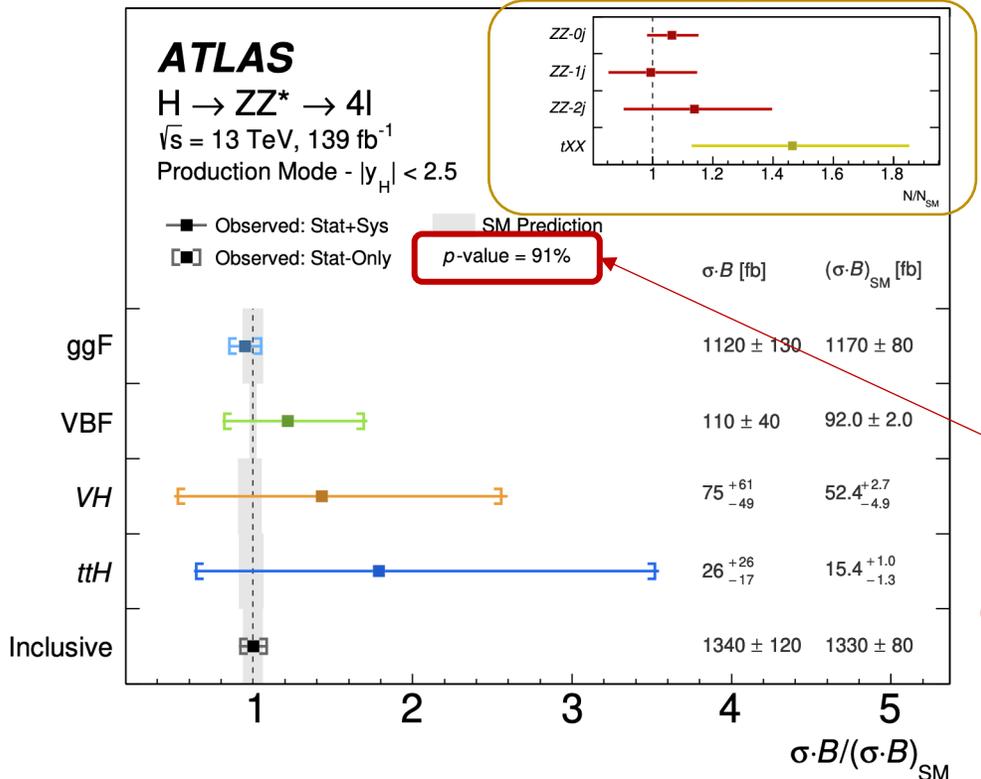


Production Cross Section measurements: STXS

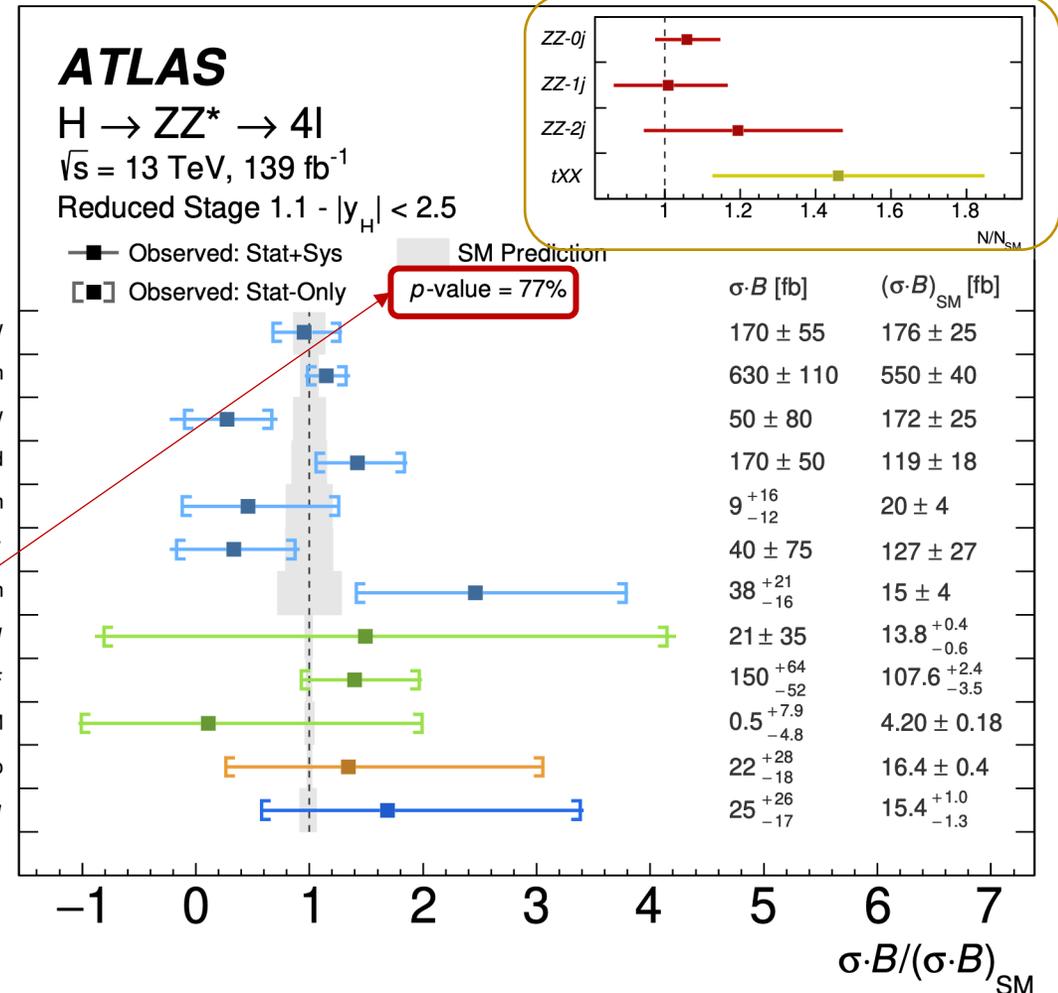
Results from Run 2

Sensitivity is statistically limited, but large Run 2 dataset allowed to probe more STXS regions than previous analyses

ZZ* and tXX background estimated from data in ad-hoc categories



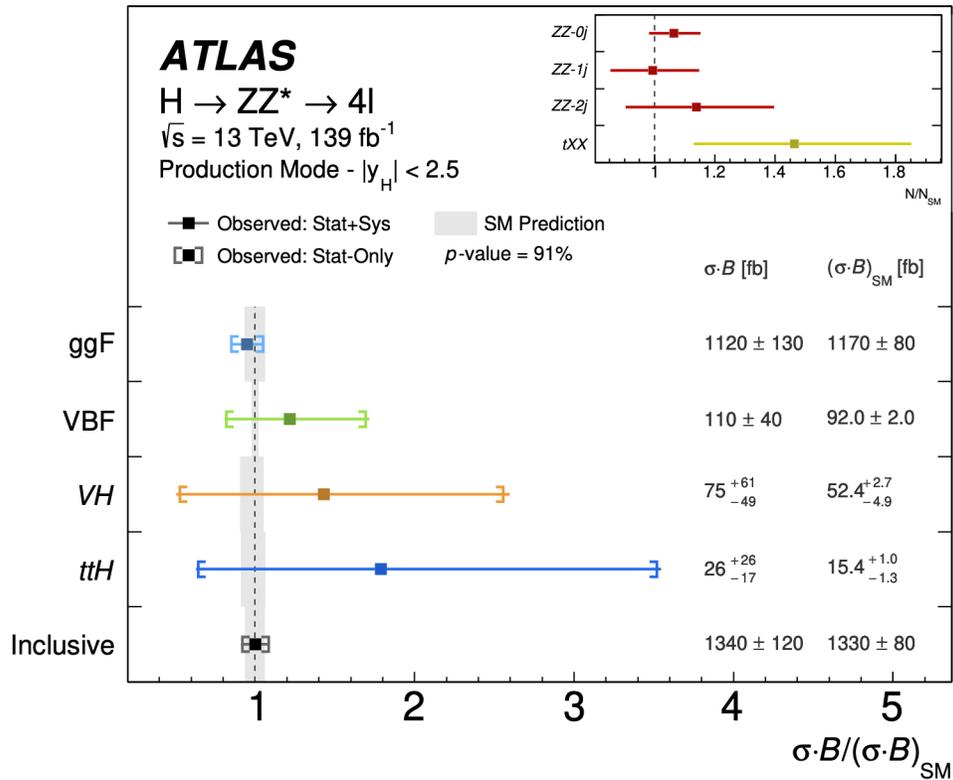
Excellent compatibility with Standard Model expectation



Production Cross Section measurements

STXS approach comes with several model dependencies

- Measurements are extrapolated to a phase space with the only selection that $|y_H| < 2.5$
- The Neural Network are trained on Standard Model predictions



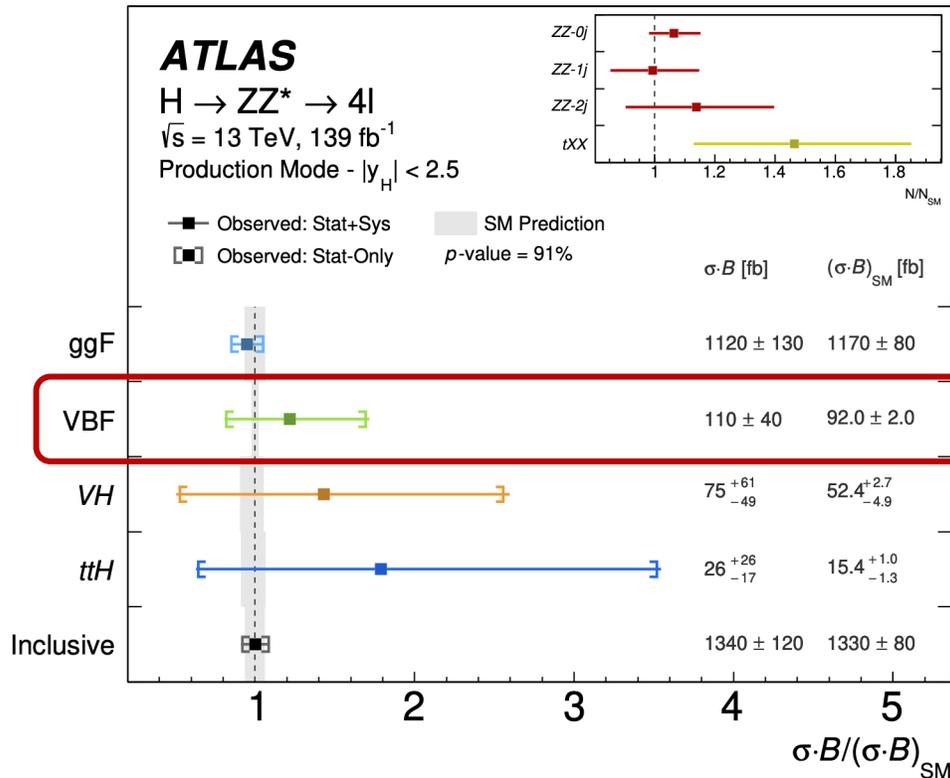
Can we be more model-independent?

Production Cross Section measurements

STXS approach comes with several model dependencies

- Measurements are extrapolated to a phase space with the only selection that $|y_H| < 2.5$
- The Neural Network are trained on Standard Model predictions

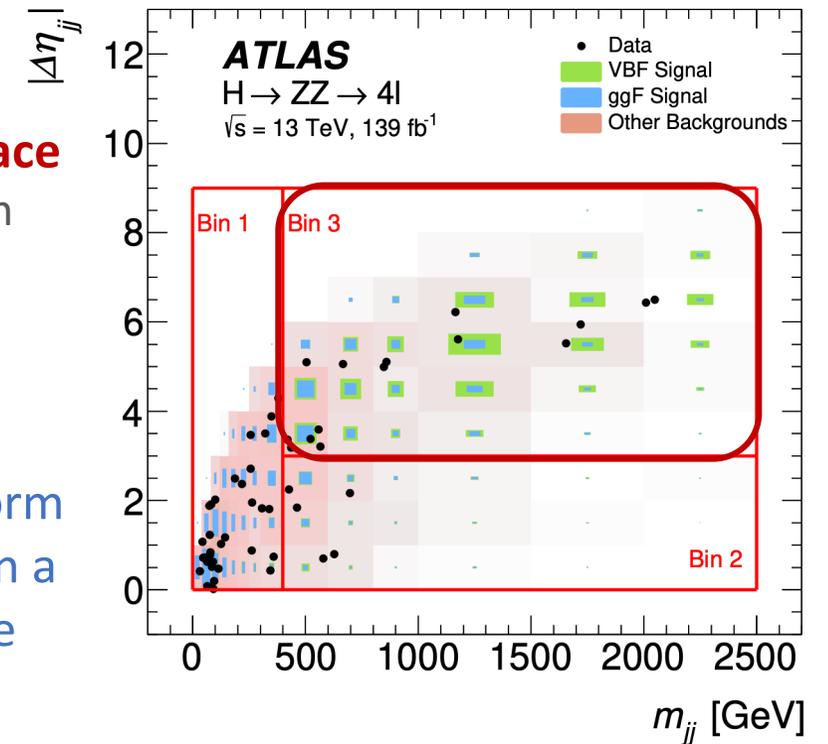
Fiducial Phase Space is the answer!



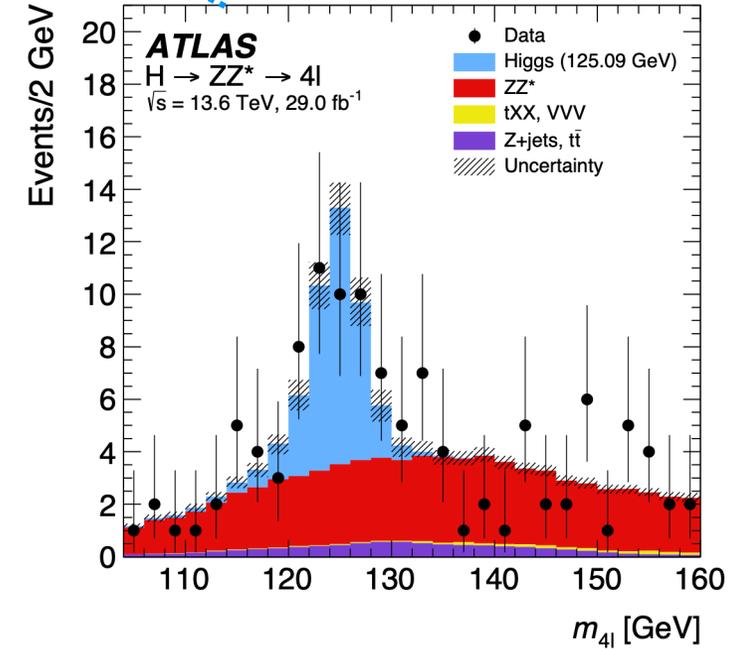
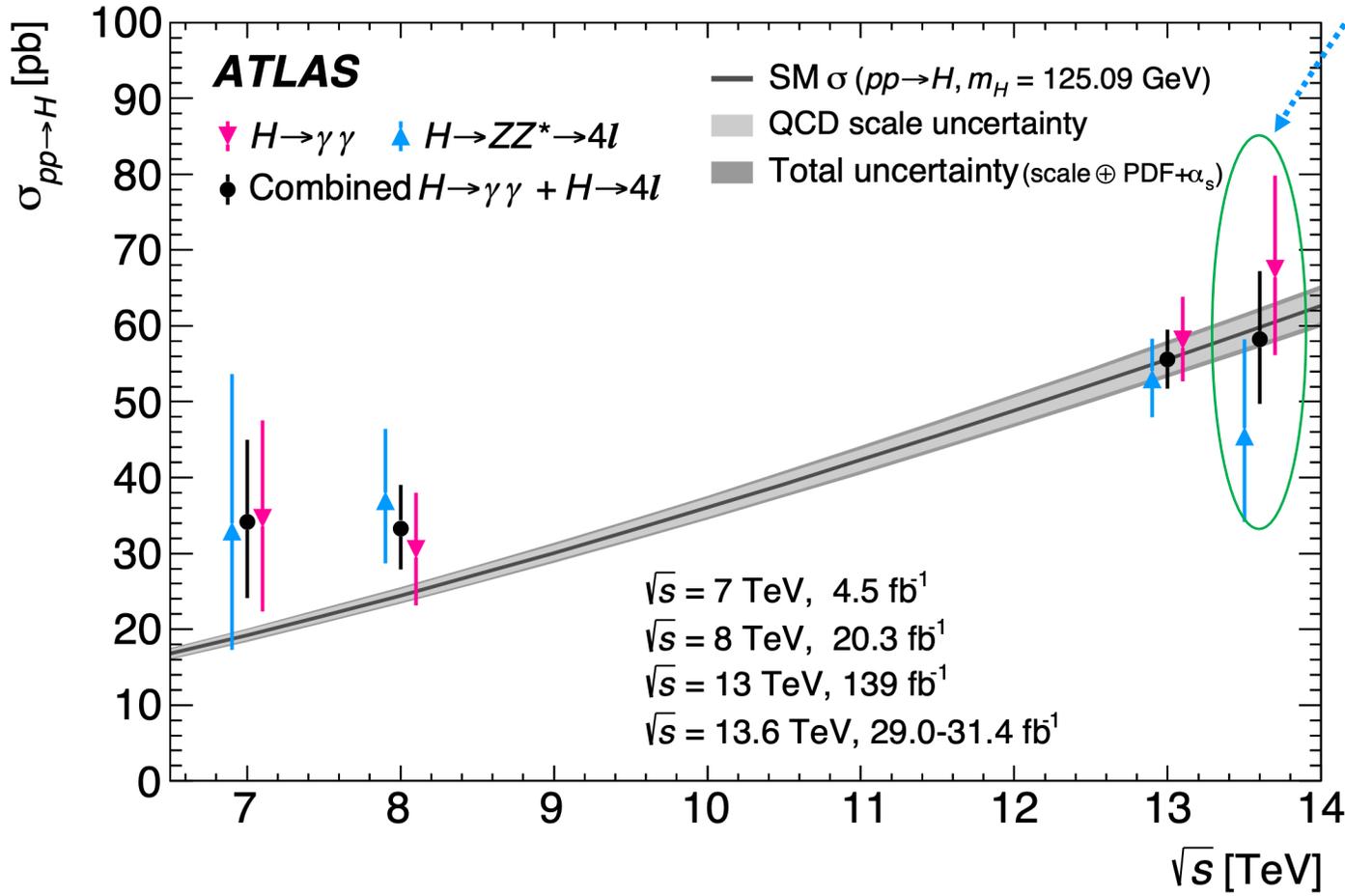
Apply “simple” kinematic selection to select a more **VBF-like fiducial phase space**

- Minimize the extrapolation
- Less assumptions

Open the possibility to perform differential measurements in a purer VBF-like phase space



Cross Section scale with collision energy



$$\sigma_{\text{fid}}^{4l} = 2.80 \pm 0.74 \text{ fb}$$

(SM: $3.67 \pm 0.19 \text{ fb}$)

$$\sigma_{\text{fid}}^{\gamma\gamma} = 76^{+14}_{-13} \text{ fb}$$

(SM: $67.6 \pm 3.7 \text{ fb}$)

$$\sigma_{\text{total}} = 58.2 \pm 8.7 \text{ pb}$$

(SM: $59.9 \pm 2.6 \text{ pb}$)

First measurement of the Higgs boson production cross section @ **13.6 TeV**

Cross Section measurements @ 13.6 TeV

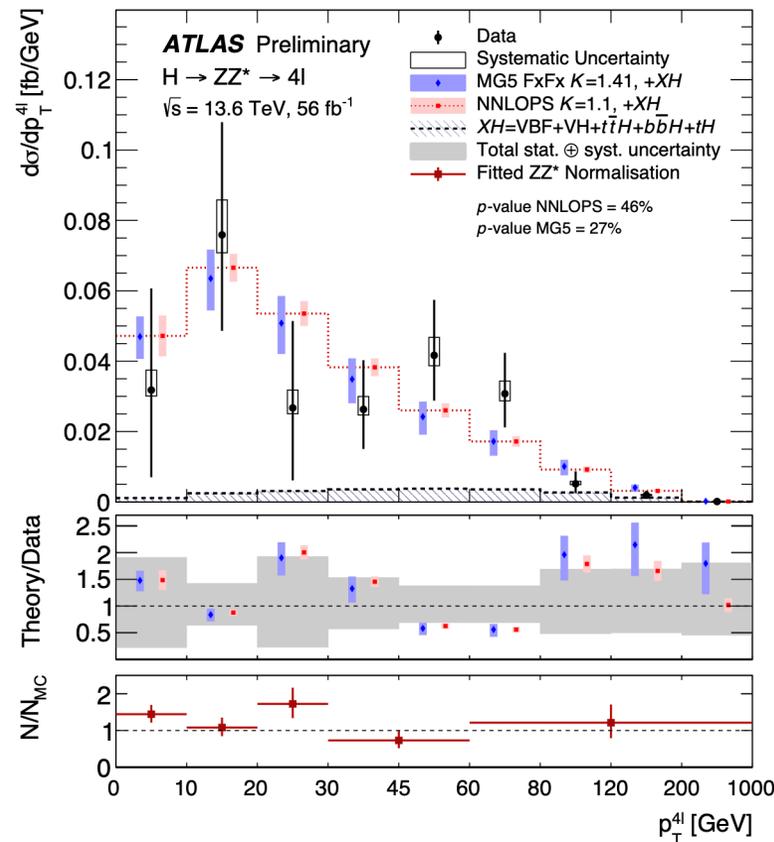
First differential measurements at 13.6 TeV, with partial Run 3 dataset of 56 fb^{-1} ($\sim 1/3$ of the full Run 2 dataset)

→ much more statistically limited

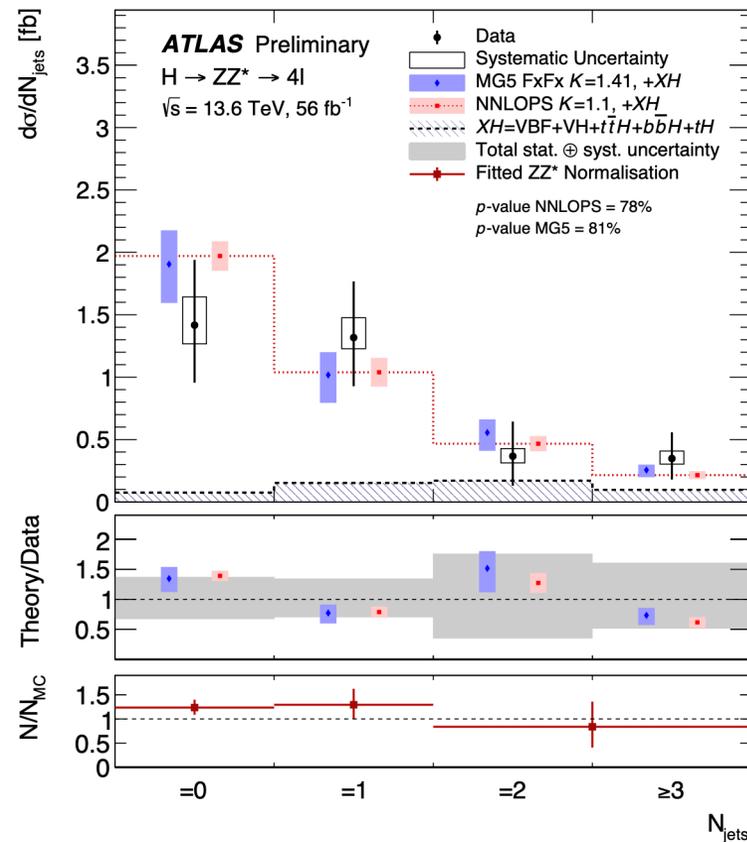
→ kinematic observables related to the Higgs production change their shape with the energy

(hard to appreciate 13 TeV \rightarrow 13.6 TeV scale due to the low statistics)

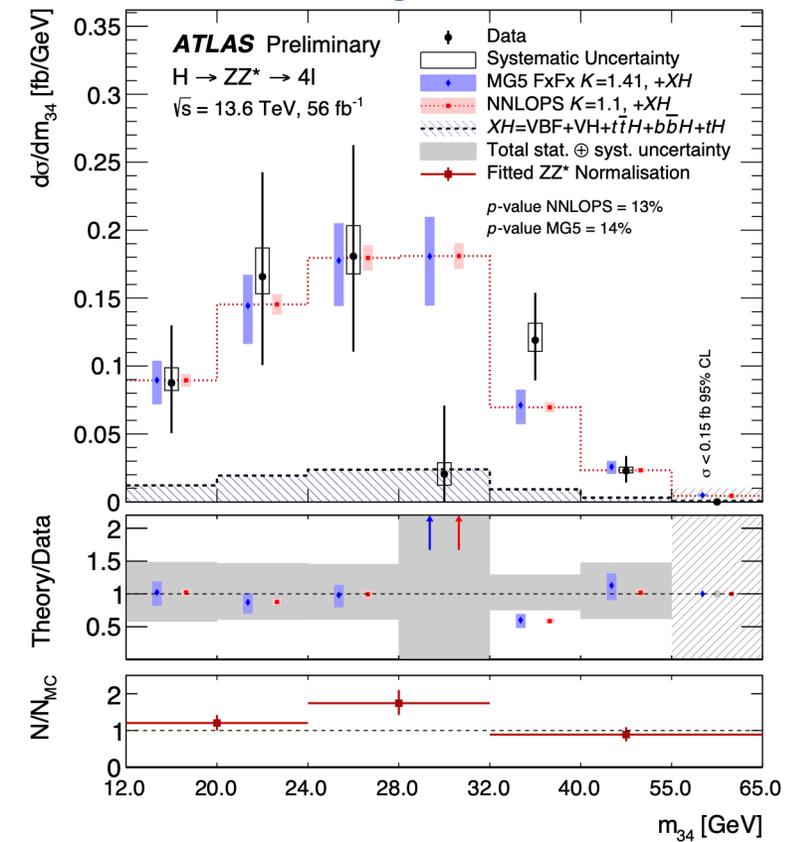
Higgs is produced at higher p_T



H + jets production changes



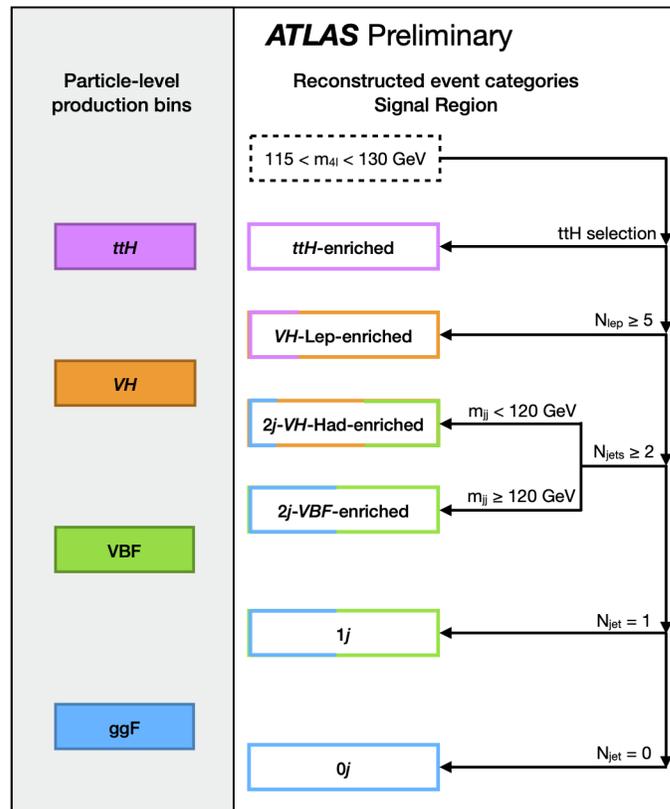
Decay kinematic observables do not change distribution



Cross Section measurements @ 13.6 TeV

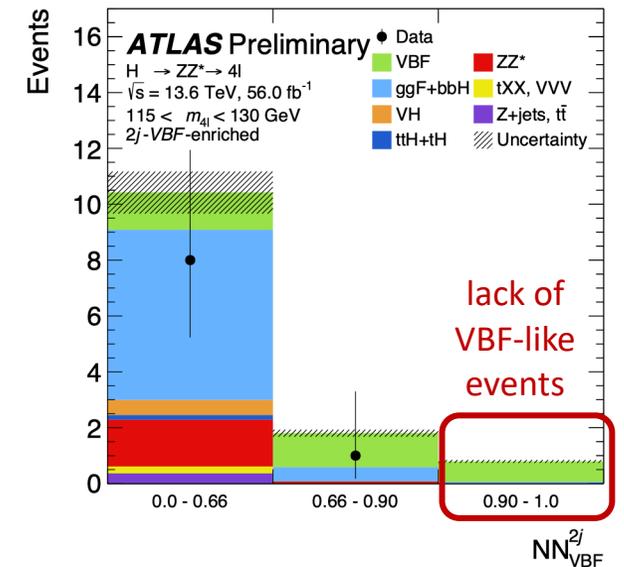
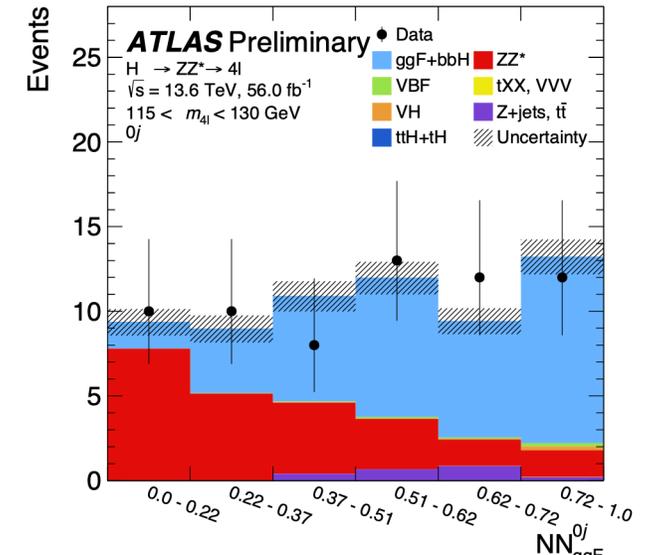
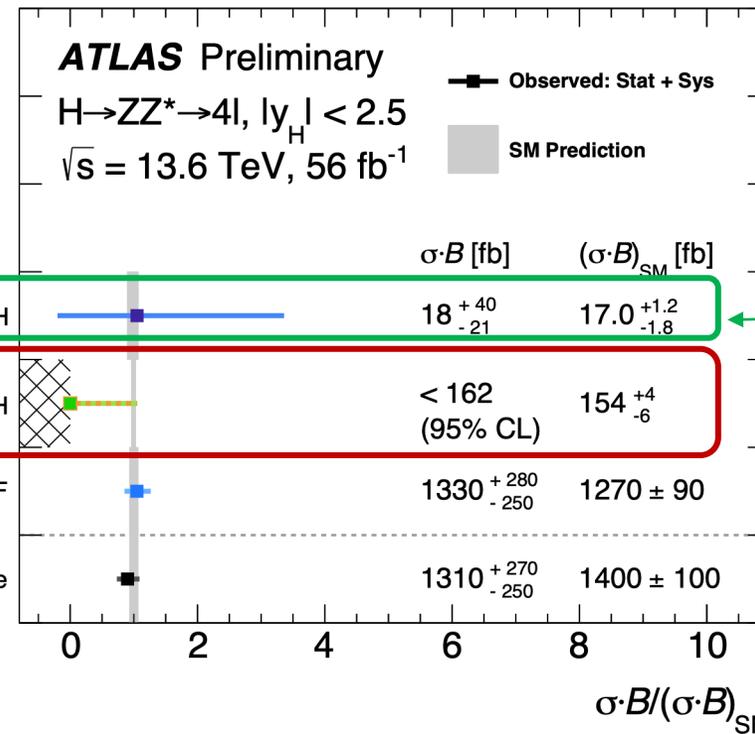
First production mode measurements at 13.6 TeV, with partial Run 3 dataset of 56 fb^{-1} ($\sim 1/3$ of the full Run 2 dataset)

- only Production Cross Section Stage (Stage 0) could be measured
- much more statistically limited



No data observed in the categories with highest sensitivity to VBF and VH

One ttH candidate!



lack of VBF-like events

...to Interpretations

(Anomalous) Higgs Couplings

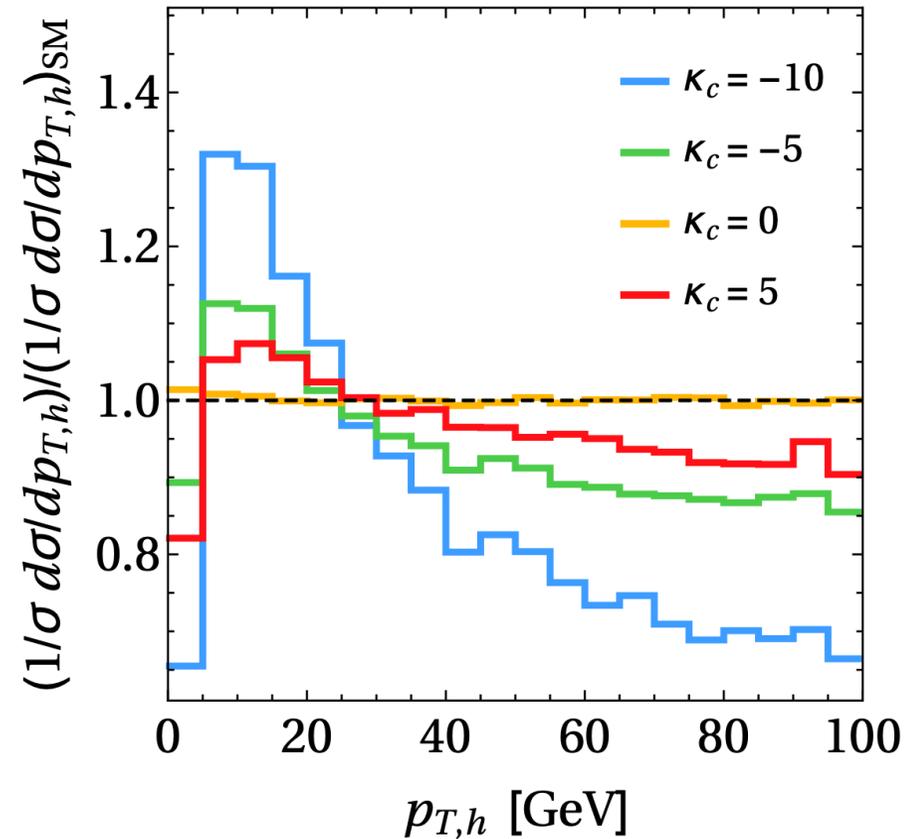
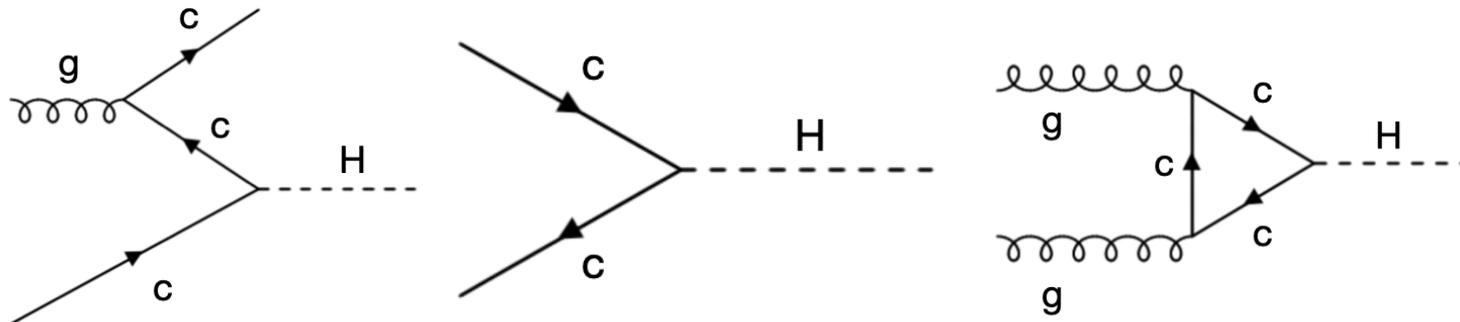
κ -framework: constraint on Yukawa couplings

Higgs coupling to light quarks (e.g. charm) may be possible to constrain without direct measurement (e.g. $H \rightarrow cc$ decay)

Higgs transverse momentum is sensitive to the Yukawa coupling of the charm, bottom quark

→ constrain the charm coupling from the bottom

- non-SM values of the coupling modifiers κ_c and κ_b are investigated



The variation of the coupling modifiers can impact both the Cross Sections, the Branching Ratio and the Shape → different tests can be performed with different level of model dependency

κ -framework: constraint on Yukawa couplings

Only the p_T^H shape can be modified

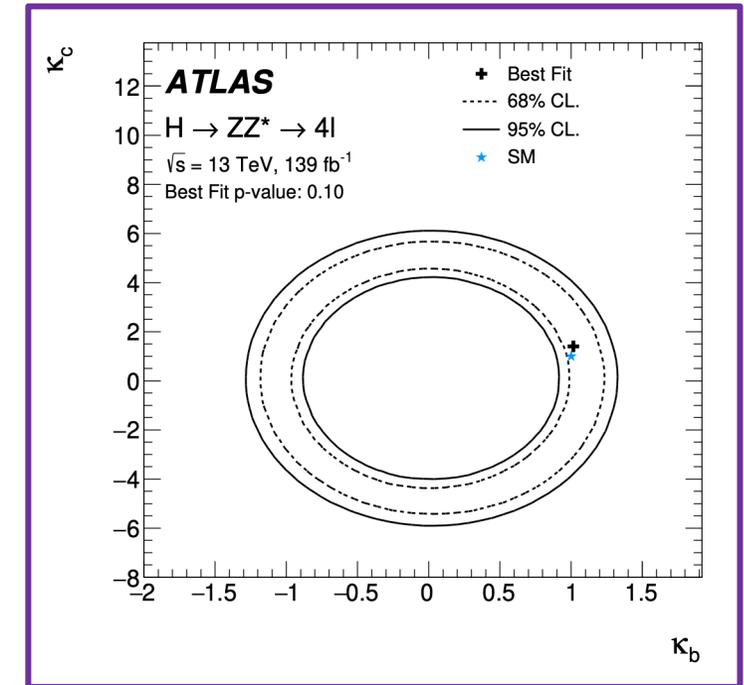
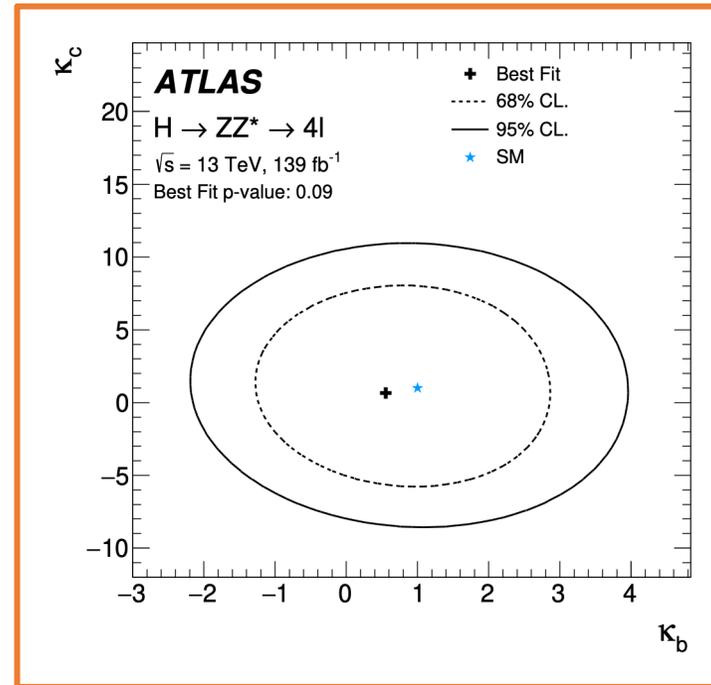
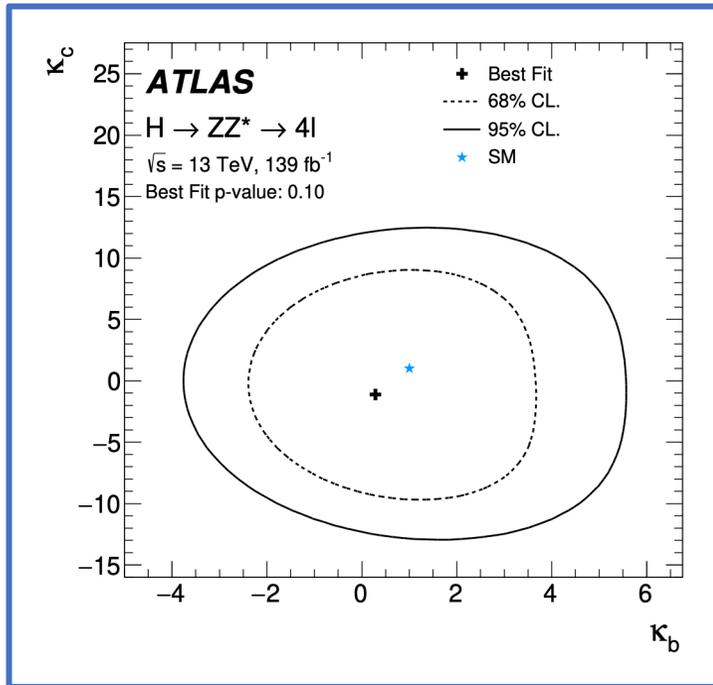
- Cross sections and BR are “free to float”
 - less assumptions

The cross section and the p_T^H shape can be modified

- Intermediate approach with no assumptions on the BR

The cross section, the p_T^H shape and the Branching Ratio can be modified

- more assumptions
- more stringent constraint

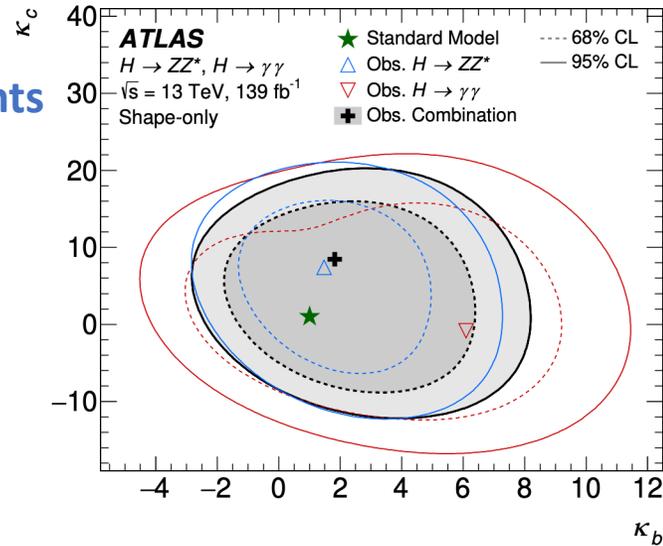
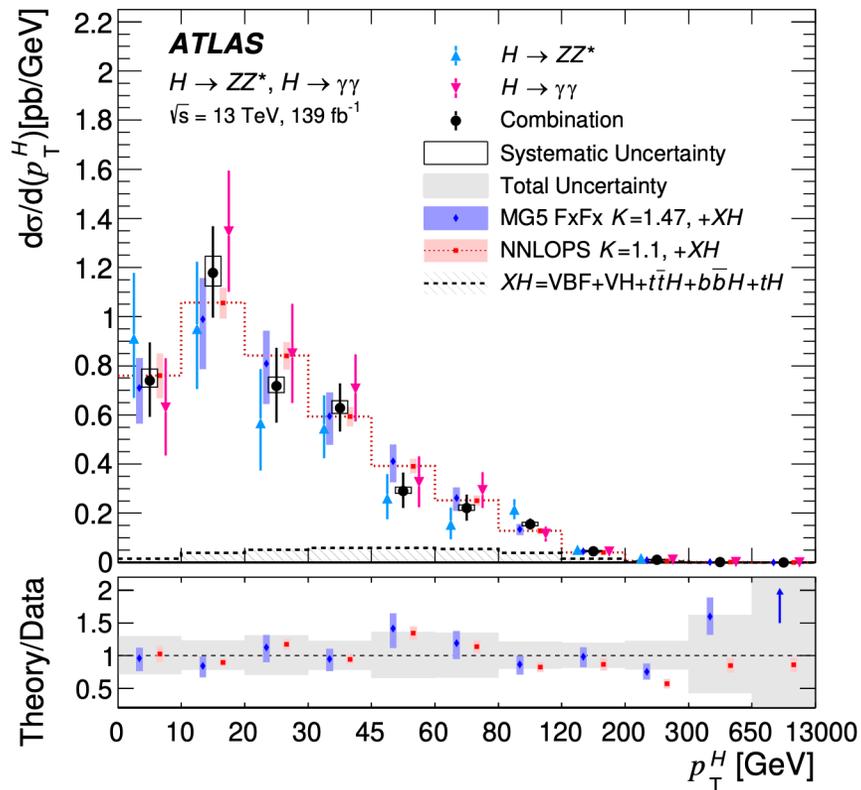


Shape dependency become subdominant
→ still sensitive to the sign!

Results from $H \rightarrow ZZ^* \rightarrow 4l$ analysis with full Run2

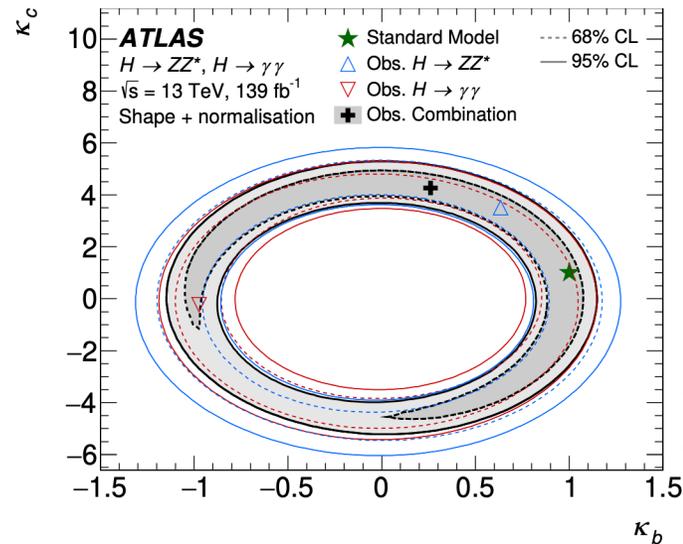
κ -framework: constraint on Yukawa couplings

This interpretation has been performed also on **combined differential cross section measurements of the Higgs p_T between $H \rightarrow ZZ^* \rightarrow 4l$ and $H \rightarrow \gamma\gamma$ decay channels**

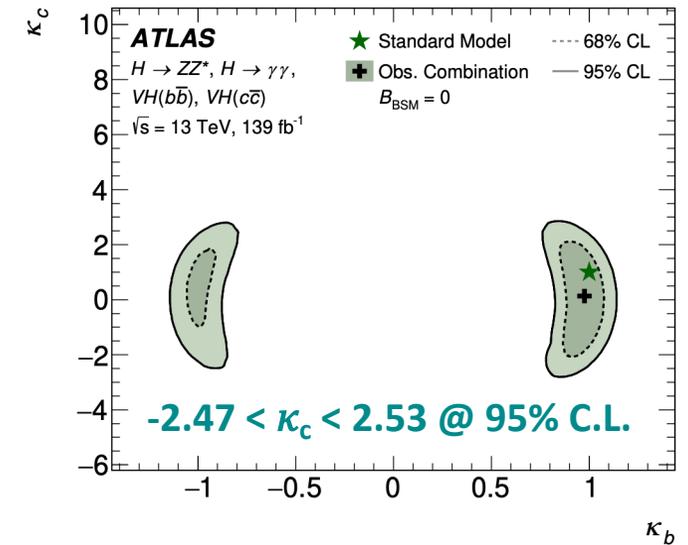


Anomalous couplings can only modify the p_T^H shape

Including direct measurements $VH(bb)$ and $VH(cc)$



Anomalous couplings can modify also the cross section and BR



κ -framework: vector vs fermion couplings

Production cross sections can be used to put constraint on the Higgs boson couplings modifiers

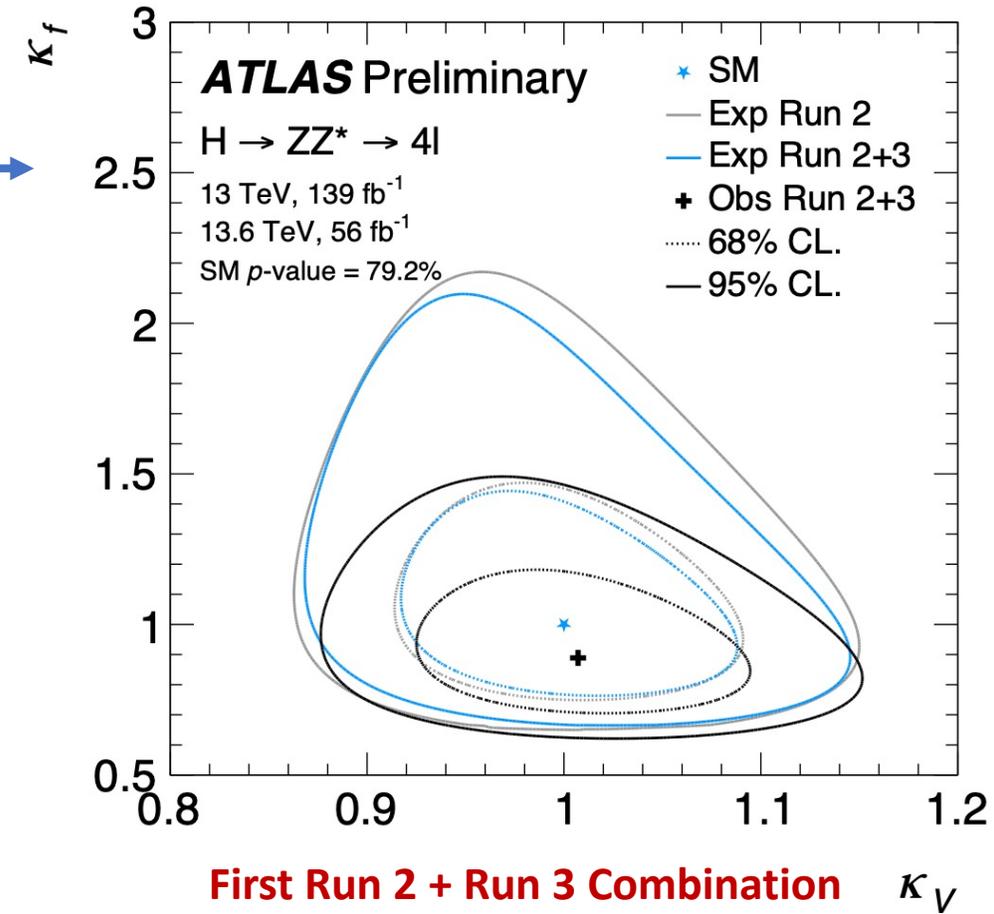
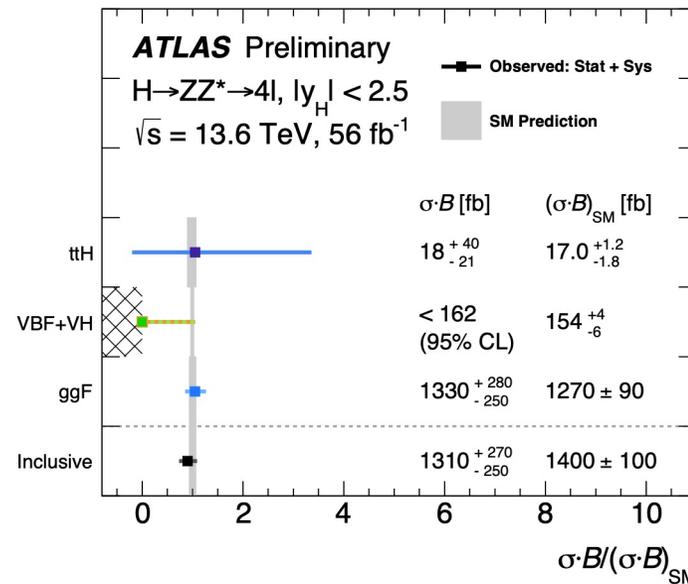
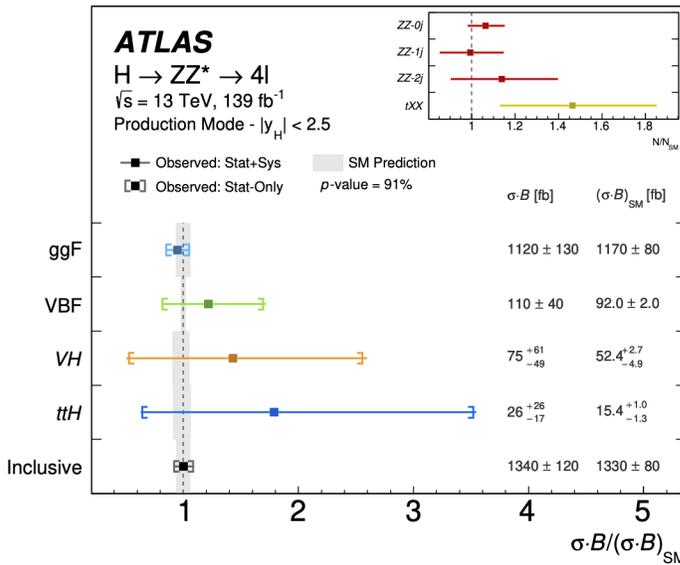
- Parametrizing the cross section as function of the κ we can extract constraints on the relative modifiers

H → ZZ* → 4l Run 2 results

Combined to increase the statistical power

H → ZZ* → 4l Run 3 results

Universal coupling strength modifiers κ_V (vector bosons) and κ_f (fermions)

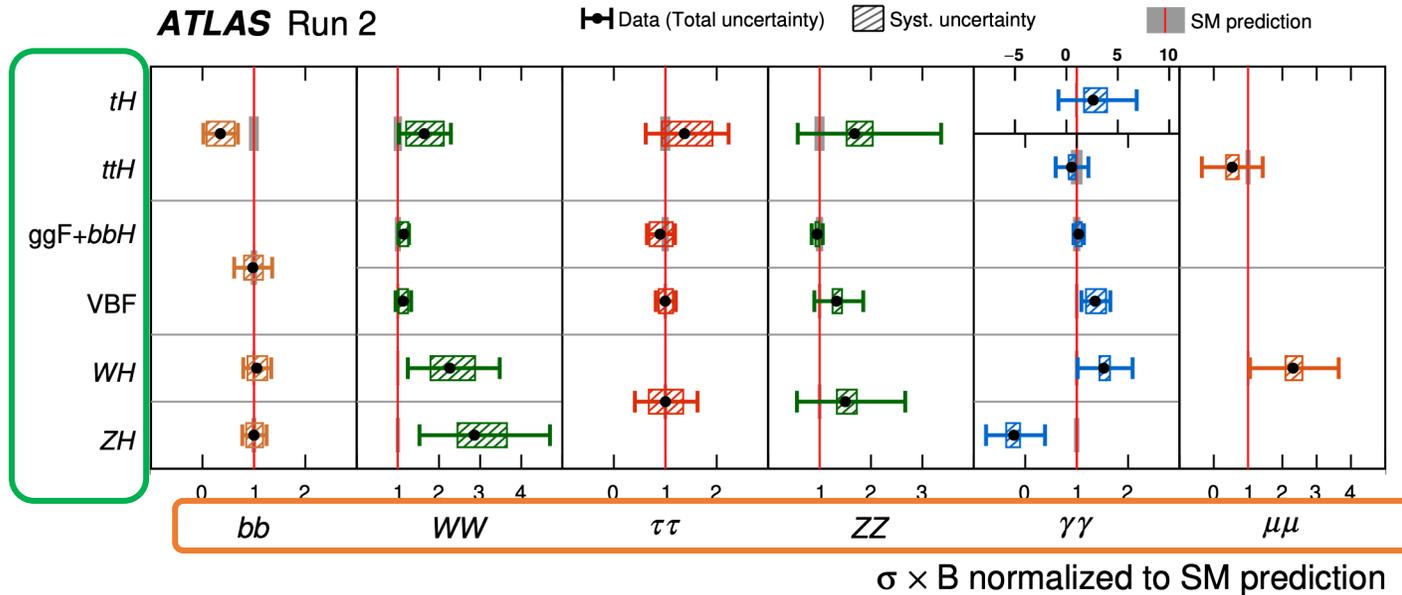


κ -framework: independent couplings

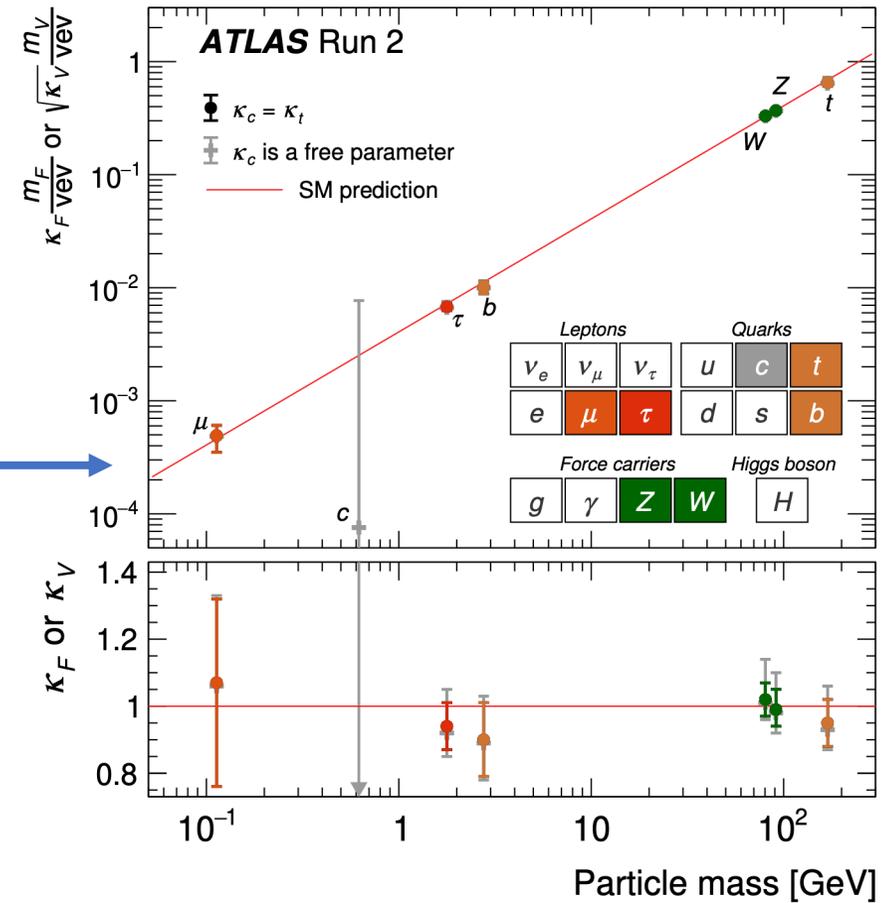
Production cross sections can be used to put constraint on the Higgs boson couplings modifiers

- Parametrizing the cross section as function of the κ we can extract constraints on the relative modifiers

Production Mode cross section measurement can be performed in different Decay Channels and then Combined



Generic parametrization with coupling strength modifiers for W, Z, t, b, c*, τ and μ treated independently



Pseudo Observables

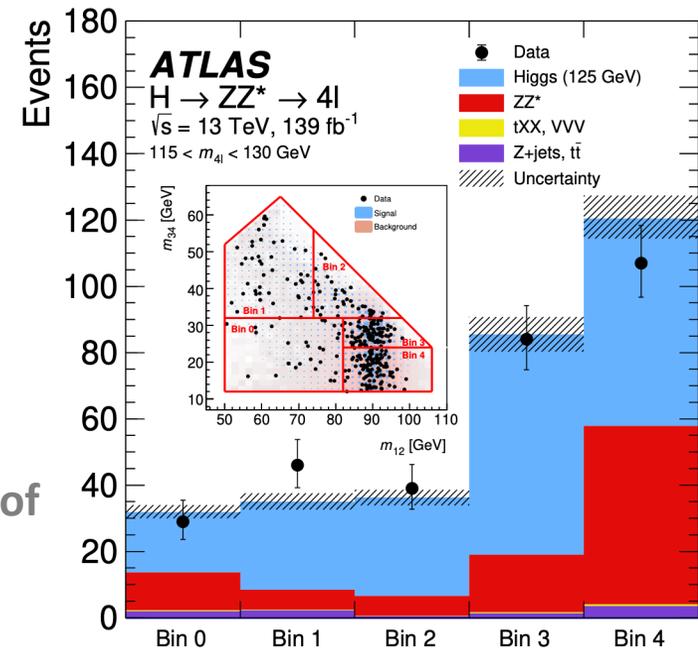
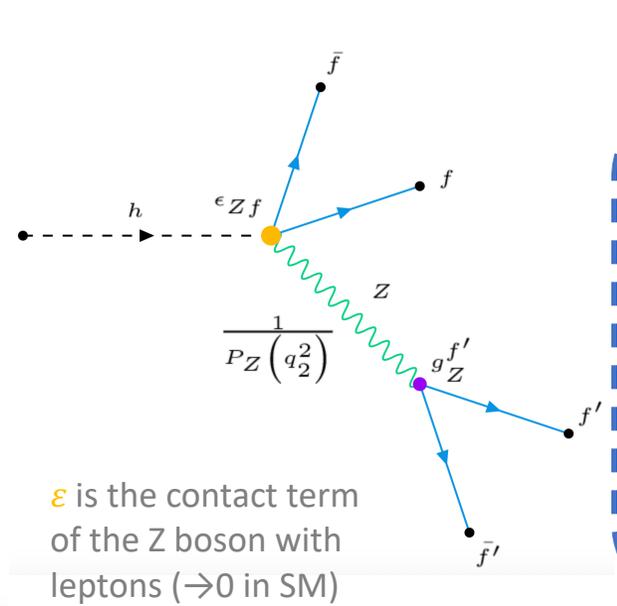
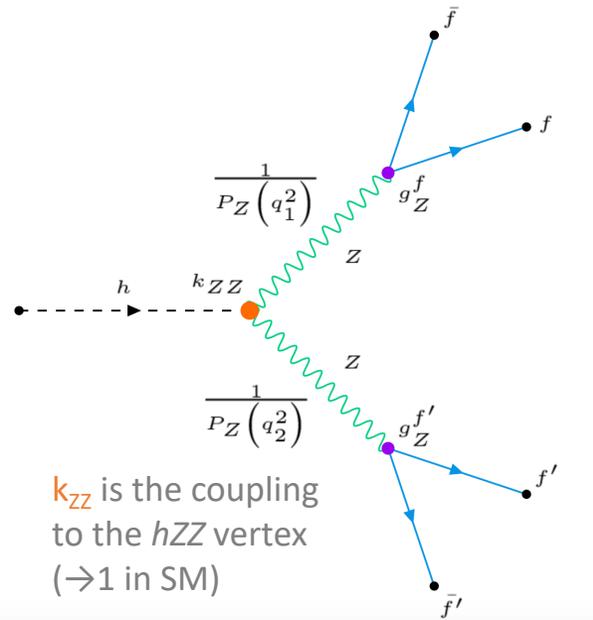
Study the **effective coupling** of the Higgs boson to the SM gauge bosons **using the invariant masses of the two Z m_{12} vs. m_{34} distribution**

- From **amplitude decomposition** the most interesting terms (assuming CP-invariance) are:

$$F_1^{ff'}(q_1^2, q_2^2) = \kappa_{ZZ} \frac{g_Z^f g_Z^{f'}}{P_Z(q_1^2) P_Z(q_2^2)} + \frac{\epsilon_{Zf}}{m_Z^2} \frac{g_Z^{f'}}{P_Z(q_2^2)} + \frac{\epsilon_{Zf'}}{m_Z^2} \frac{g_Z^f}{P_Z(q_1^2)}$$

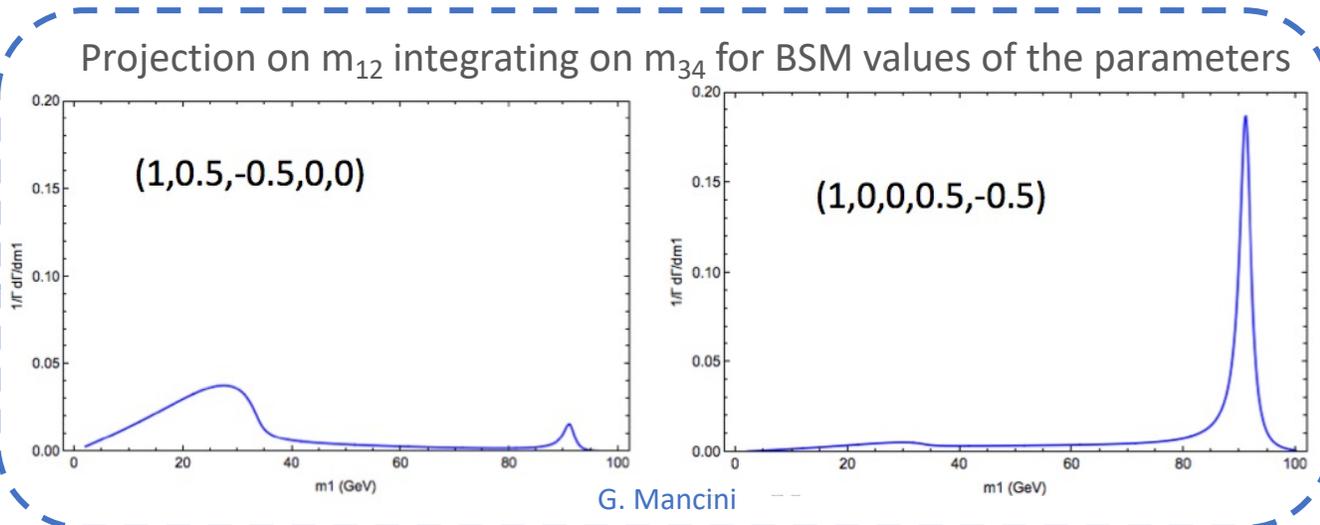
- $\kappa_{ZZ} \sim$ coupling modifier
- ϵ_{Zf} only possible source of **flavor non-universality!**

Momentum expansion around physical pole $P_Z(q^2) = q^2 - m_Z^2 + im_Z\Gamma_Z$



5 parameters can be studied: $\kappa_{ZZ}, \epsilon_{ZeL}, \epsilon_{Z\mu L}, \epsilon_{ZeR}, \epsilon_{Z\mu R}$

m_{12} vs. m_{34}



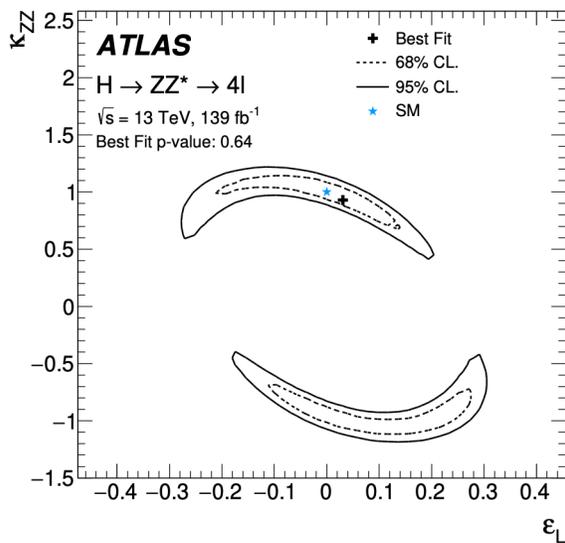
Pseudo Observables: Run 2 results

Different scenarios have been investigated with Run 2 analysis

Lepton Flavor Universal scenarios

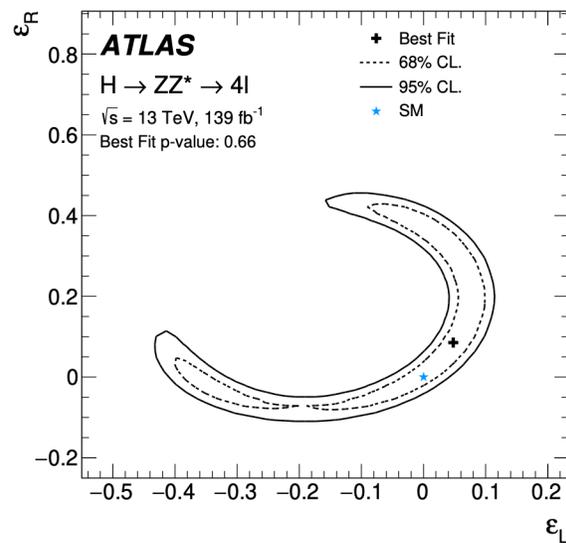
Linear EFT – inspired

Assume Higgs as part of a doublet No assumption on ϵ_R and ϵ_L
 $\epsilon_R \sim 0.48 \epsilon_L$



Contact Terms

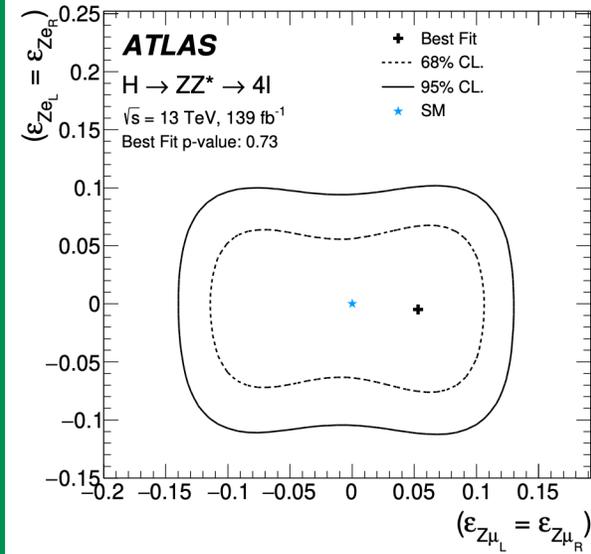
No assumption on ϵ_R and ϵ_L
 Assume $\epsilon_\mu = \epsilon_e$



Lepton Flavor Non – Universal scenarios

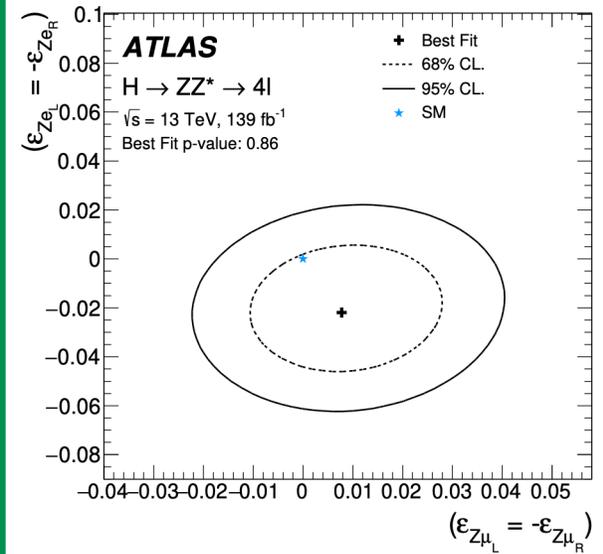
Vector Contact Terms

No assumption on ϵ_μ and ϵ_e
 Assume $\epsilon_L = \epsilon_R$



Axial Contact Terms

No assumption on ϵ_μ and ϵ_e
 Assume $\epsilon_L = -\epsilon_R$



Effective Field Theory

A powerful tool to interpret results in terms of New Physics without assumptions of the underlying theory (or just partially) → Perturbative expansion of the Standard Model Lagrangian adding new operators suppressed at a new high energy scale Λ

$$\sum_n \mathcal{L}_{\text{LE}}^{(n)} \rightarrow \mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \frac{c^{(5)}}{\Lambda} \mathcal{O}^{(5)} + \sum_i \frac{c^{(6)}}{\Lambda^2} \mathcal{O}^{(6)} + \dots$$

Wilson coefficients (free parameter of the theory) → constrained from the experimental measurements

Operators from SM fields with higher mass dimension

Effective Field Theory

A powerful tool to interpret results in terms of New Physics without assumptions of the underlying theory (or just partially) → Perturbative expansion of the Standard Model Lagrangian adding new operators suppressed at a new high energy scale Λ

$$\sum_n \mathcal{L}_{\text{LE}}^{(n)} \rightarrow \mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \frac{c^{(5)}}{\Lambda} \mathcal{O}^{(5)} + \sum_i \frac{c^{(6)}}{\Lambda^2} \mathcal{O}^{(6)} + \dots$$

Wilson coefficients (free parameter of the theory) → constrained from the experimental measurements

Dimension-5 (and Dim-7) operator violates the Lepton (and B-L) Number → neglected because suppressed by previous experiments

Operators from SM fields with higher mass dimension

Standard Model Effective Field Theory (SMEFT) Lagrangian

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)}$$

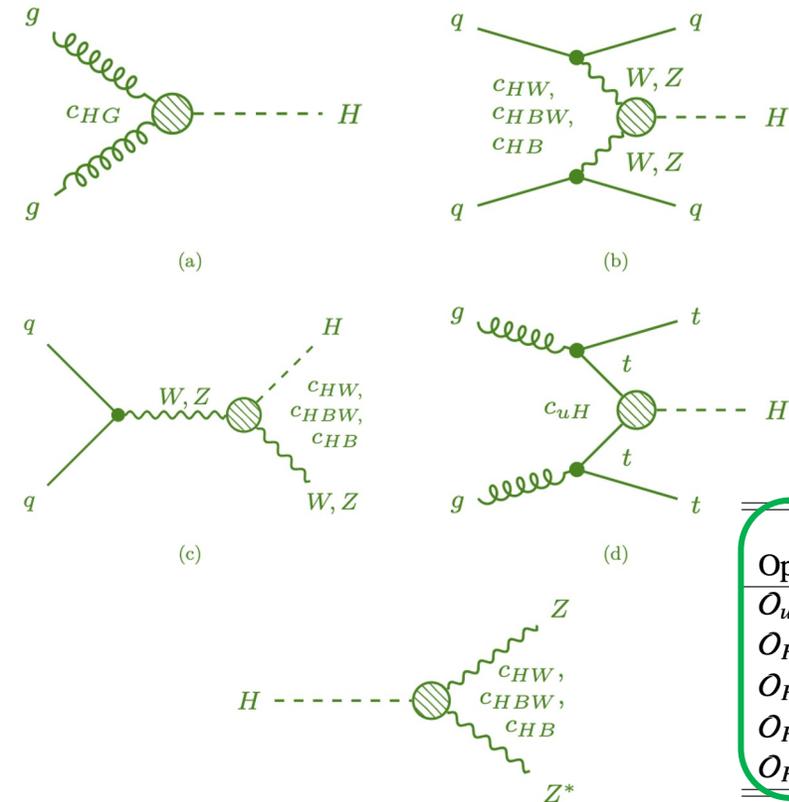
Leading contributions of New Physics are from dimension-six operators

Effective Field Theory

A powerful tool to interpret results in terms of New Physics without assumptions of the underlying theory (or just partially) → Perturbative expansion of the Standard Model Lagrangian adding new operators suppressed at a new high energy scale Λ

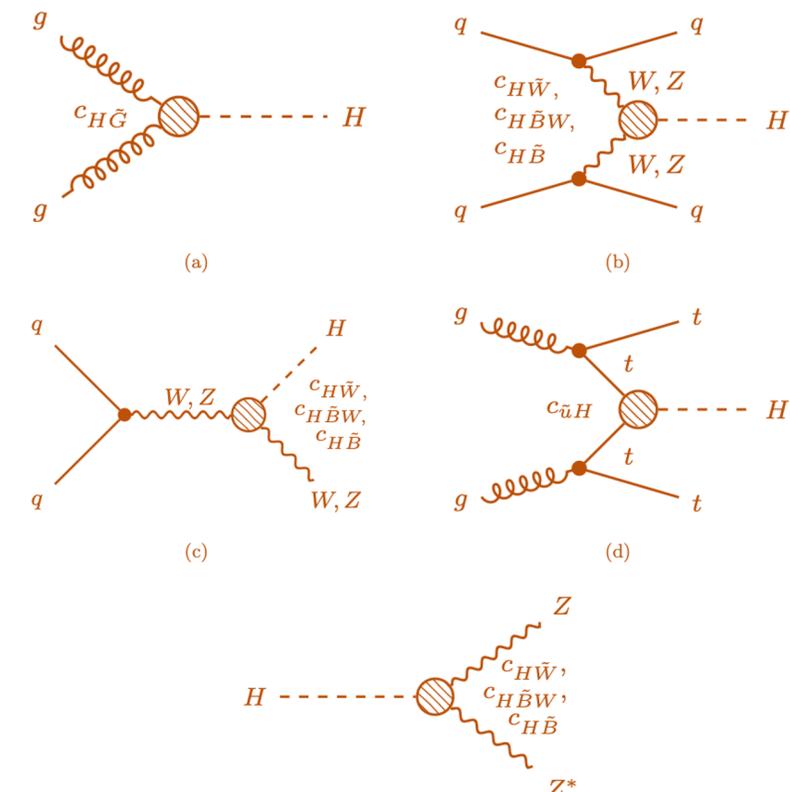
Standard Model Effective Field Theory (SMEFT) Lagrangian

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{c_i}{\Lambda^2} O_i^{(6)}$$



Production modes and $H \rightarrow ZZ^* \rightarrow 4l$ decay channel are sensitive just to a subset of all the Wilson coefficient

| CP-even | | | CP-odd | | |
|-----------|---|-----------|-------------------|---|-------------------|
| Operator | Structure | Coeff. | Operator | Structure | Coeff. |
| O_{uH} | $HH^\dagger \bar{q}_p u_r \tilde{H}$ | c_{uH} | O_{uH} | $HH^\dagger \bar{q}_p u_r \tilde{H}$ | $c_{\bar{u}H}$ |
| O_{HG} | $HH^\dagger G_{\mu\nu}^A G^{\mu\nu A}$ | c_{HG} | $O_{H\tilde{G}}$ | $HH^\dagger \tilde{G}_{\mu\nu}^A G^{\mu\nu A}$ | $c_{H\tilde{G}}$ |
| O_{HW} | $HH^\dagger W_{\mu\nu}^l W^{\mu\nu l}$ | c_{HW} | $O_{H\tilde{W}}$ | $HH^\dagger \tilde{W}_{\mu\nu}^l W^{\mu\nu l}$ | $c_{H\tilde{W}}$ |
| O_{HB} | $HH^\dagger B_{\mu\nu} B^{\mu\nu}$ | c_{HB} | $O_{H\tilde{B}}$ | $HH^\dagger \tilde{B}_{\mu\nu} B^{\mu\nu}$ | $c_{H\tilde{B}}$ |
| O_{HWB} | $HH^\dagger \tau^l W_{\mu\nu}^l B^{\mu\nu}$ | c_{HWB} | $O_{H\tilde{W}B}$ | $HH^\dagger \tau^l \tilde{W}_{\mu\nu}^l B^{\mu\nu}$ | $c_{H\tilde{W}B}$ |



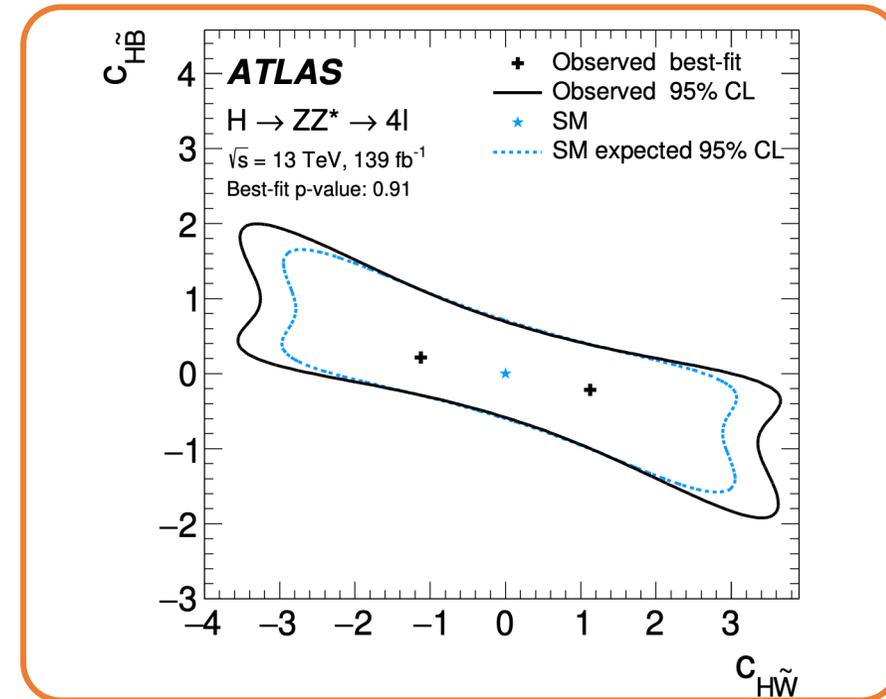
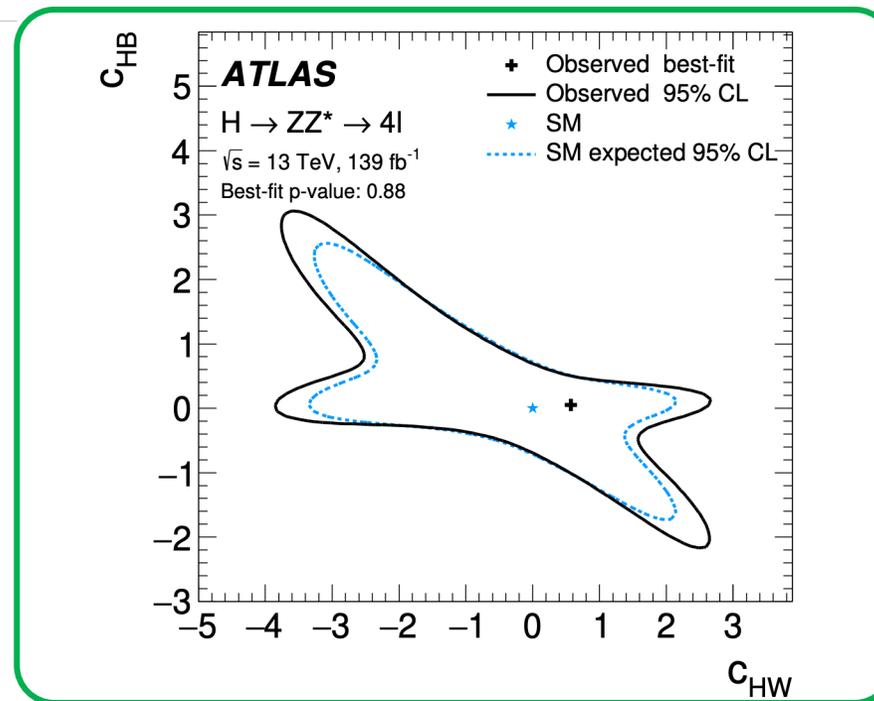
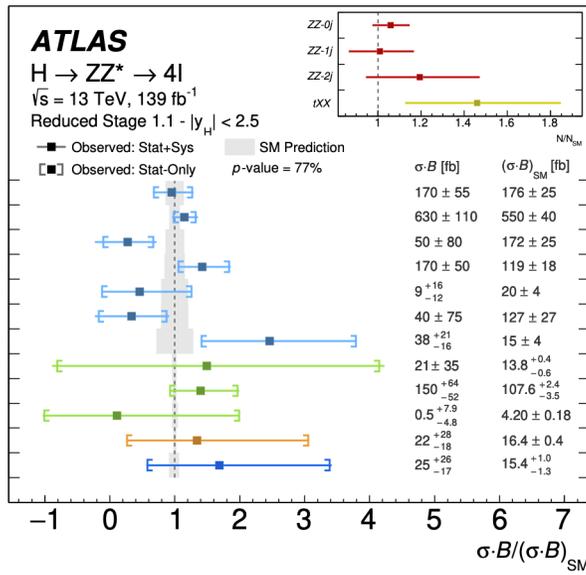
Effective Field Theory in Run 2

STXS measurements give enough sensitivity to probe anomalous Higgs boson couplings in EFT framework

Parametrize the STXS cross sections and Branching Ratios as function of the Wilson coefficients

Constraint CP-even couplings

Constraint CP-odd couplings
(not sensitive to the sign)



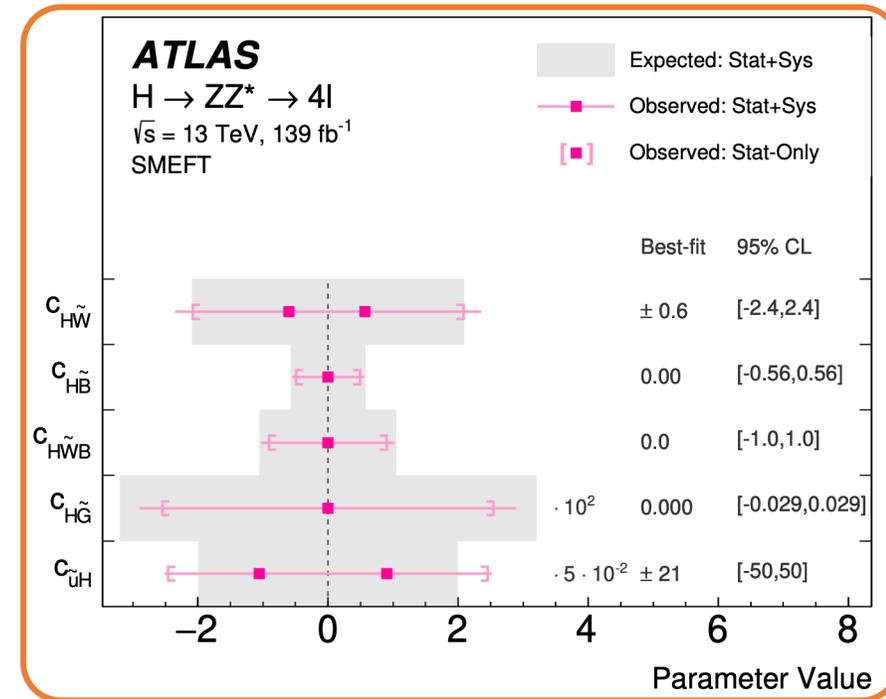
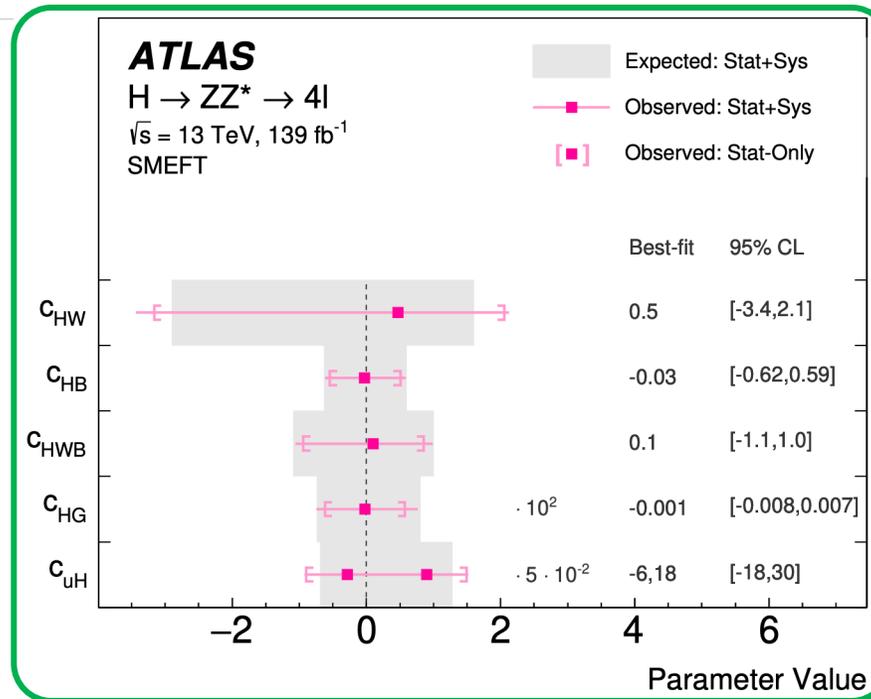
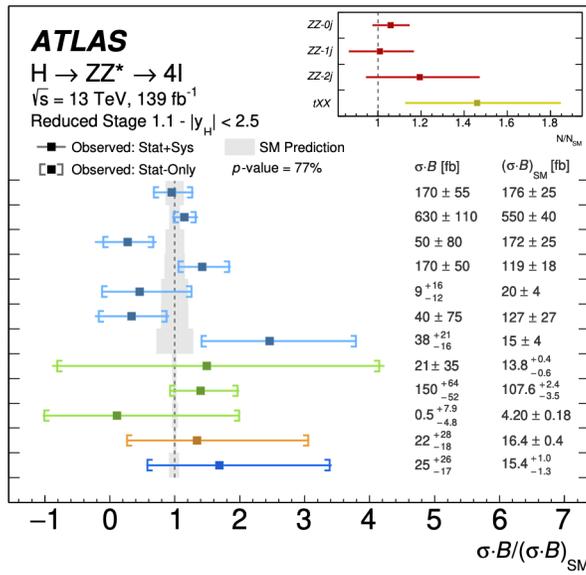
Effective Field Theory in Run 2

STXS measurements give enough sensitivity to probe anomalous Higgs boson couplings in EFT framework

Parametrize the STXS cross sections and Branching Ratios as function of the Wilson coefficients

Constraint CP-even couplings

Constraint CP-odd couplings
(not sensitive to the sign)

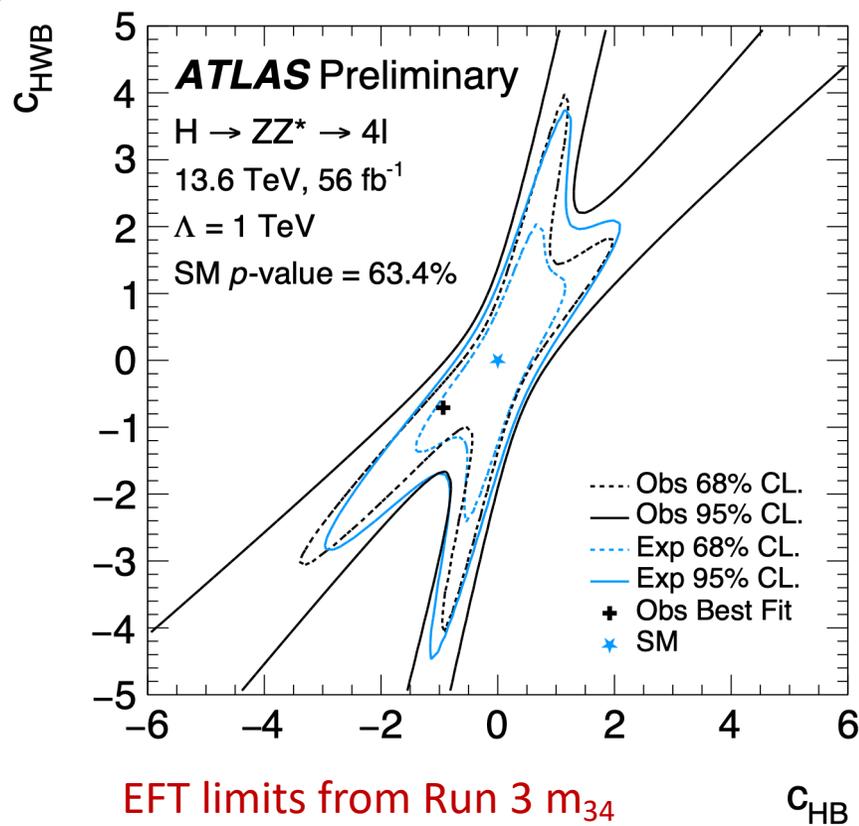
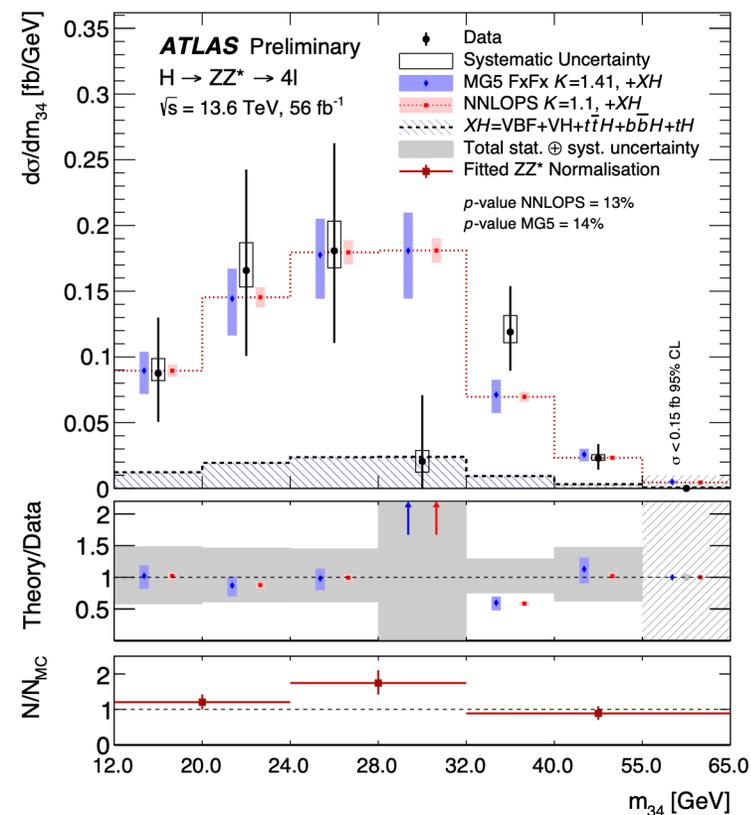


Effective Field Theory in Run 3

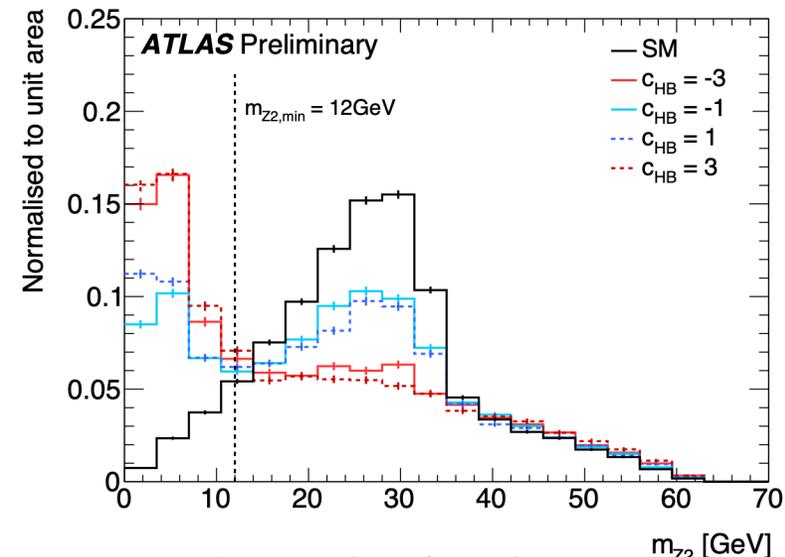
Off-shell Z distribution of m_{34} sensitive to BSM effects

Parametrize the fiducial cross sections as function of the Wilson coefficients

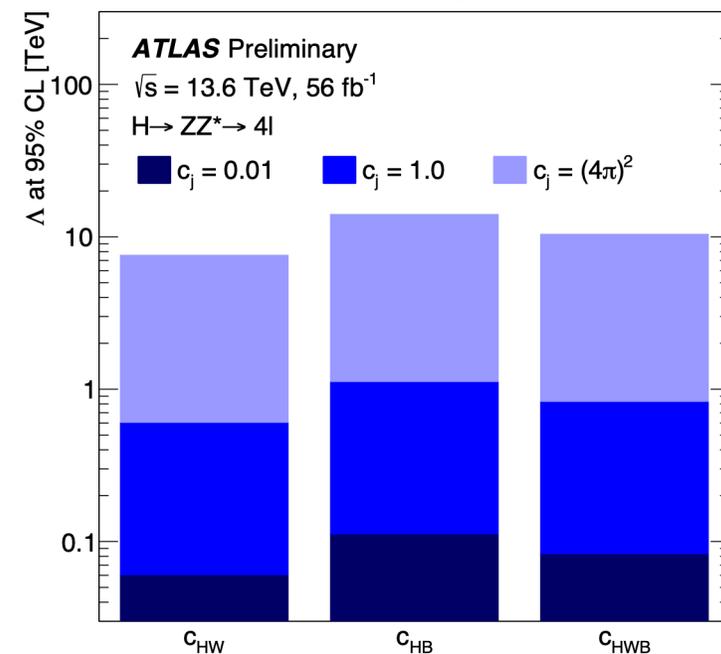
Constraint on CP-even couplings



EFT limits from Run 3 m_{34} interpretation comparable with one from Run 2 STXS



Limits on the Λ scale based on coupling values

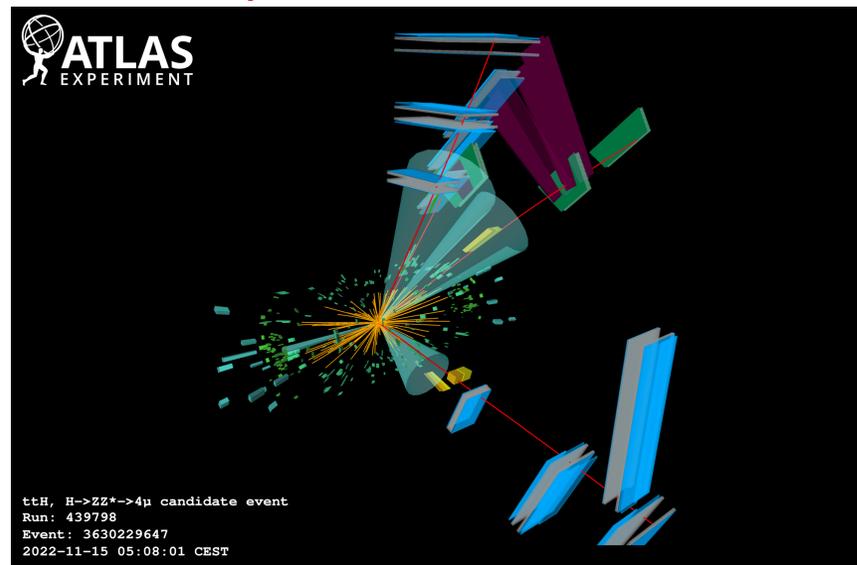


Conclusions

Cross section measurements to study the Higgs boson couplings with other Standard Model particles and to probe possible New Physics phenomena

- Presented results from full Run 2 analyses and preliminary picture using Run 3 dataset

Experimental Data



Observable Measurements



Lagrangian parameters

$$\begin{aligned}\mathcal{L} = & -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \\ & + i\bar{\psi} \not{D} \psi \\ & + \chi_i Y_{ij} \chi_j \phi + \text{h.c.} \\ & + |D_m \phi|^2 - V(\phi)\end{aligned}$$

More Run 3 data coming,
Stay Tuned!



Thanks for the attention

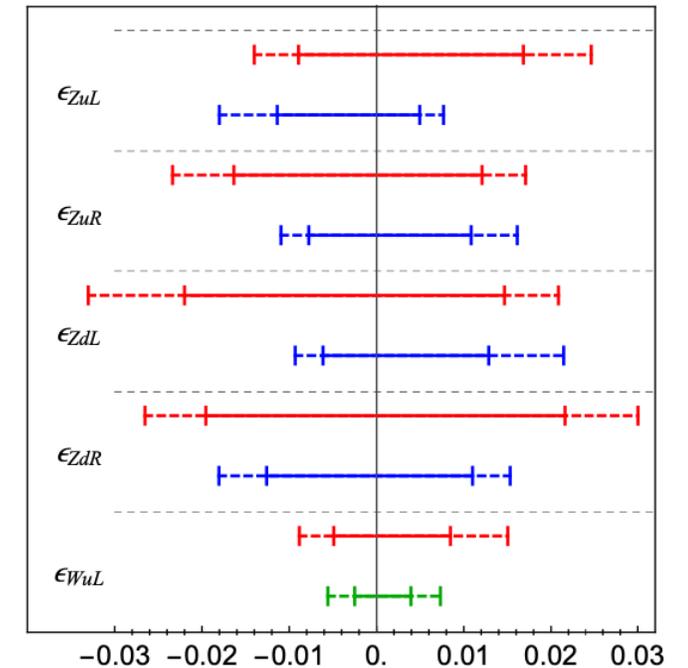
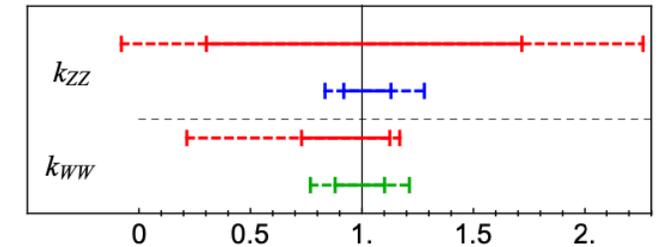
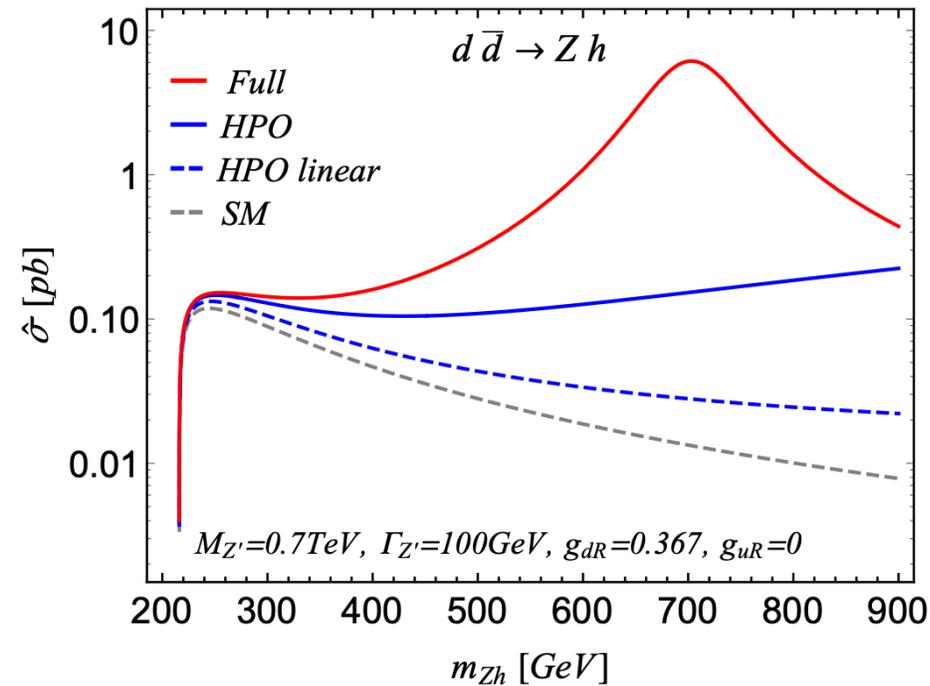
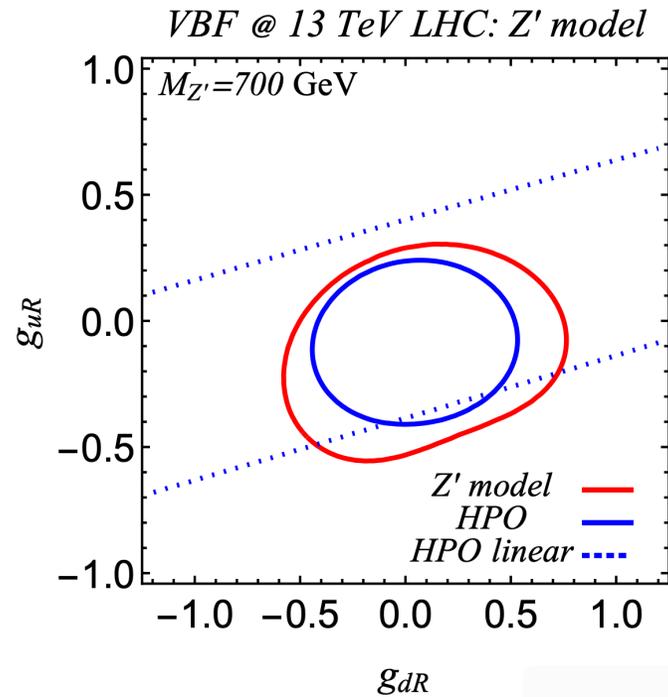
Pseudo Observables: Prospects

[arXiv:1512.06135](https://arxiv.org/abs/1512.06135)

Constrain PO using also EW production modes \rightarrow probe new PO $k_{ZZ}, k_{WW}, \epsilon_{ZuL}, \epsilon_{ZdL}, \epsilon_{ZuR}, \epsilon_{ZdR}, \epsilon_{WuL}$

Prospects for Higgs PO in EW production @ the HL-LHC

- **VBF**: double-differential distribution on p_T^{j1} vs p_T^{j2} to access the $F(q_1^2, q_2^2)$
- **ZH** (or **WH**): differential distribution of p_T^Z or $m_{Zh} \sim q^2$



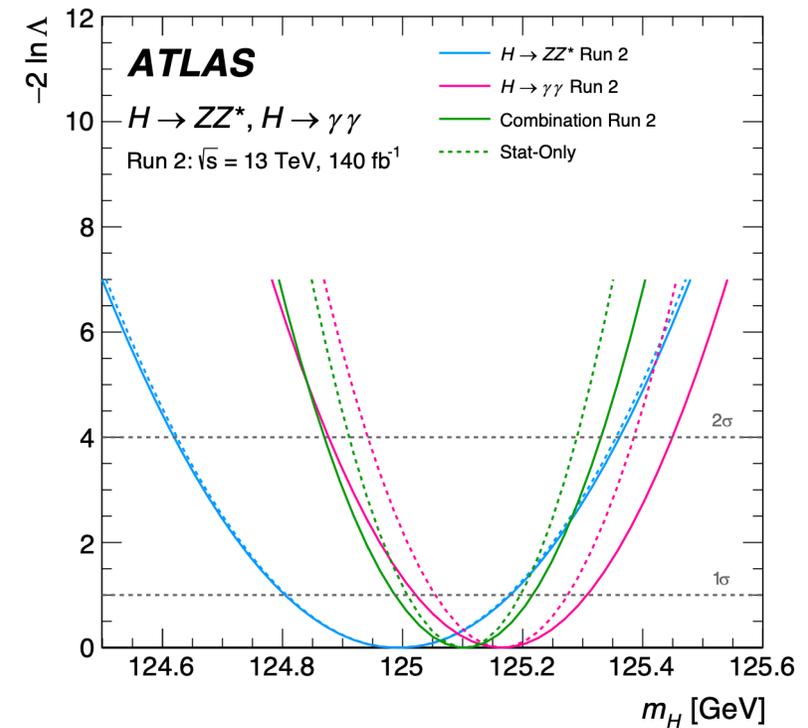
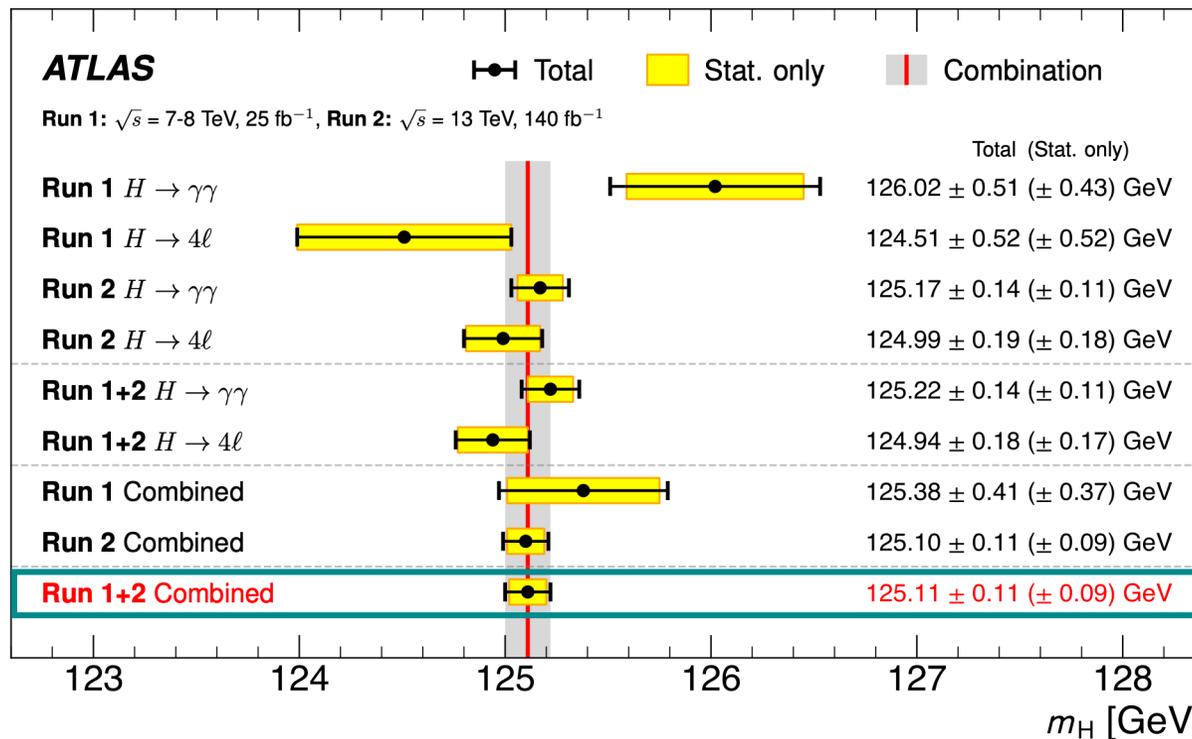
\rightarrow Define fiducial volume targeting specific production modes can improve sensitivity

The Higgs boson Mass

$H \rightarrow ZZ^* \rightarrow 4\ell$ and $H \rightarrow \gamma\gamma$ are the most sensitive channels

- Clear signature final states
- High mass resolution 1-2 %
- Main uncertainties: **Electron/Photon** energy scale and **Muon** momentum scale
- Combination of the two channels and of the two runs lead to the **most precise measurement of the Higgs boson mass!**

[Phys. Rev. Lett. 131 \(2023\) 251802](#)



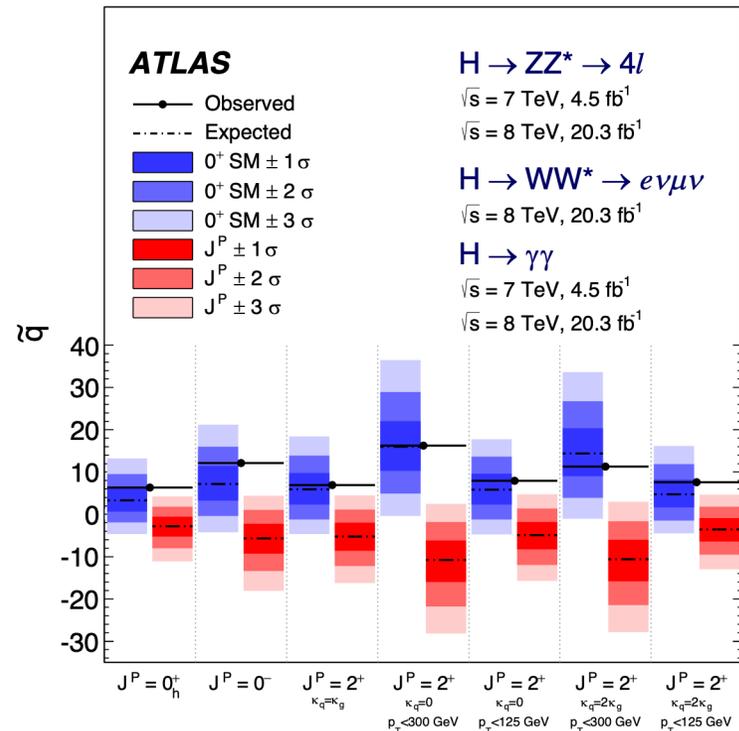
$$m_H = 125.11 \pm 0.11 \text{ GeV}$$

0.09% precision achieved on this fundamental parameter of the Standard Model of particle physics.

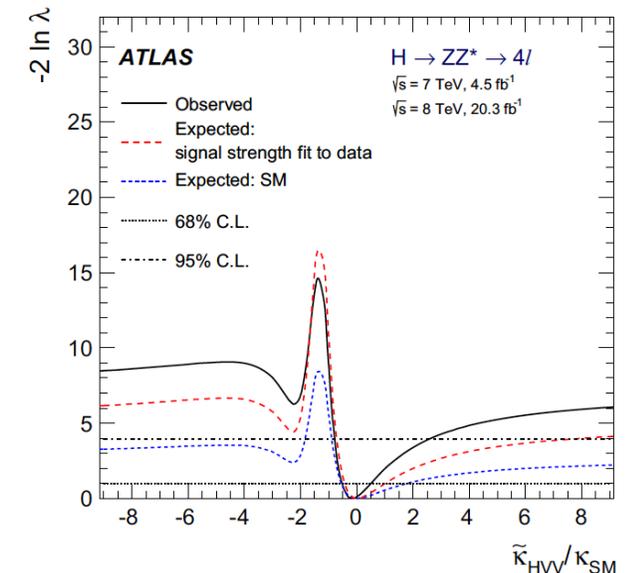
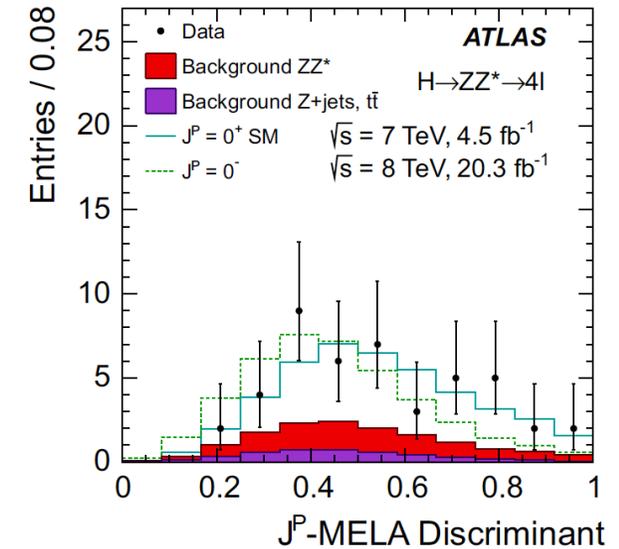
The Higgs boson CP structure

In Run 1 analyses aimed to assert that the Higgs boson is CP-even

[Eur. Phys. J. C 75 \(2015\) 476](#)



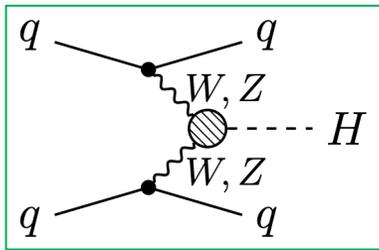
- Test of fixed spin and parity hypotheses
 - Using kinematic 4-lepton information
 - J^P MELA or BDT to discriminate different hypotheses
 - $J^P = 0^+$ compared with alternative spin models \rightarrow non-SM hypothesis excluded with at least 99.9% CL in favor of SM Higgs boson with Spin/Parity 0^{++}
- Investigation of mixing CP-even and CP-odd state looking at HVV tensor structure
 - First use of the Optimal Observables...
 - EFT model based on Higgs Characterization
- In Run 2 the focus is mainly on production modes rather than decay only.



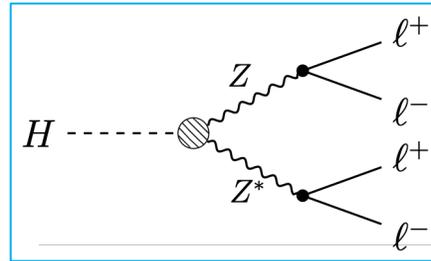
The Higgs boson CP structure

Looking for signs of CP-violation in the Higgs sector [arXiv:2304.09612](https://arxiv.org/abs/2304.09612)

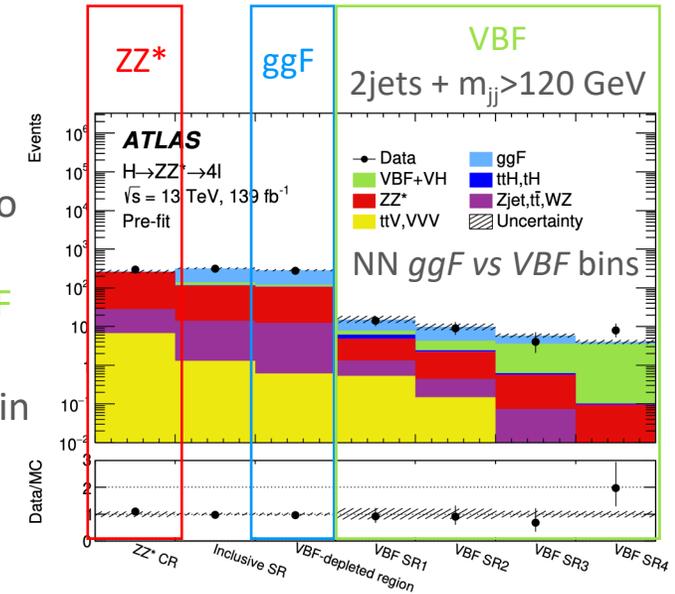
- Study the coupling with vector bosons (HVV) both at production level with **Vector Boson Fusion (VBF)** production and at decay level in the $H \rightarrow ZZ^* \rightarrow 4l$ decay



VBF \rightarrow high Q^2 process BSM effects expected to be higher

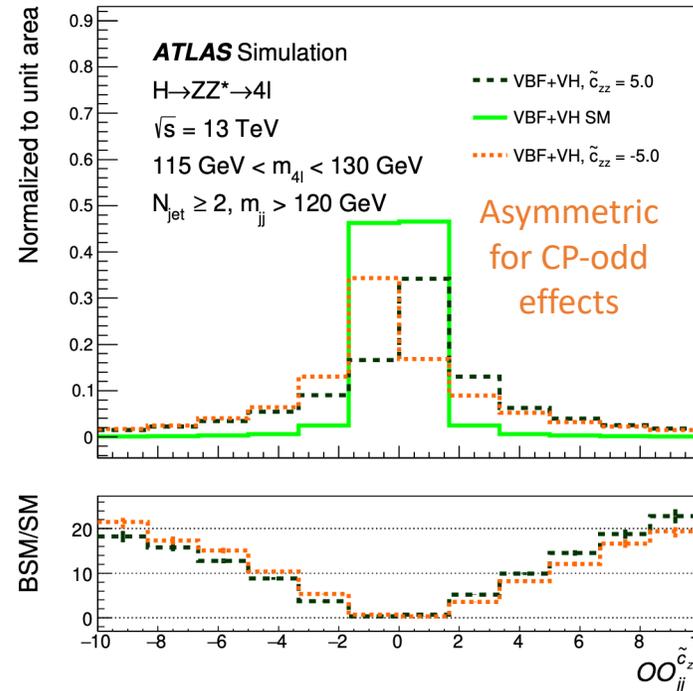


- Categorization** to maximize the sensitivity to VBF production and estimate the main backgrounds



- Use of **observables** optimized to discriminate different CP hypothesis
 - Rate cannot disentangle anomalous CP-even or CP-odd effects, **observable shapes** does
- \rightarrow Matrix – Element based variable called **Optimal Observable (OO)**

$$OO_1(c) = \frac{|\mathcal{M}_{\text{Mix}}(c)|^2 - |\mathcal{M}_{\text{SM}}|^2 - |\mathcal{M}_{\text{BSM}}(c)|^2}{|\mathcal{M}_{\text{SM}}|^2}$$



Constraint on EFT CP-odd couplings

