

A Light Dark Neutrino Sector

Matheus Hostert
(IPPP, Durham University)

based on:

1903.07589 / 1903.07590

Peter Ballett and Silvia Pascoli

1812.08768

Carlos Argüelles and Yu-Dai Tsai

Outline

Dark neutrino models — what and why?

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Neutrino masses from our dark sector

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Phenomenology — MiniBooNE anomaly

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Dark neutrino models — what and why?

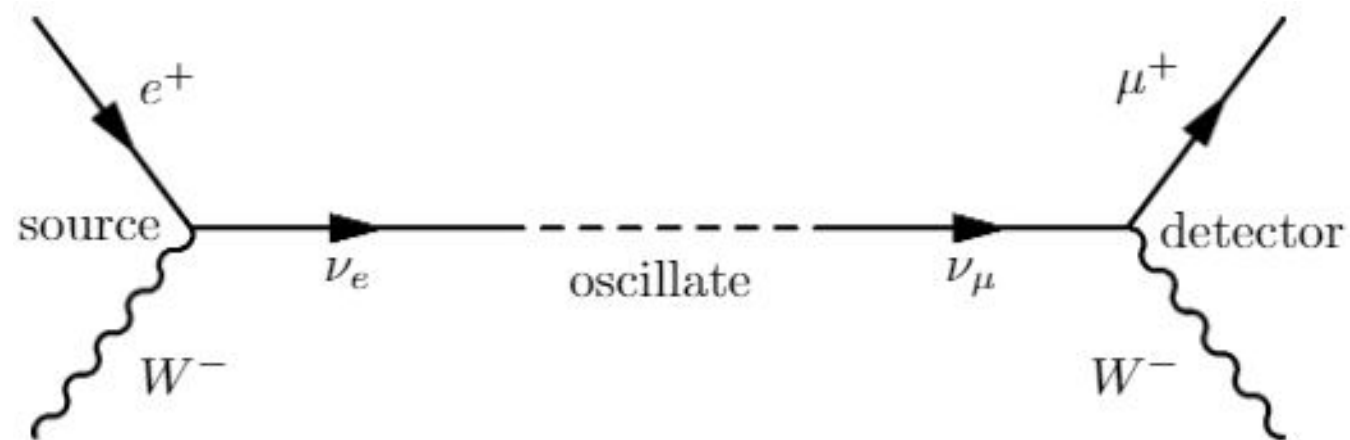
Neutrino masses from our dark sector

Phenomenology — MiniBooNE anomaly

MINERvA and CHARM-II

New physics is out there

At the electroweak scale, the Standard Model does too good a job, but...



Neutrino mixing and masses

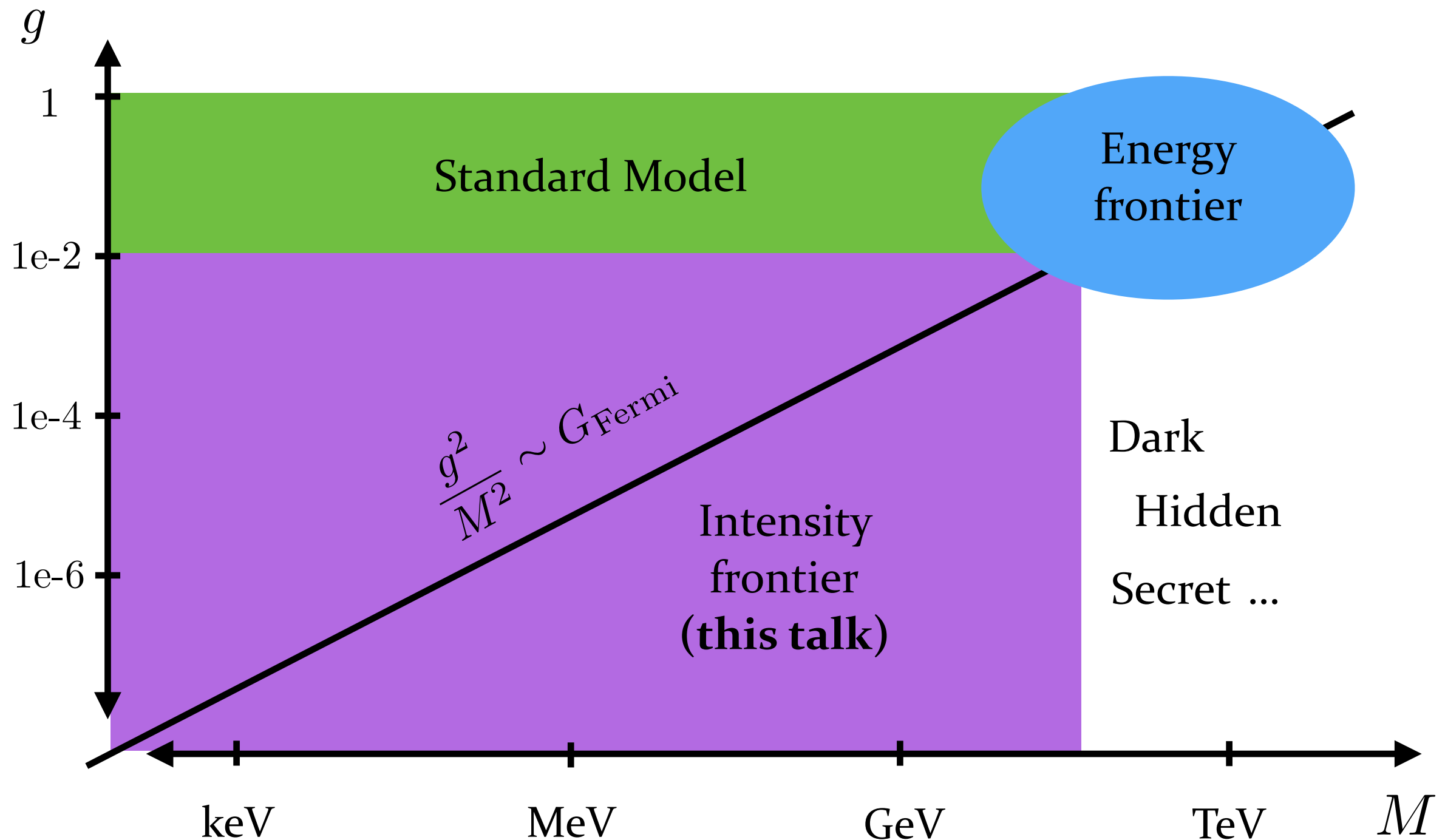
heavy neutral leptons?



Dark Matter

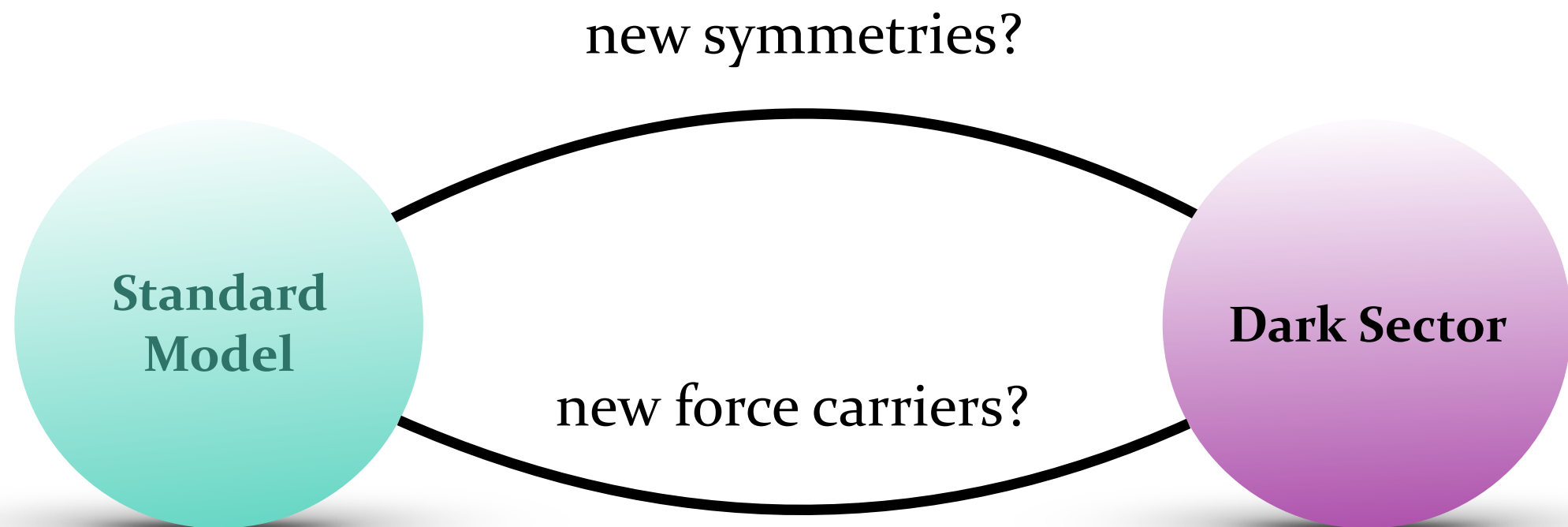
stable and neutral particles?

May be right under our noses?



Dark sector separate from the SM

Perhaps new degrees of freedom in a dark sector of their own?



DM, heavy neutrinos, ...

My dark sector wish list



- 1) Would like to have a **testable** neutrino mass model.
- 2) Dark sector: minimal, but consistent (e.g. dark photons need a mass mechanism).
- 3) Assume new interactions are confined to the dark sector (as opposed to B-L, etc).

Dark sector separate from the SM

A new gauge symmetry confined to the dark sector

$U(1)'$ with mediator X_μ

broken by the vev of a new complex scalar Φ

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A pair of heavy neutrino fields:

N

ν_D

Completely sterile

Charged under
new dark force

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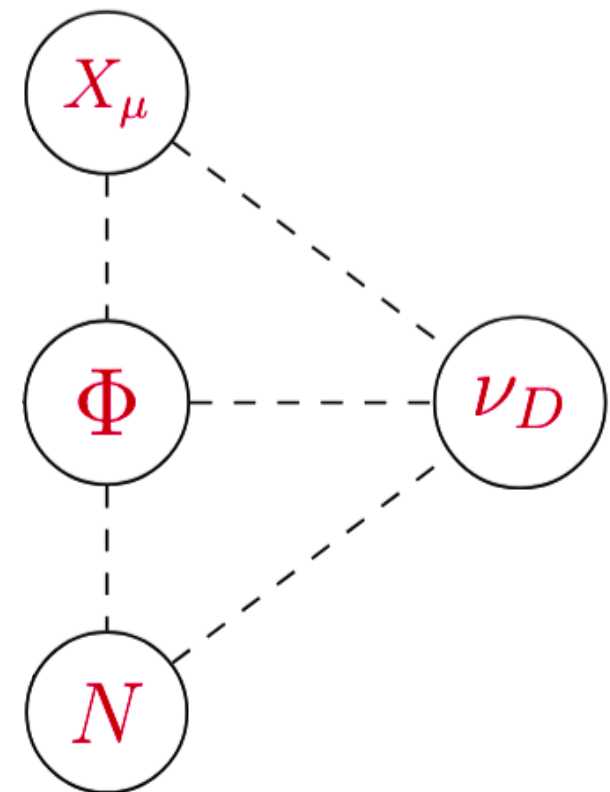
A pair of heavy neutrino fields:

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	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$	$U(1)'$
N	1	1	0	0
ν_D	1	1	0	Q
Φ	1	1	0	Q

Neutrino masses

Neutrino mass matrix with a single generation of active neutrinos

— Pair of heavy neutrinos resembles **Inverse, Extended and Linear seesaw**—

$$\mathcal{L}_{\text{mass}} \supset \frac{1}{2} \begin{pmatrix} \bar{\nu}_\alpha & \bar{N} & \bar{\nu}_D \end{pmatrix} \begin{pmatrix} 0 & m_D & \varepsilon' \\ m_D & \mu' & \Lambda \\ \varepsilon' & \Lambda & \mu \end{pmatrix} \begin{pmatrix} \nu_\alpha^c \\ N^c \\ \nu_D^c \end{pmatrix}$$

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symmetry protected

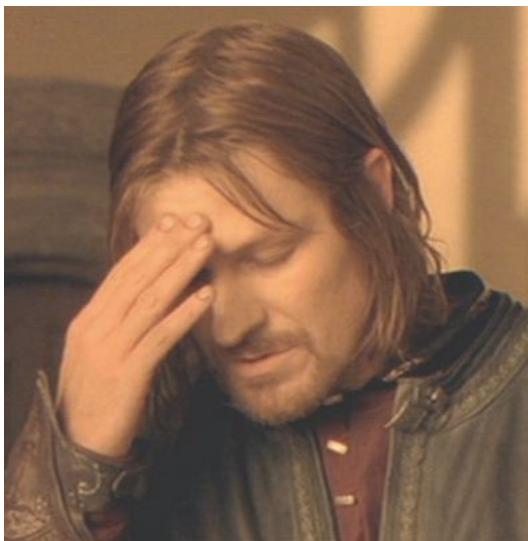
μ'
Lepton number violation

Neutrino masses

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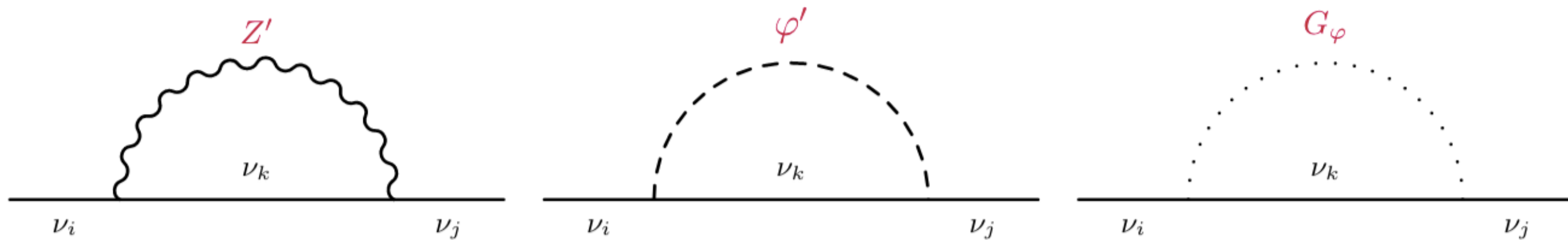
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$$m_3 = 0$$

$$m_{4,5} = \frac{\mu' \mp \sqrt{\mu'^2 + 4(\Lambda^2 + m_D^2)}}{2}$$

Neutrino masses at one-loop level



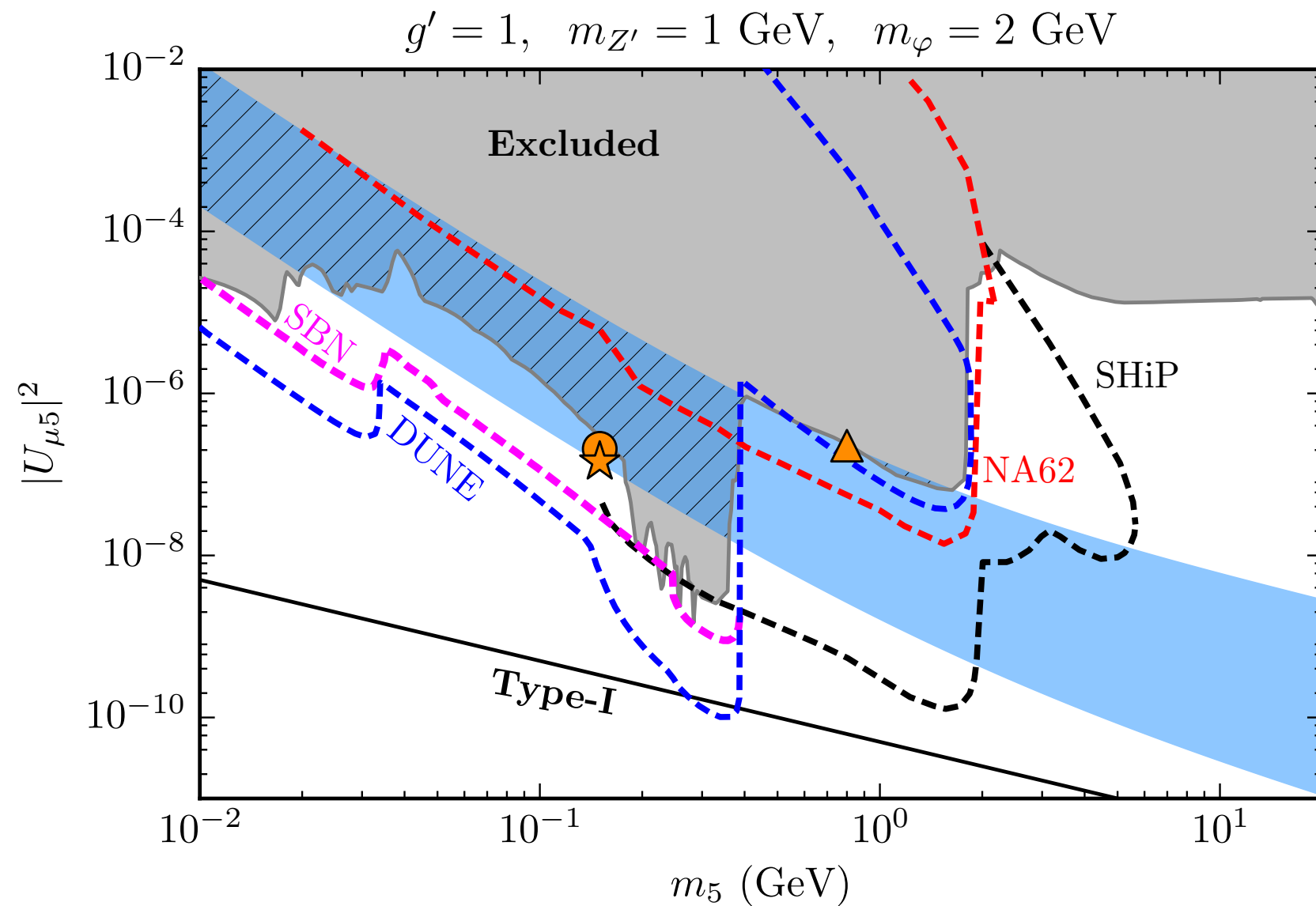
For instance, if in an Extended Seesaw scenario ($m_5 \gg m_4$) and $m_{Z'}, m_{\varphi'} \ll m_5$



$$m_3 \simeq \frac{g'^2}{16\pi^2} \frac{m_D^2}{\Lambda^2 + m_D^2} \frac{\Lambda^2}{m_{Z'}^2} \mu' \left(3 \ln \frac{m_{Z'}^2}{\mu'^2} + \ln \frac{m_{\varphi'}^2}{\mu'^2} \right)$$

New particles generate light neutrino masses radiatively!
Receive SM and BSM contributions.

Looking for the heavy neutrinos



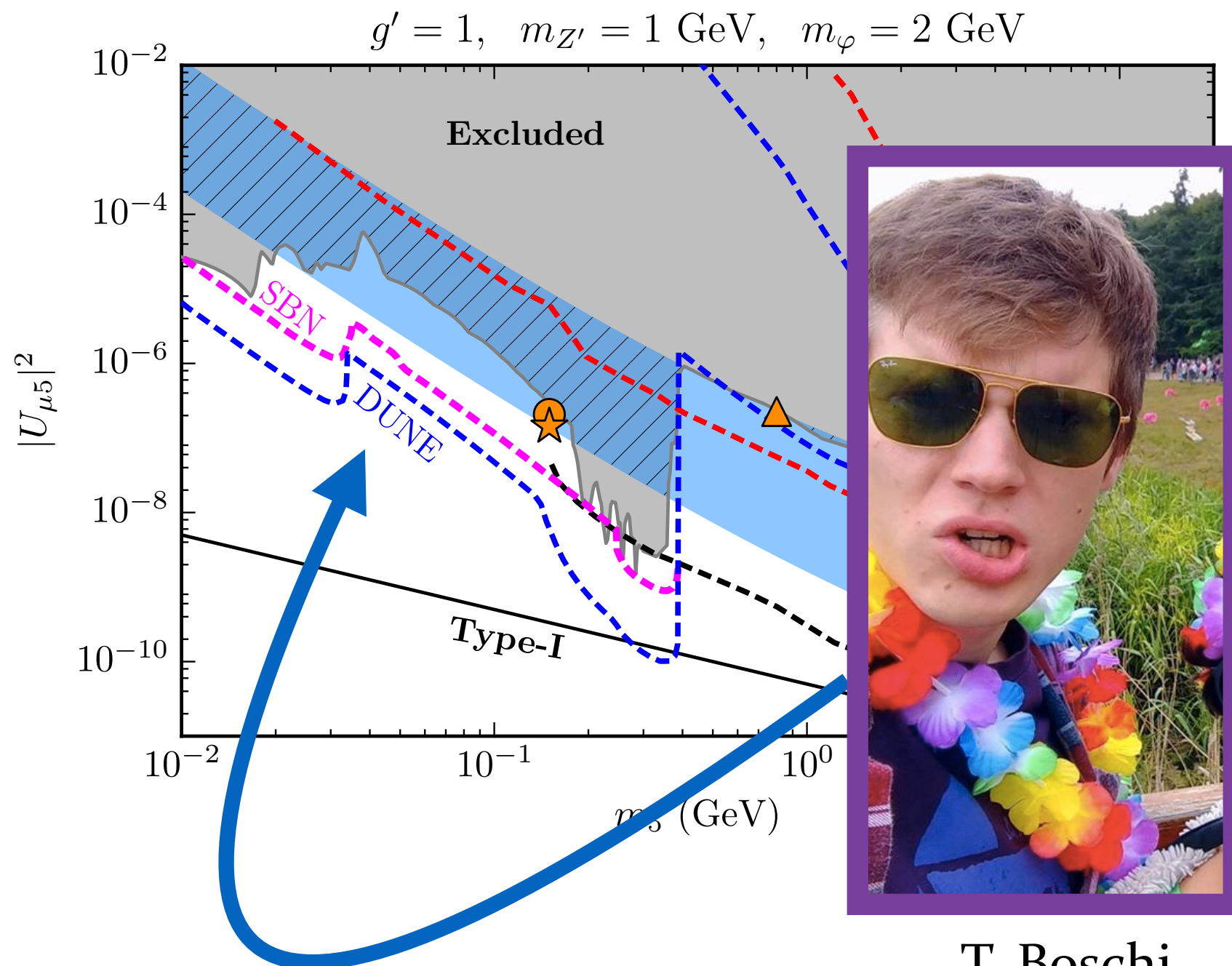
$$R = \frac{m_4}{m_5} = -\frac{U_{\alpha 5}^2}{U_{\alpha 4}^2}$$

Blue band is from:

$$m_3 = \sqrt{\Delta m_{\text{atm}}^2}$$

$$1\% < R < 99\%$$

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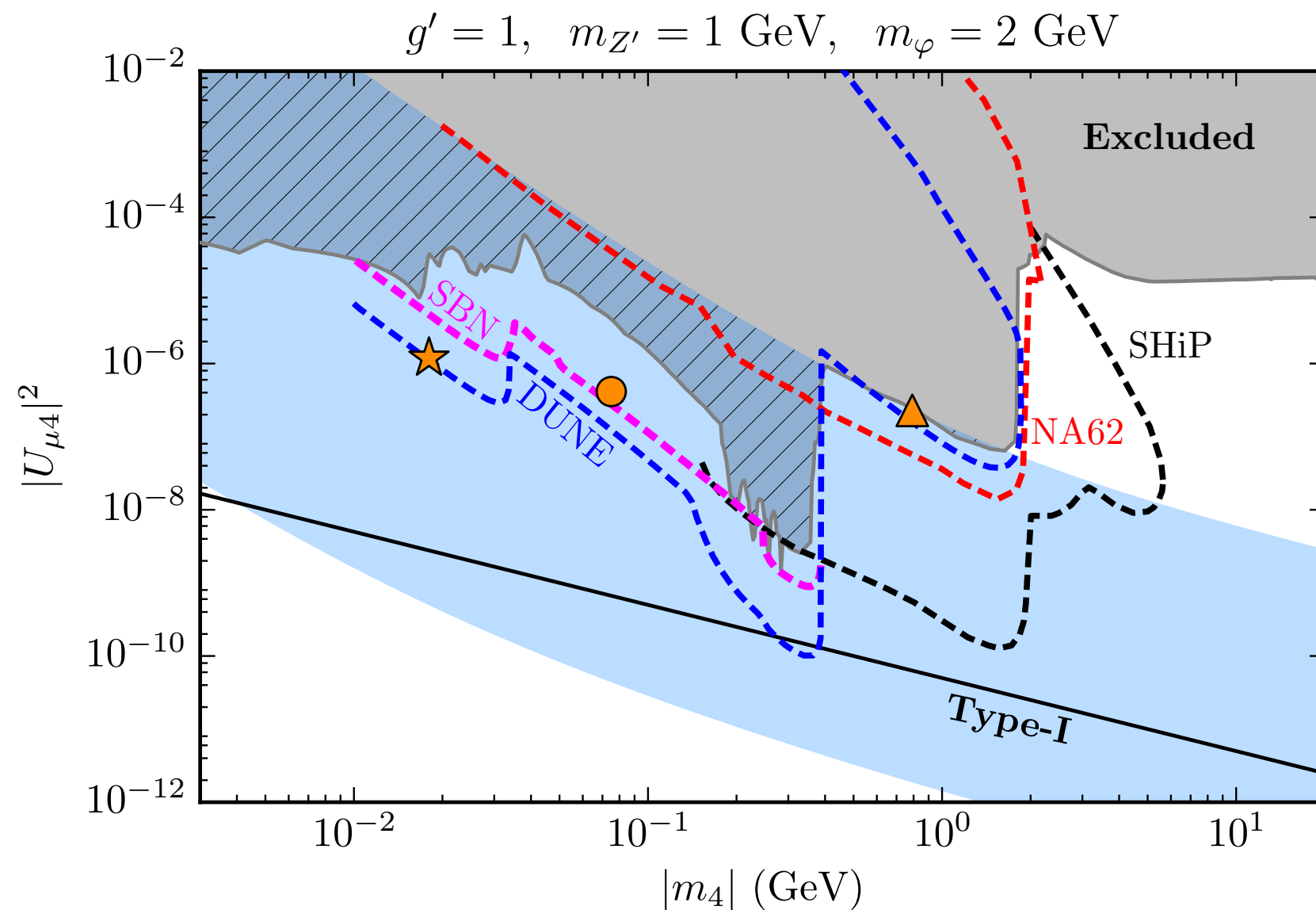
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T. Boschi

Looking for the heavy neutrinos



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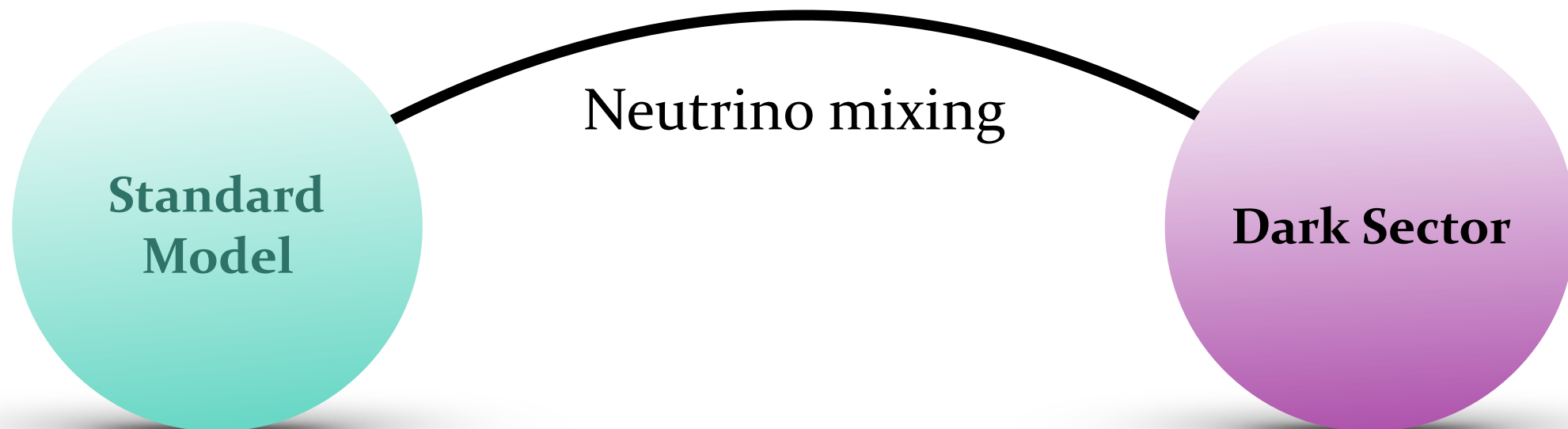
$$m_3 = \sqrt{\Delta m_{\text{atm}}^2}$$

$$1\% < R < 99\%$$

What we have so far

Heavy neutrinos interact strongly in a dark sector

Light neutrino masses arise radiatively



See also C. Diaz et al, 1712.05433, Bertuzzo et al, 1808.02500

Portals to the dark sector

Neutral BSM particle
+
SM singlet terms

$(\bar{L}\tilde{H})$	N^c	—	Neutrino portal (neutrino mixing)	} dim 4
$B_{\mu\nu}$	$X^{\mu\nu}$	—	Vector portal (kinetic mixing)	
$H^\dagger H$	$ \Phi ^2$	—	Higgs portal (scalar mixing)	

Portals to the dark sector

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+
SM singlet terms

$(\bar{L}\tilde{H}) N^c$ — Neutrino portal (neutrino mixing)

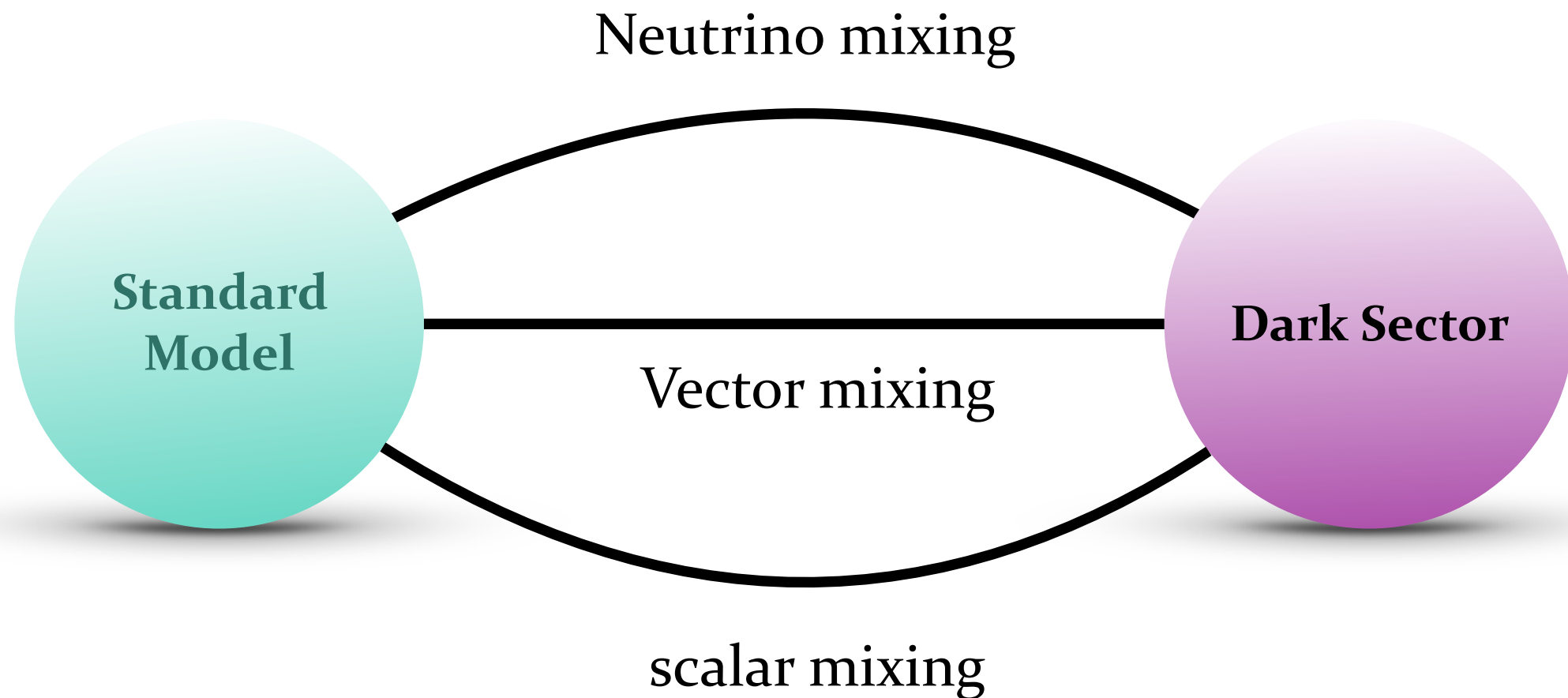
No reason why should not be there!

$B_{\mu\nu} X^{\mu\nu}$ — Vector portal (kinetic mixing)

$H^\dagger H |\Phi|^2$ — Higgs portal (scalar mixing)

dim 4

Full picture



The dark sector interactions “leak” to the SM via the three portals.

Portal couplings are **theoretically** unconstrained and expected to be large!

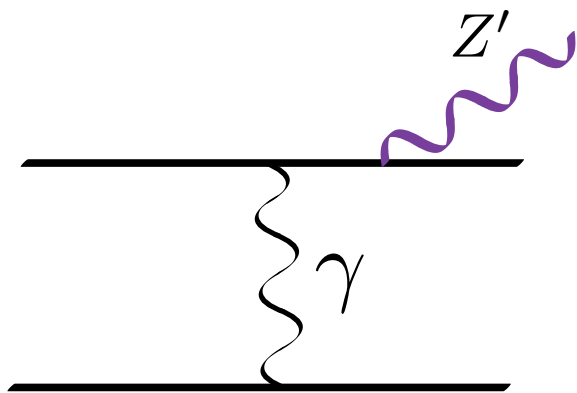
Vector portal

$$\frac{\sin \chi}{2} B_{\mu\nu} X^{\mu\nu}$$

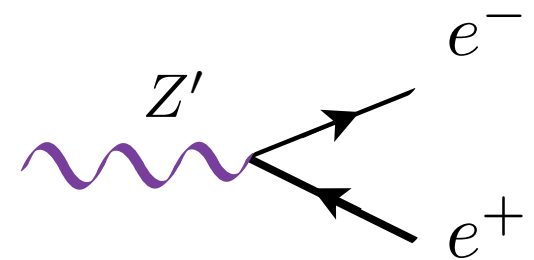
For a light boson $\rightarrow$ couples only to electric charge (dark photon).



At beam dumps and colliders:



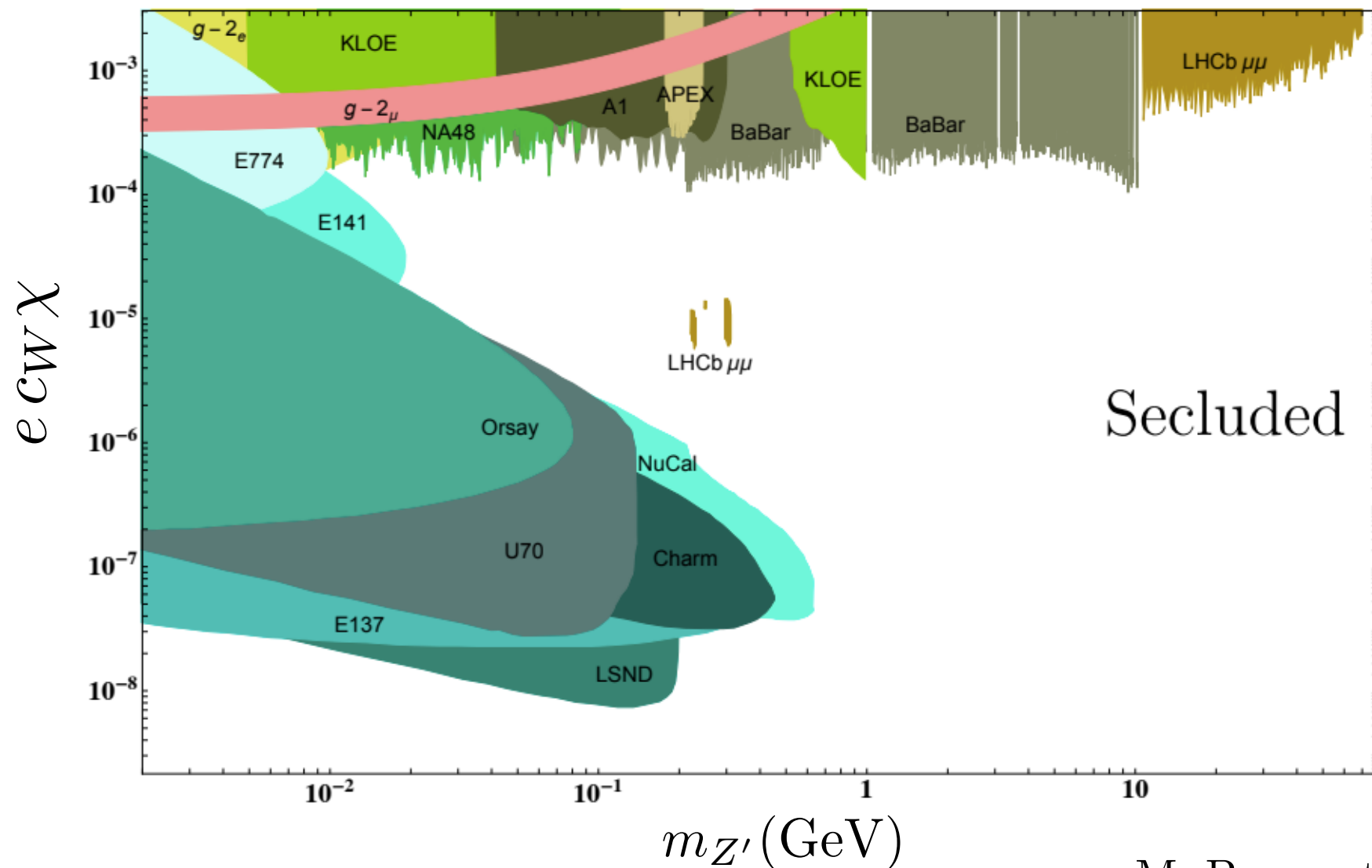
.....



Dark photon

$$\frac{\sin \chi}{2} B_{\mu\nu} X^{\mu\nu}$$

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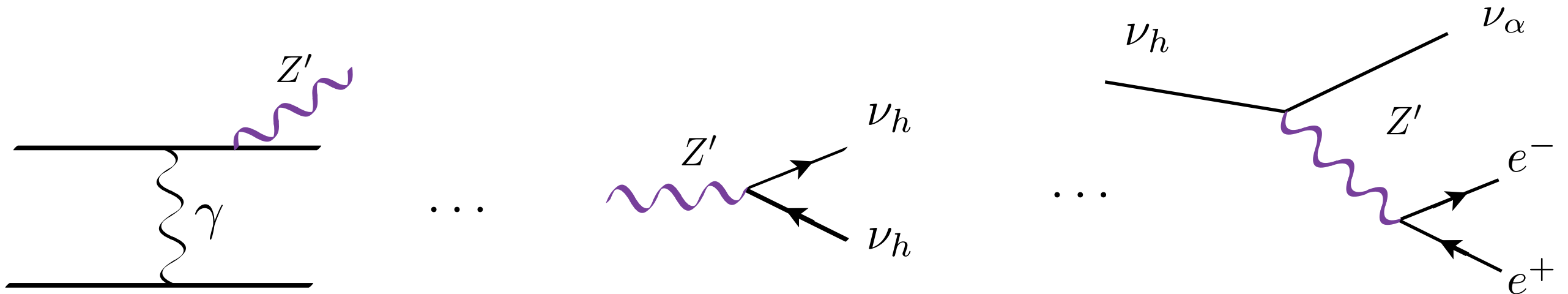
M. Bauer et al., 1803.05466

A *new* dark photon

$$\frac{\sin \chi}{2} B_{\mu\nu} X^{\mu\nu}$$

For a light boson $\rightarrow$ couples only to electric charge (dark photon).

But in our full model, the dark photon mainly decays to heavy neutrinos!

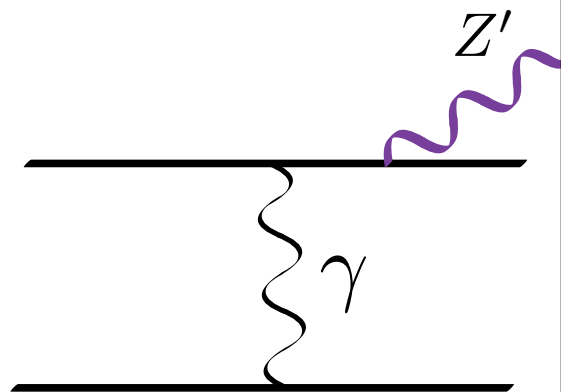


Three-body kinematics, no longer reconstructs the invariant mass!

A *new* dark photon

For a light

But in our



Three-b

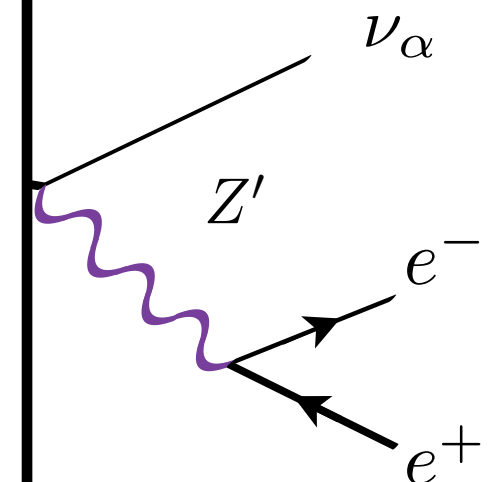
A **neutrinophilic dark photon**
hiding in plain sight...

Weakened bounds
— final states do not pass the current vetoes —

Similar considerations apply for dark scalar.

photon).

neutrinos!



variant mass!

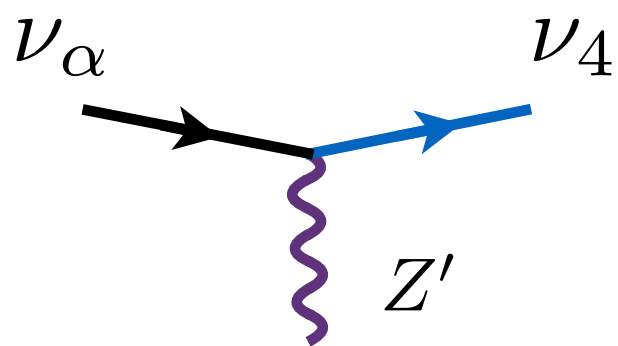
New signatures at neutrino experiments

For simplicity, let us take a phenomenological approach and consider a single **dark** neutrino.

This can be achieved in an extended seesaw limit of the model discussed before.

New signatures at neutrino experiments

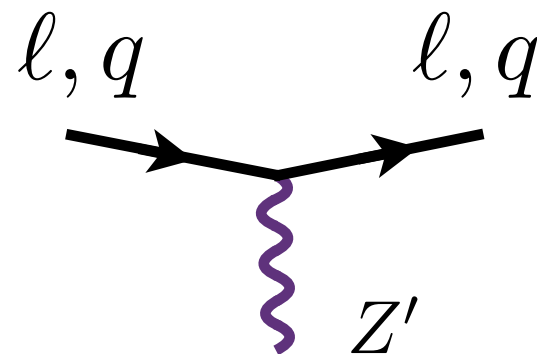
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$$U_{\alpha 4}^* g'$$

SM + heavy neutrino

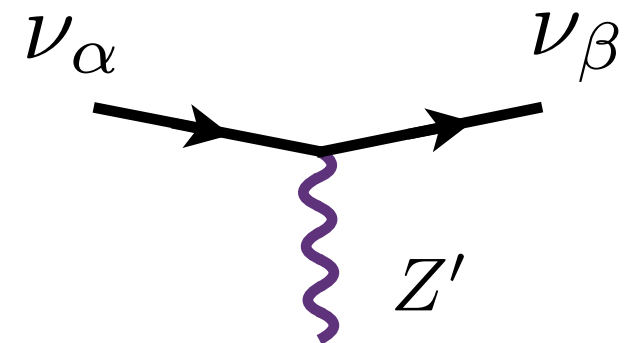
Mixing + $O(1)$ coupling



$$e \chi q_f$$

Quarks and charged leptons

Kinetic mixing



$$U_{\alpha 4}^* U_{\beta 4} g'$$

SM neutrinos only

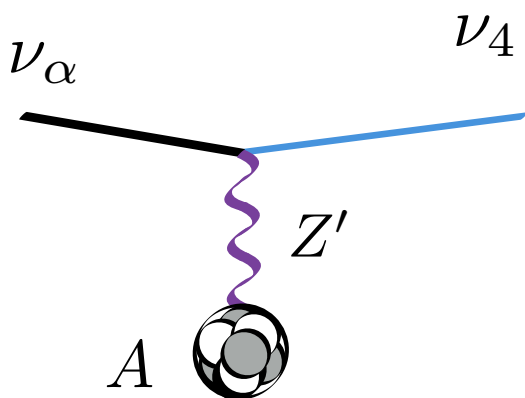
Two mixings needed

Neutrino scattering

Interest from the point of view of neutrino scattering and experimental anomalies

Neutrino NC cross sections are enhanced due to the light mediator:

“Upscattering”


$$\propto g'^2 \chi^2 |U_{\alpha 4}|^2 \frac{1}{m_{Z'}^4}$$

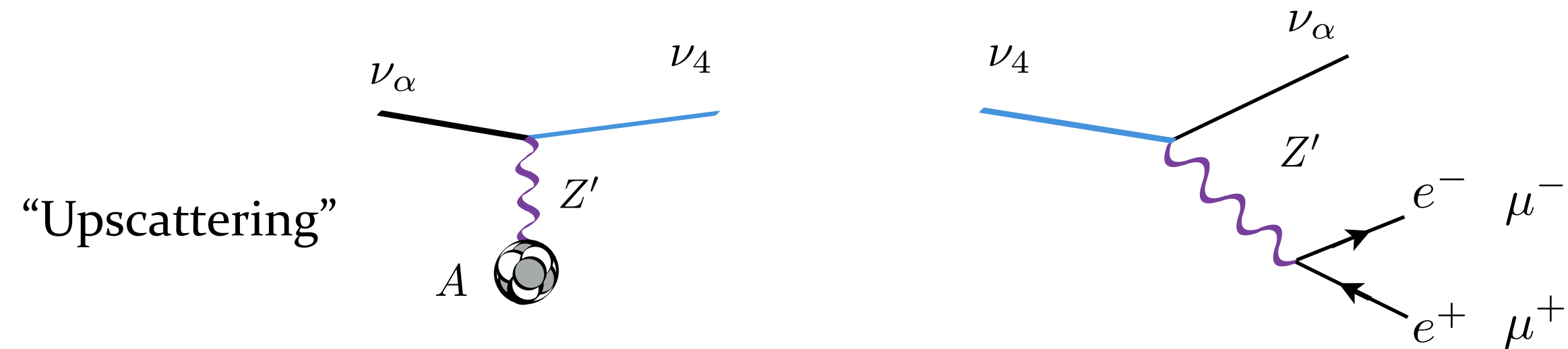
E. Bertuzzo et al, PhysRevLett.121.241801

P. Ballett et al, 1808.02915

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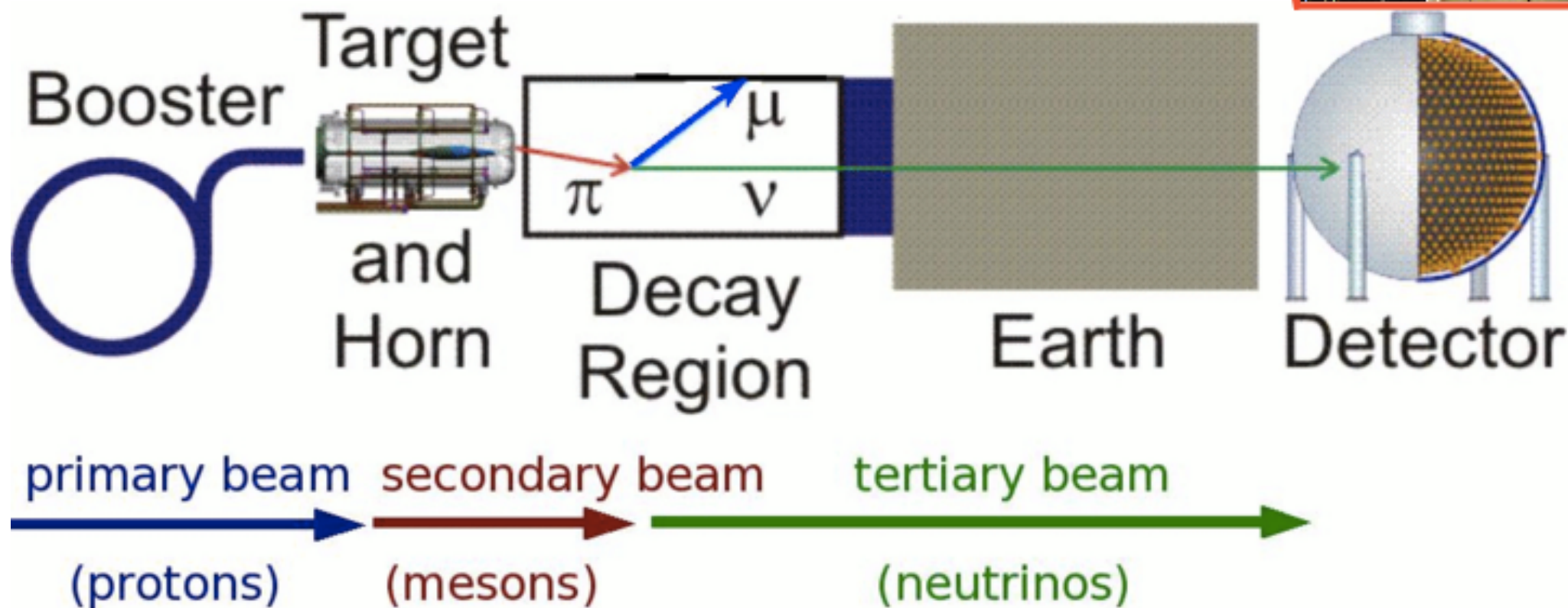
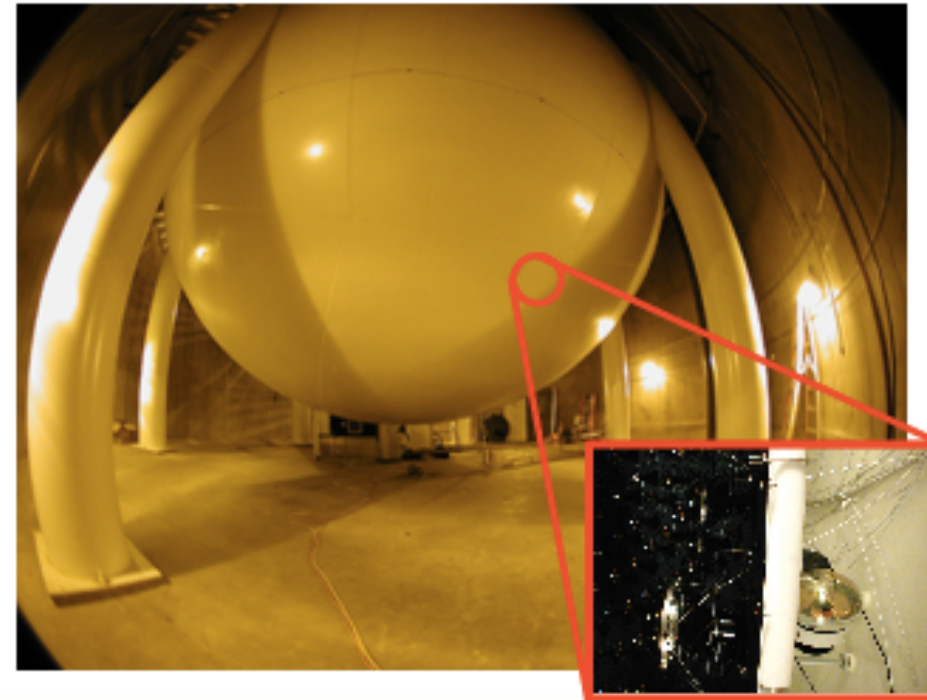
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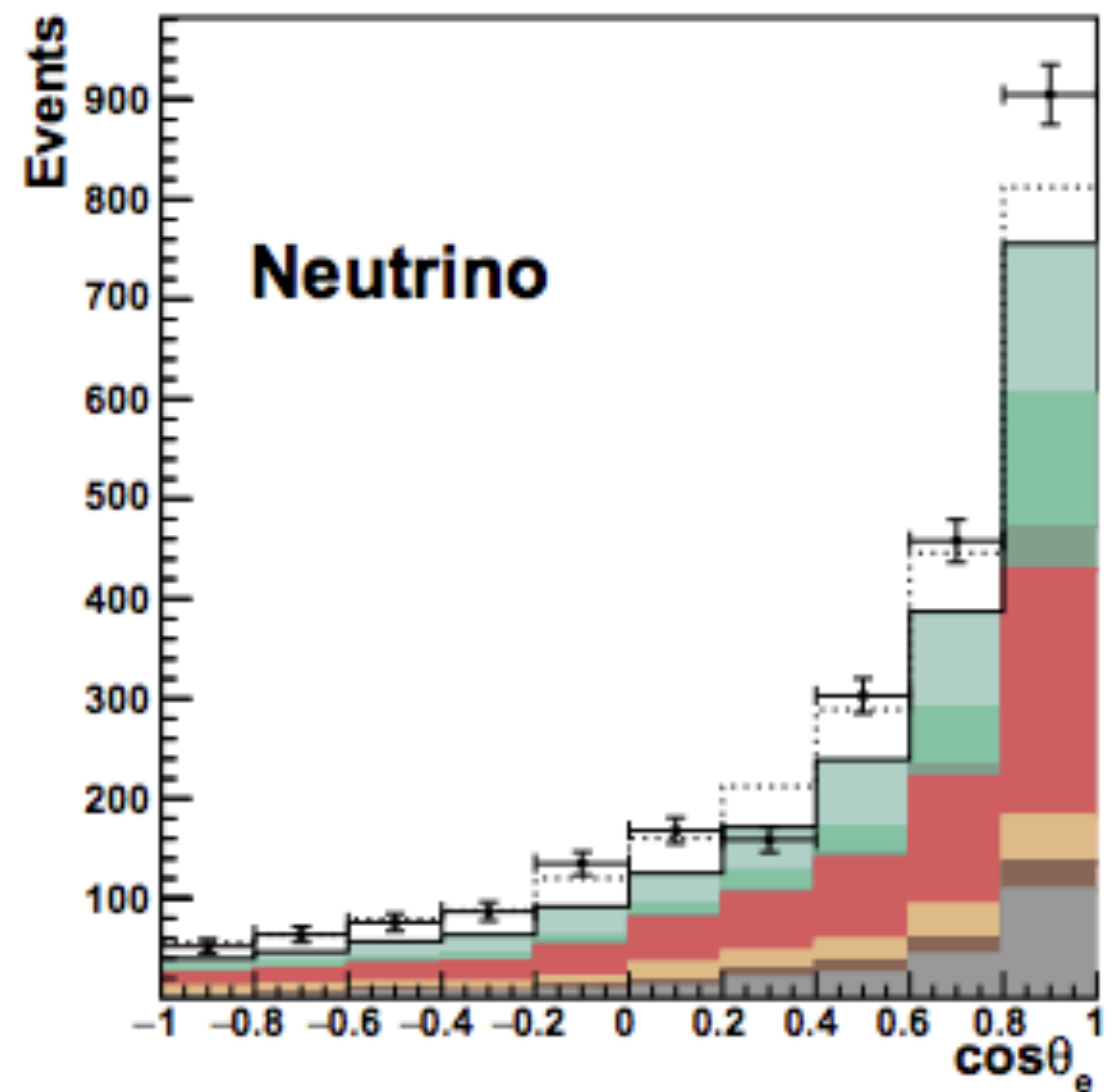
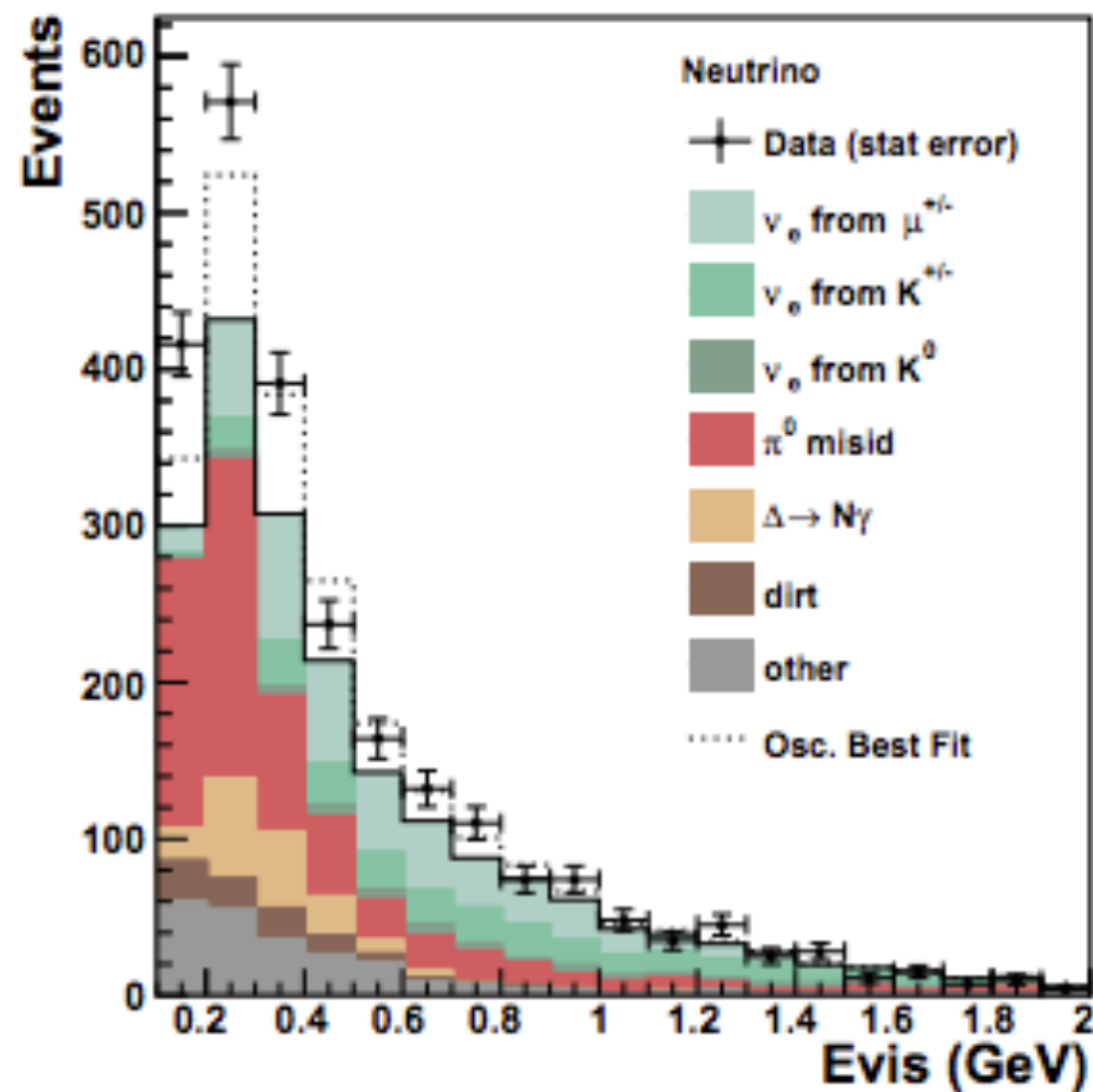
MiniBooNE



818 t

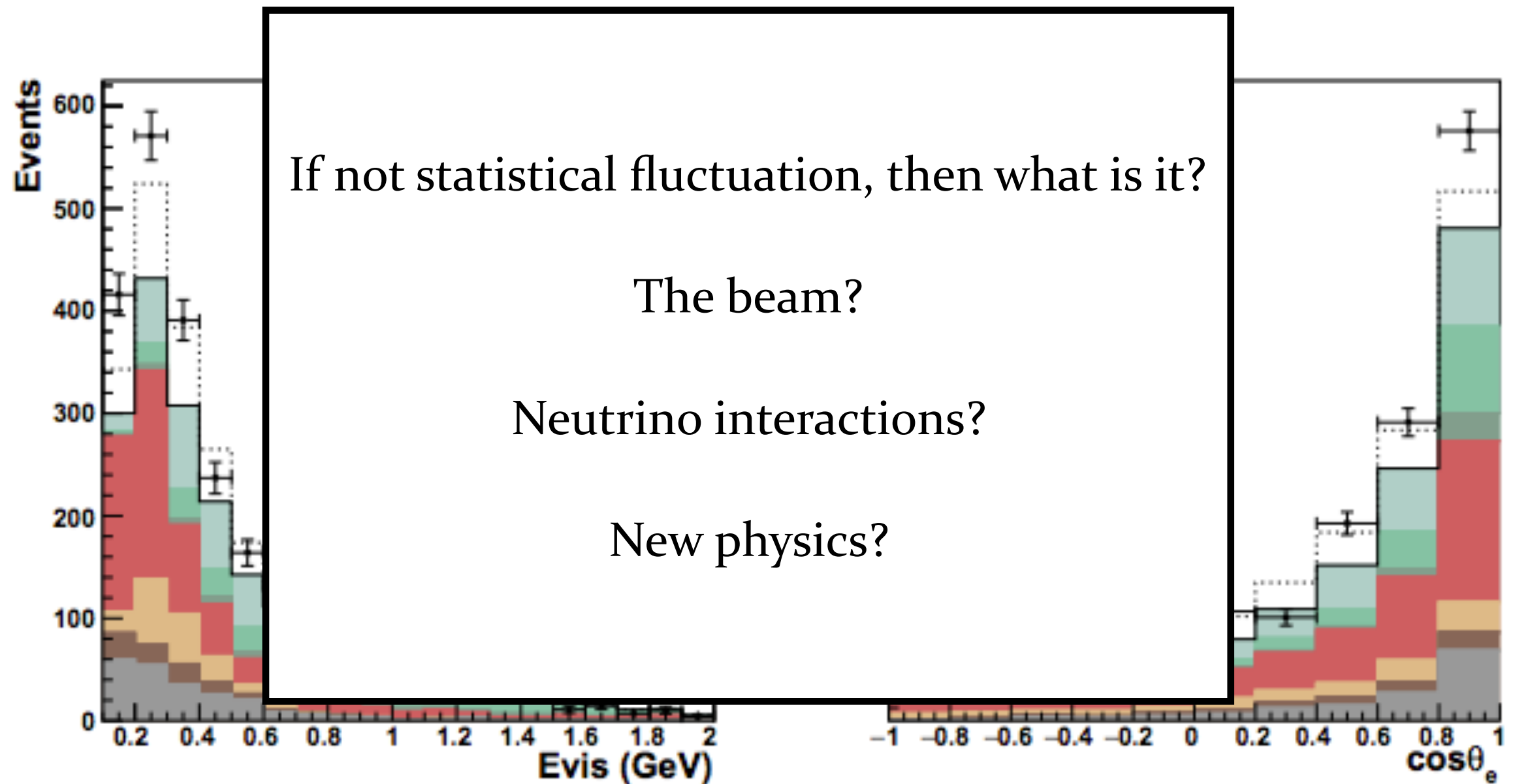
MiniBooNE

4.7 σ excess observed in neutrino + antineutrino modes
— data/MC disagreement beyond statistical doubt —



MiniBooNE

4.7 σ excess observed in neutrino + antineutrino modes
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We need something new

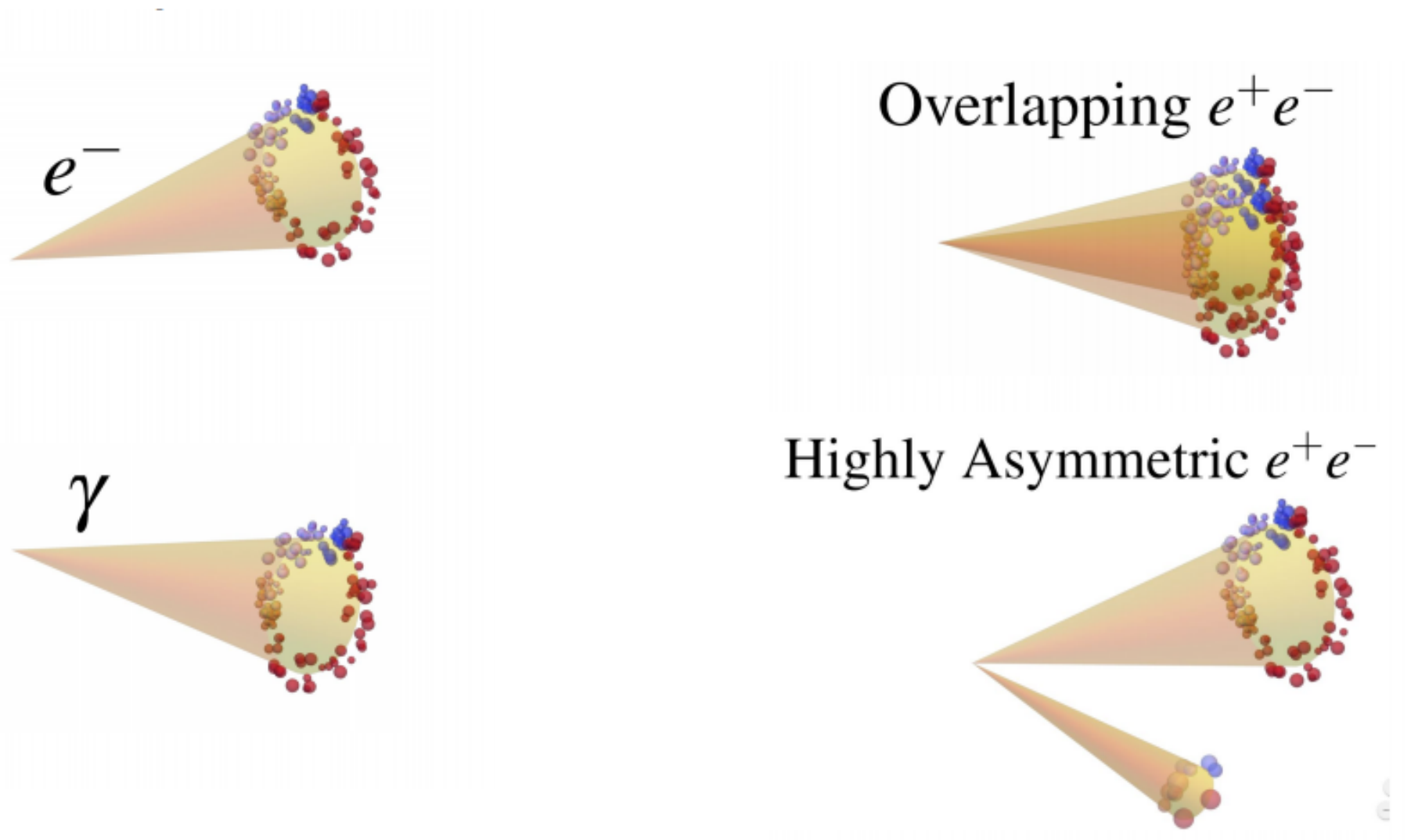
“If the data doesn’t agree with the null hypothesis or the alternative hypothesis, some say you need more data, while some say you need more hypotheses.”

Maury Goodman, arXiv:1901.07068

At this point the anomaly requires **non-minimal scenarios**.

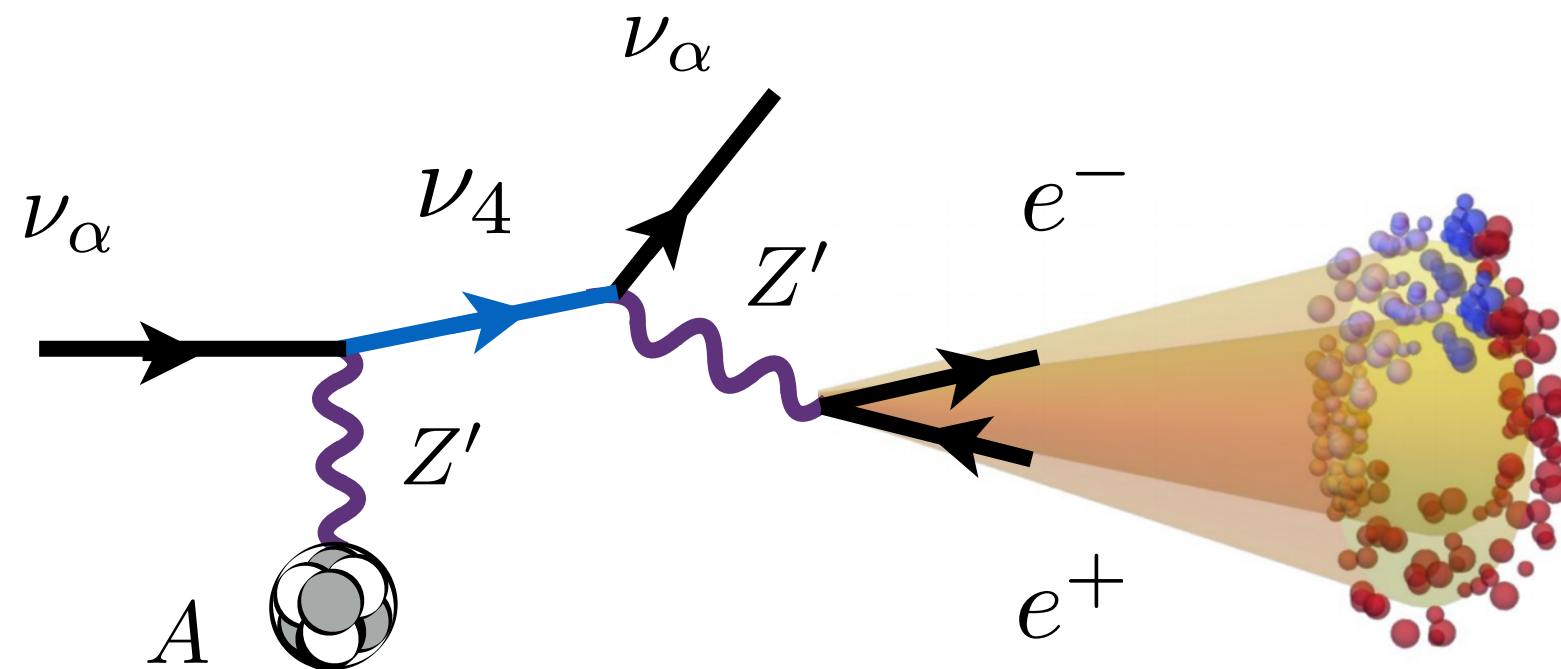
MiniBooNE

What kind of events contribute to the excess:



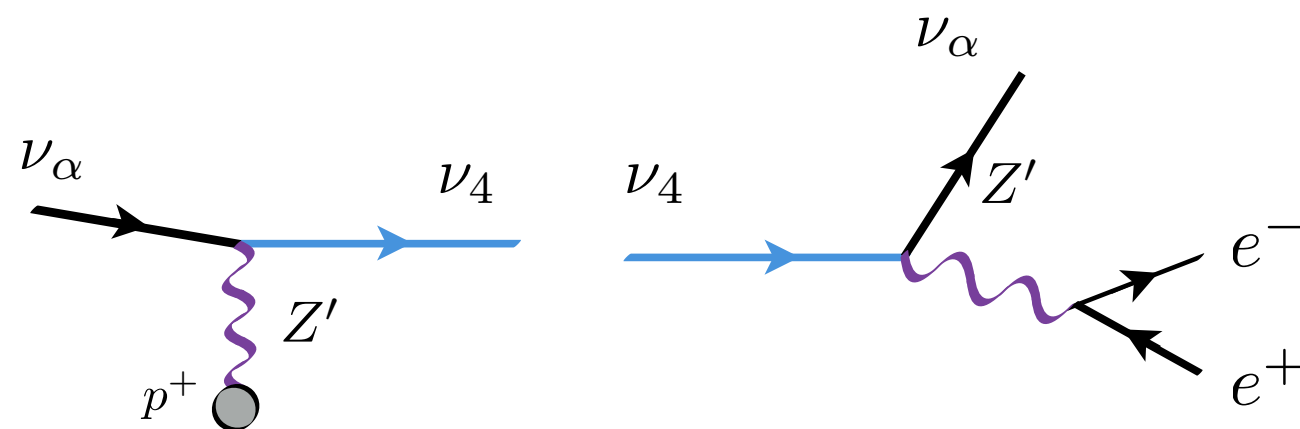
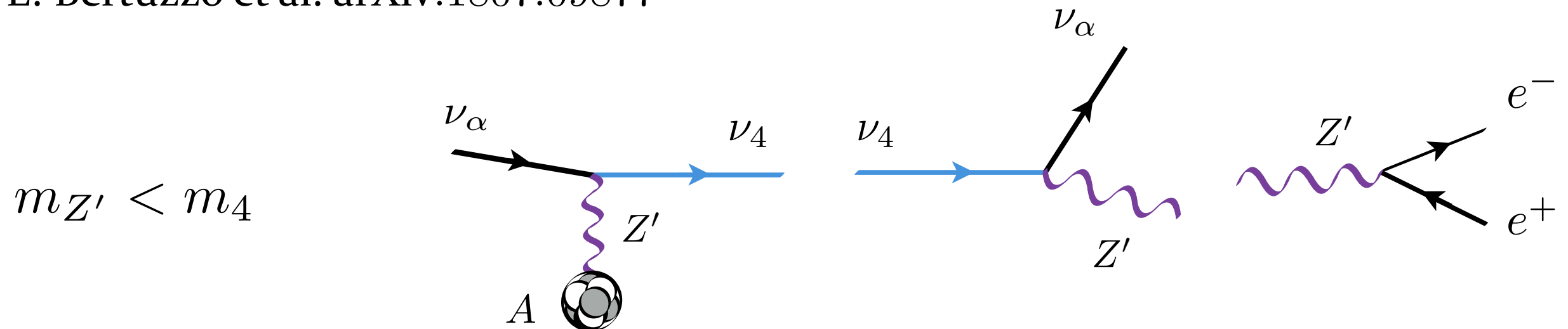
The Signal at MiniBooNE

Neutrinos upscatter into heavy state, which immediately decays into an overlapping pair of electrons.



Two realisations

E. Bertuzzo et al. arXiv:1807.09877



P. Ballett et al. arXiv:1808.02915

$$m_{Z'} > m_4$$

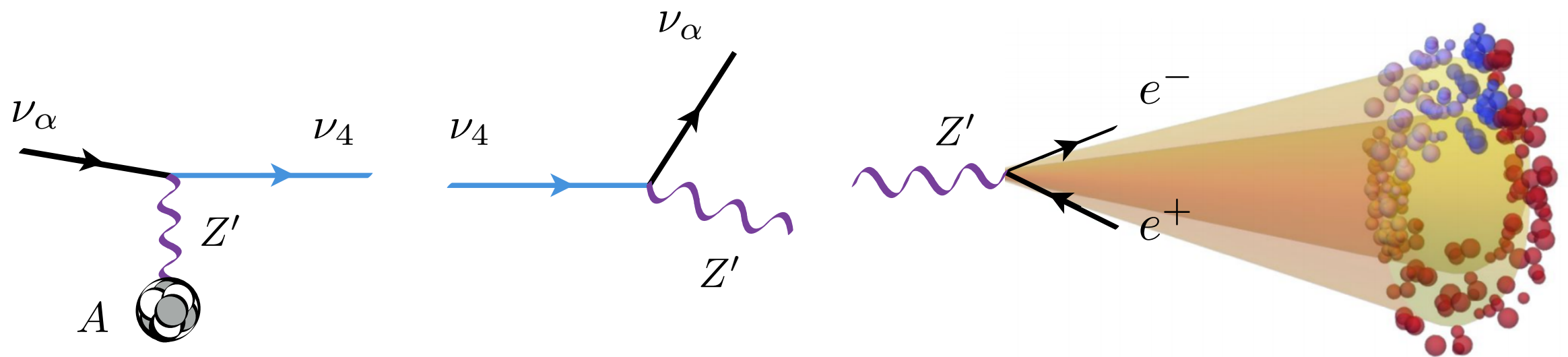
Light dark photon case

E. Bertuzzo et al. arXiv:1807.09877

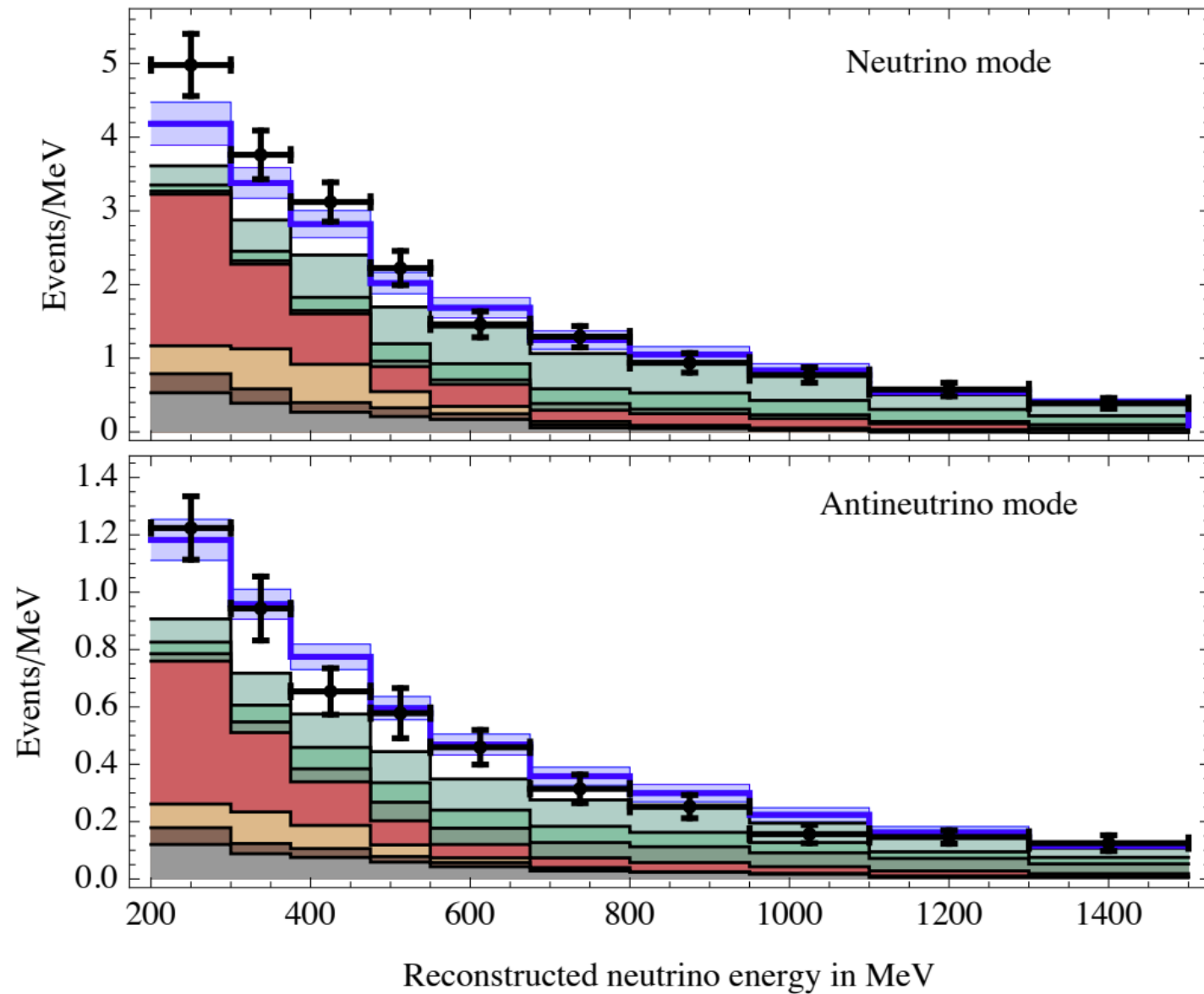
$$m_{Z'} < m_4$$

Coherent signal

— single boosted EM shower —



Light dark photon case

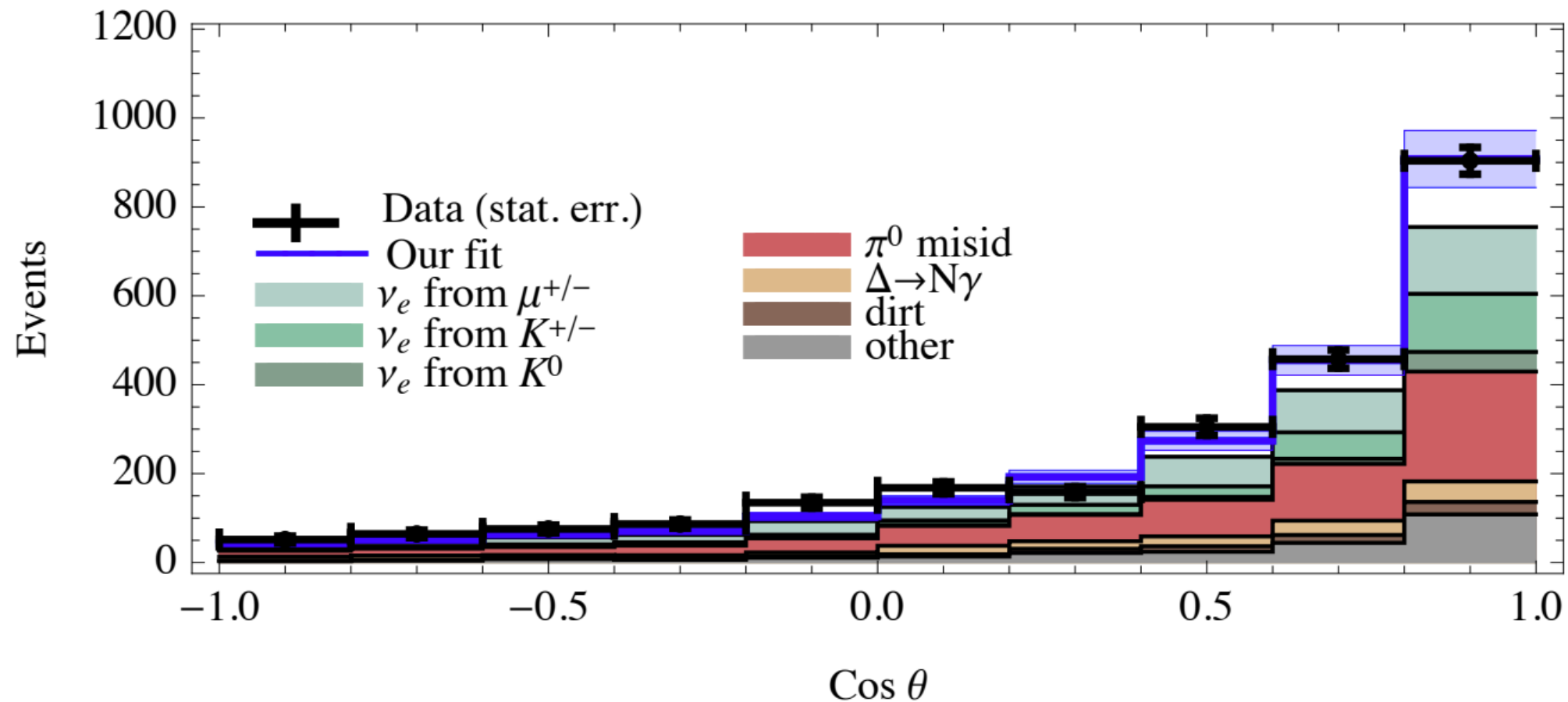


$$m_{Z'} < m_4$$

The predictions of our benchmark point $m_{N_D} = 420$ MeV, $m_{Z_D} = 30$ MeV, $|U_{\mu 4}|^2 = 9 \times 10^{-7}$, $\alpha_D = 0.25$ and $\alpha \epsilon^2 = 2 \times 10^{-10}$ are also shown as the blue lines.

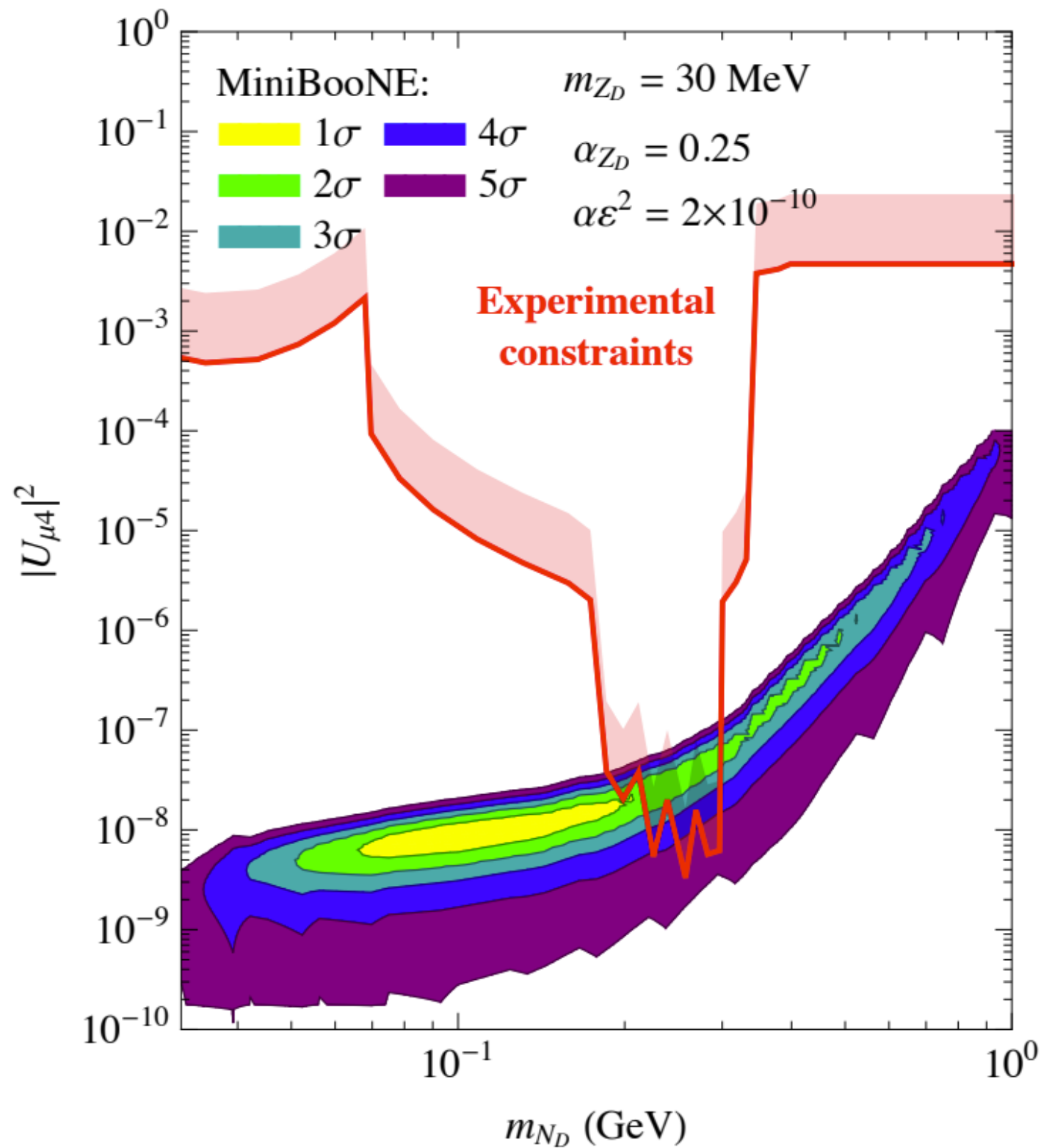
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Light dark photon case



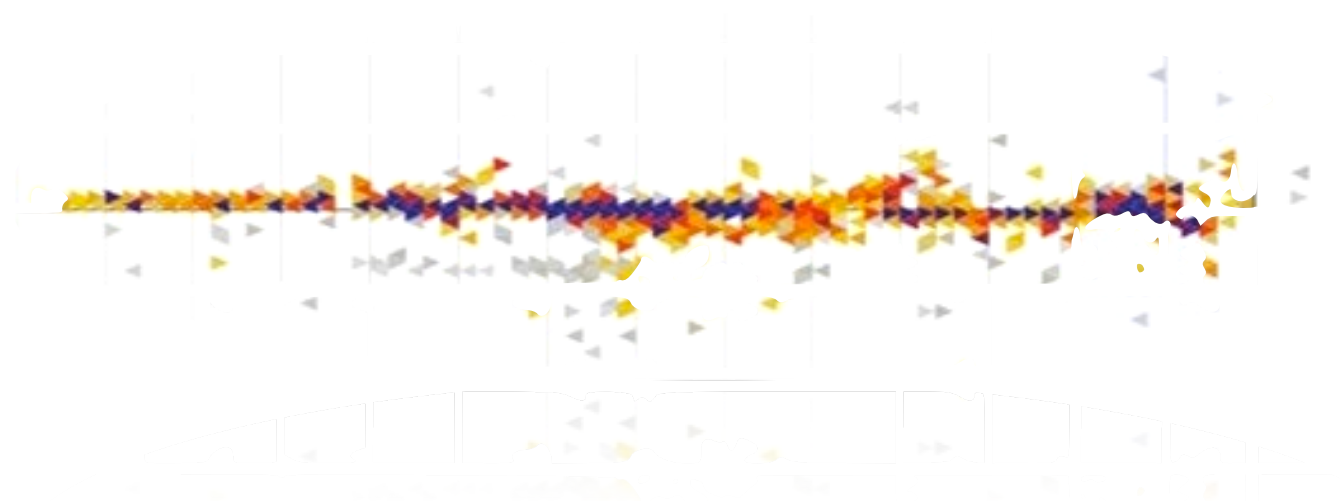
$$m_{Z'} < m_4$$

Testing light dark photon case

Coherent signal

$$m_{Z'} < m_4$$

— single boosted EM shower —



See π^0 and coherent gammas. Also recent search for coherent gammas in T2K.

Testing light dark photon case

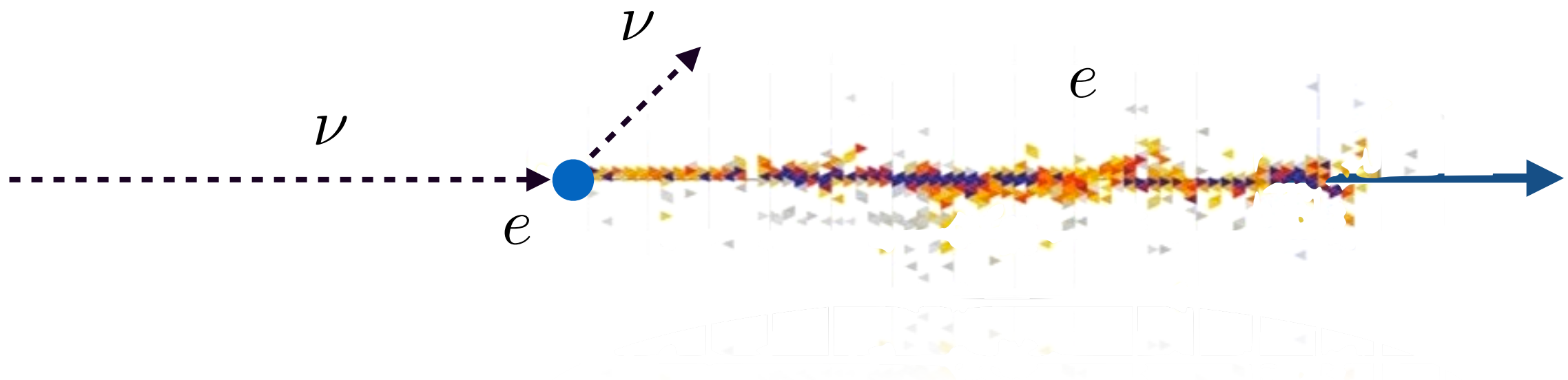
Coherent signal

$$m_{Z'} < m_4$$

— single boosted EM shower —

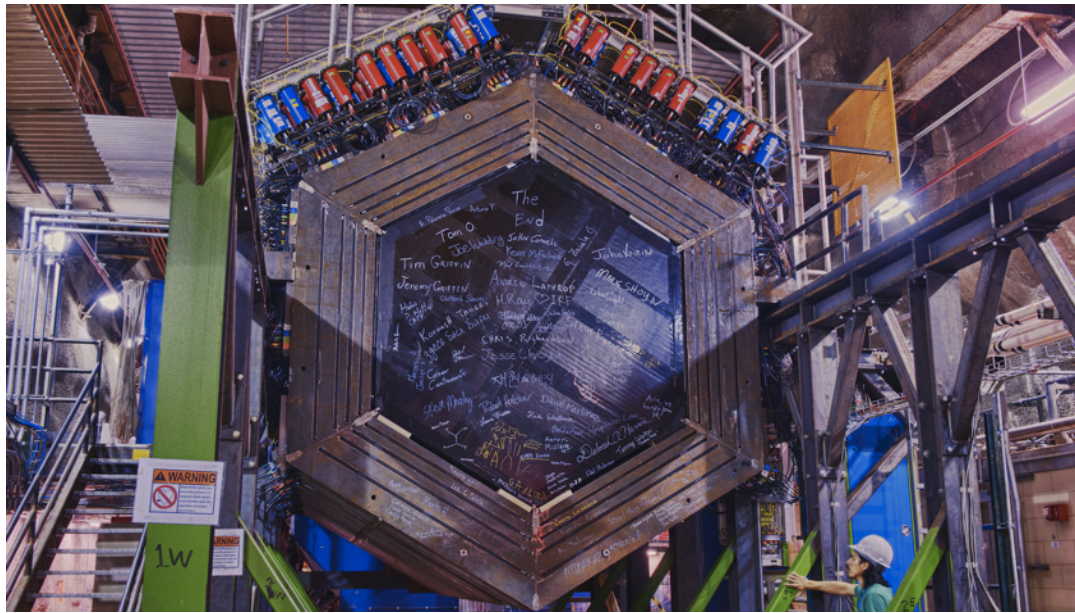


neutrino-electron scattering.



Rare SM process with similar signature.

Neutrino-electron scattering



MINERvA

NuMI beam FNAL

Low-Energy data on nu-e scattering

8 t

CHARM-II

WANF beam at CERN back in the 90s

692 t

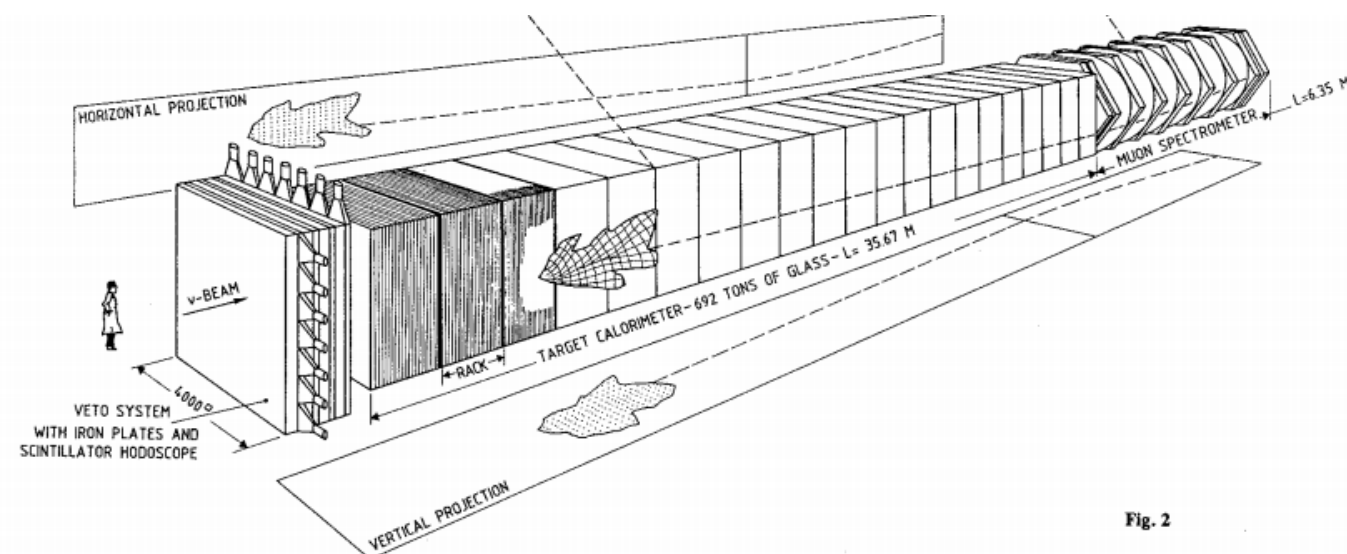
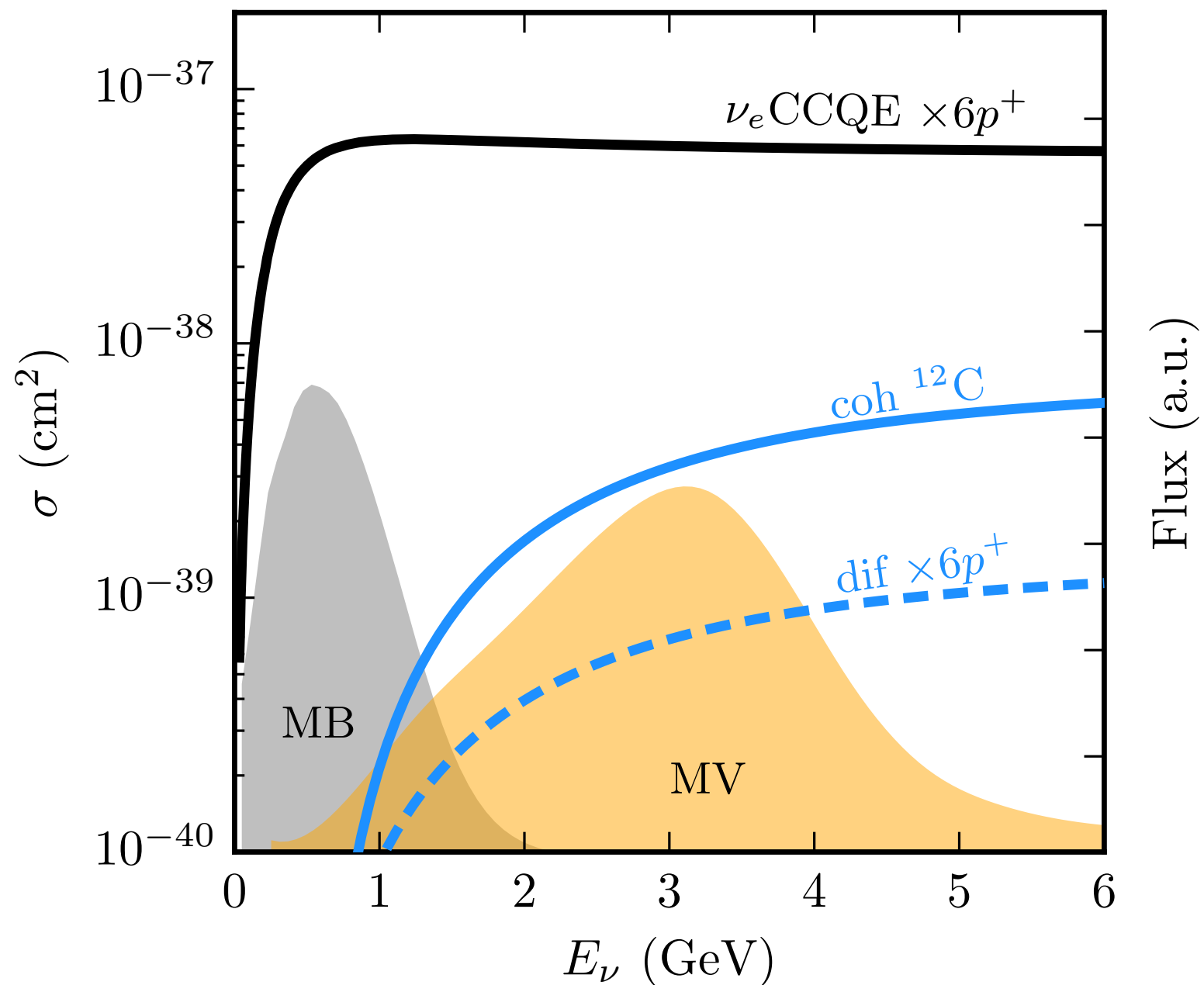
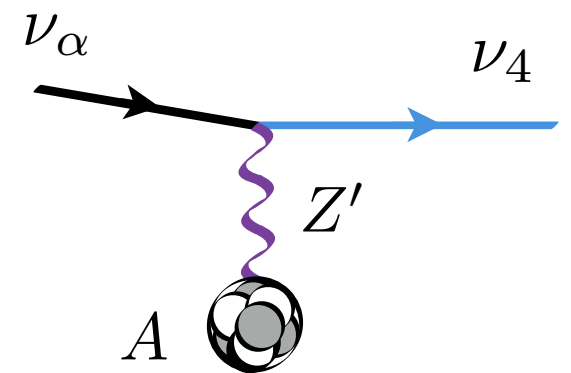


Fig. 2

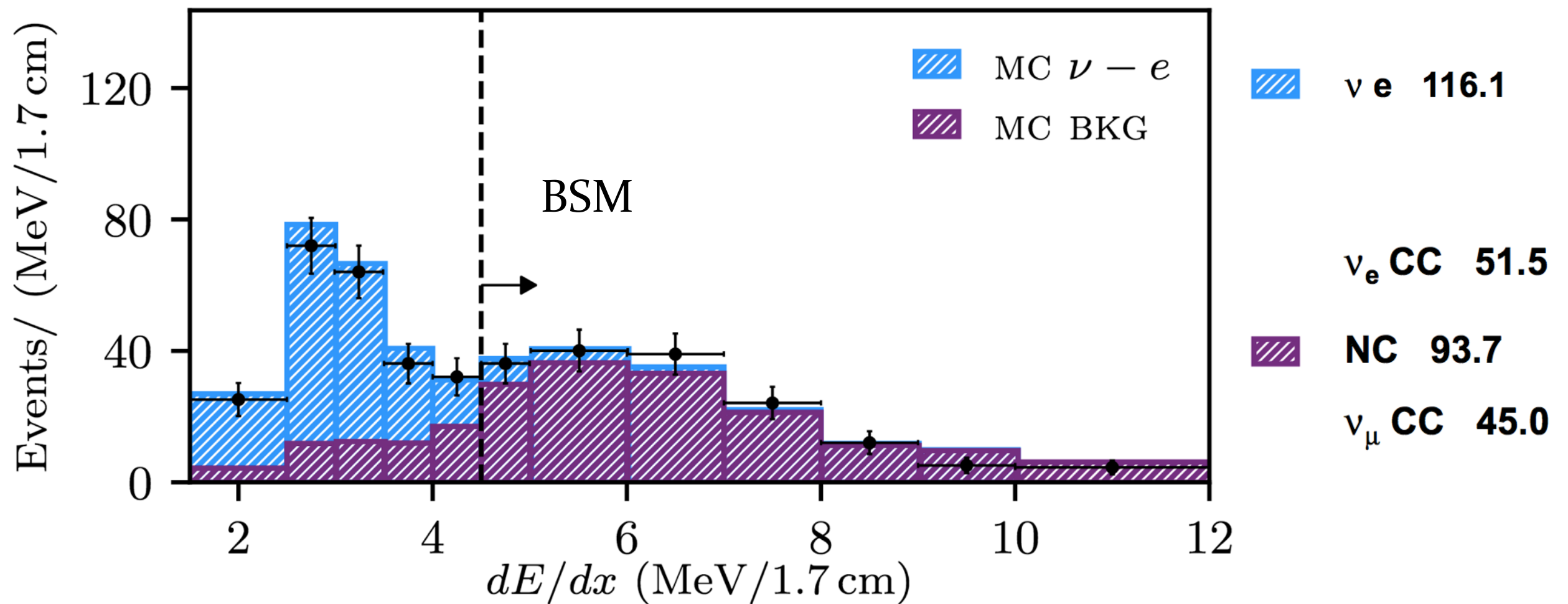
BSM cross sections



New physics is large
at larger energies!



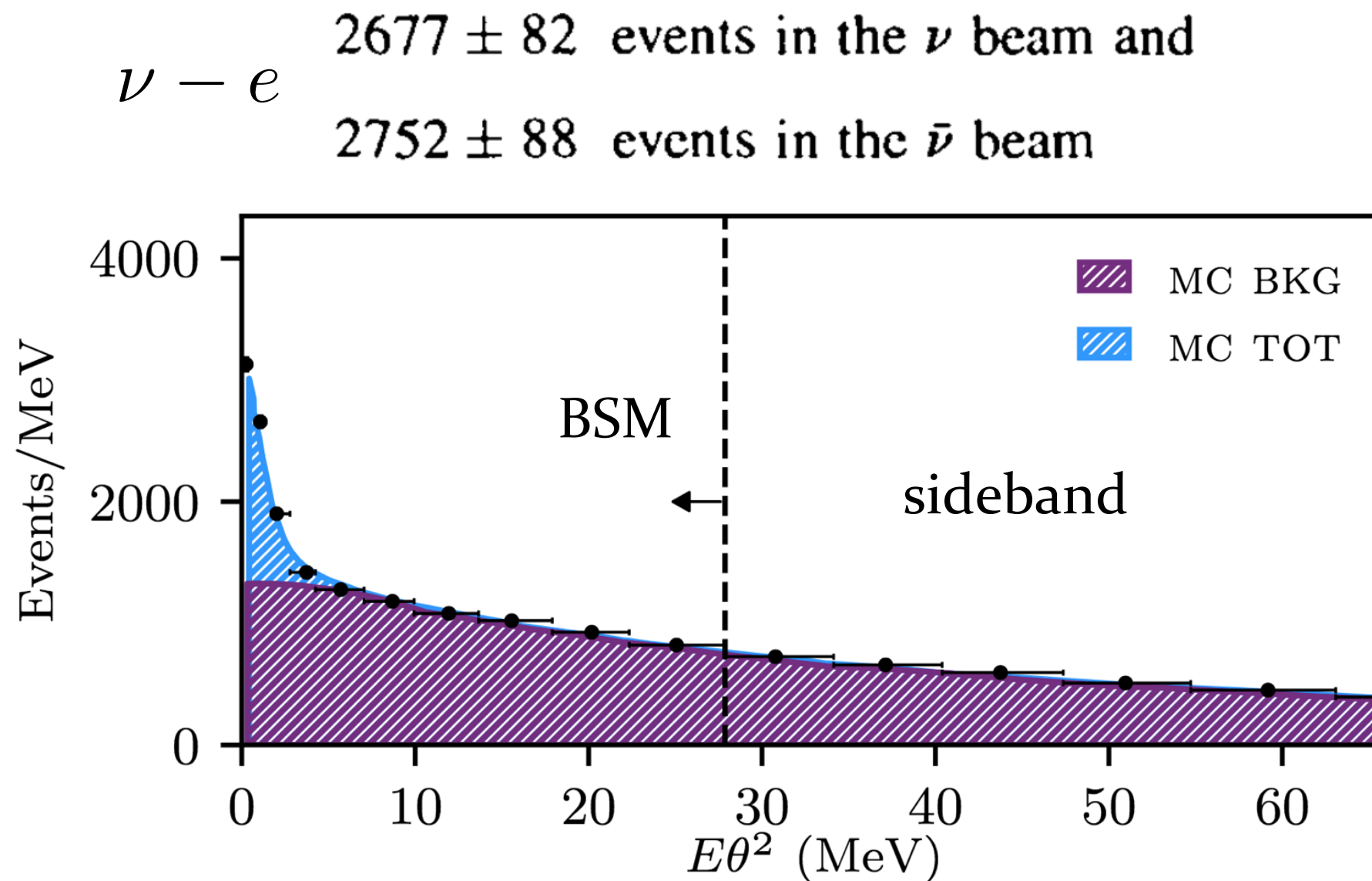
MINERvA



Backgrounds are small but are tuned! Based on tuning, assign a 30% uncertainty.

BSM signal typically overshoots bkg:
232 BSM events!

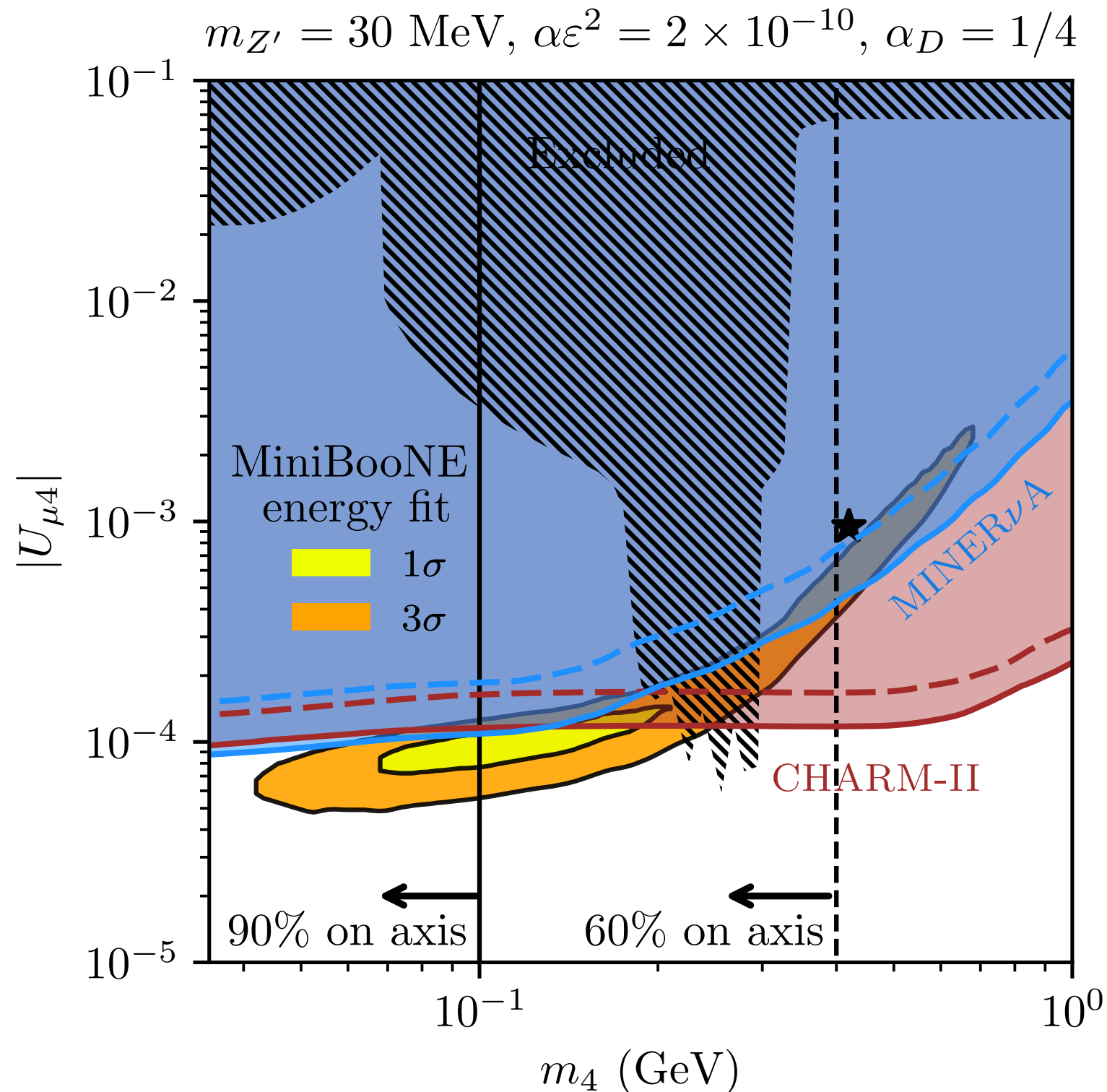
CHARM-II



Measure backgrounds at large angles (3% uncertainties).

1.43×10^5 BSM events in signal region!

Constraints on light boson case



$$m_{Z'} < m_4$$

Vary background systematics:

from 30% — 100% for MINERvA

from 3 — 10% for CHARM-II

Small mass region:
disagreement with angular data!

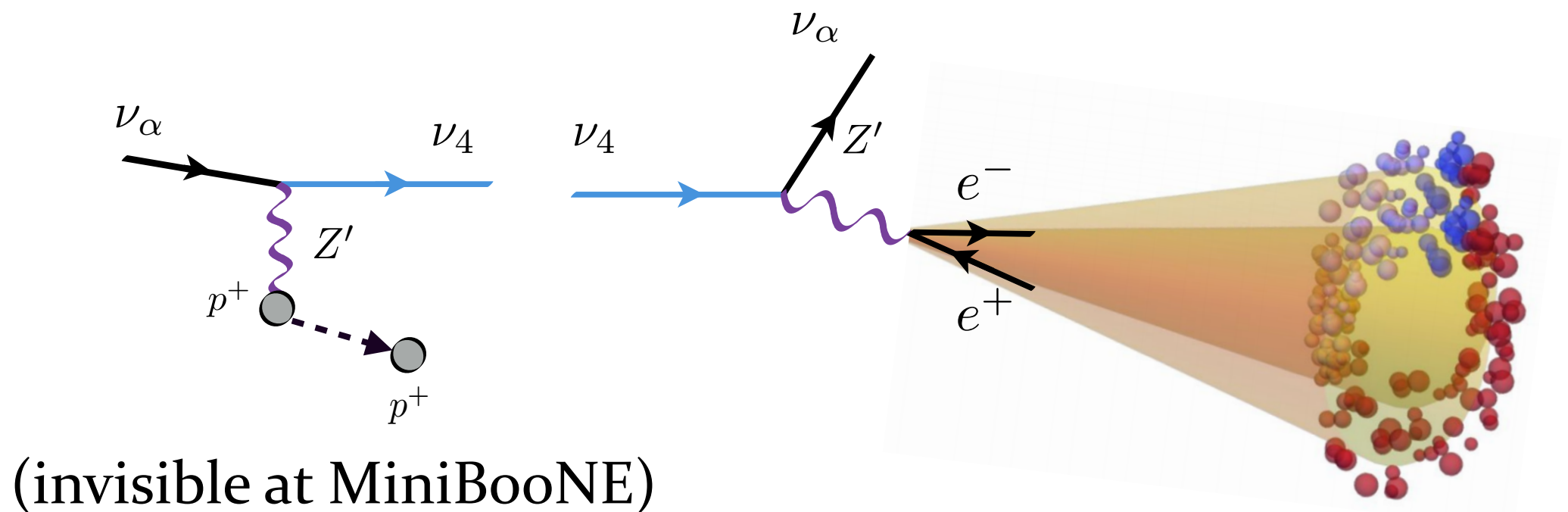
Heavy dark photon case

P. Ballett et al. arXiv:1808.02915

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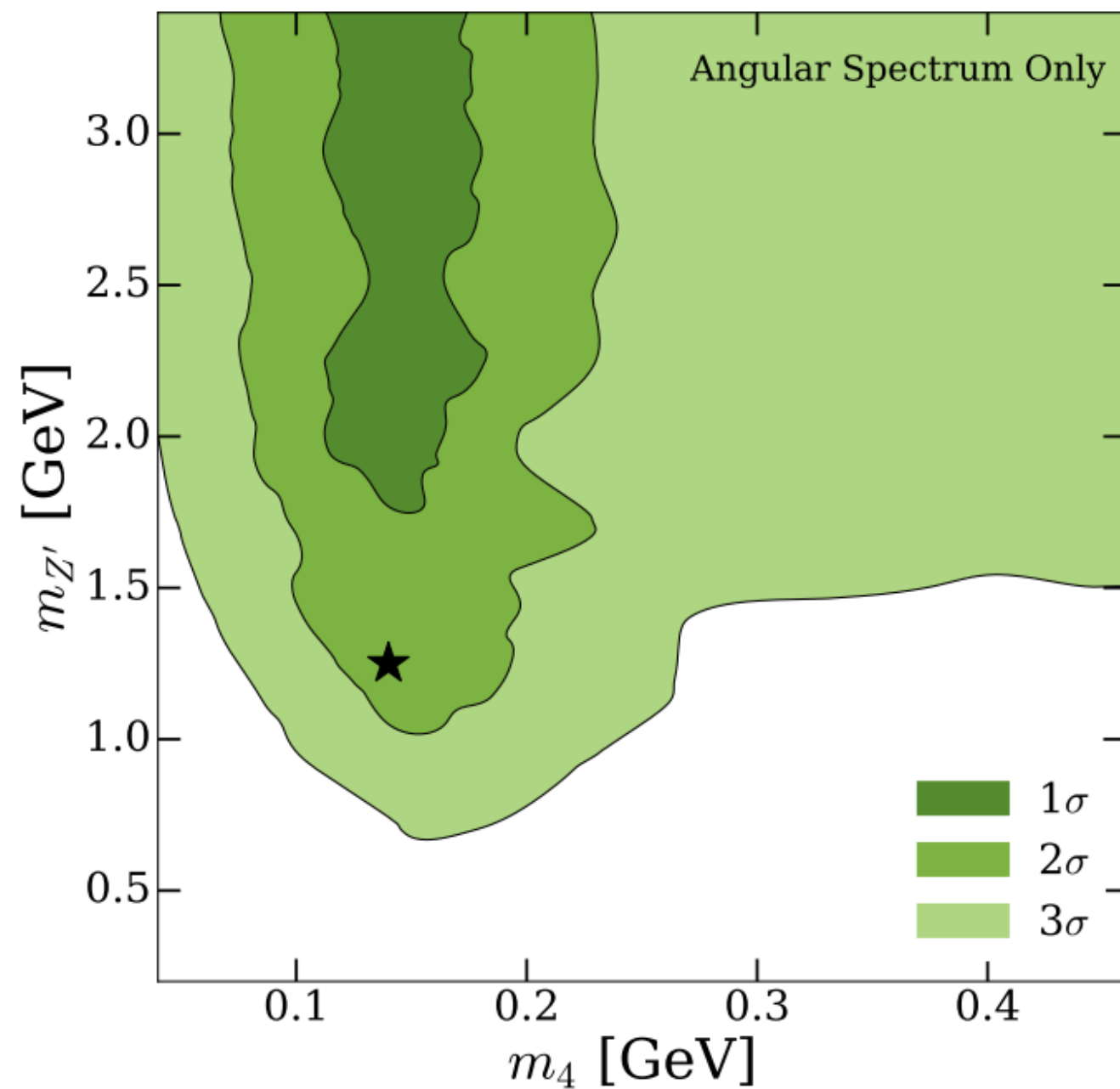
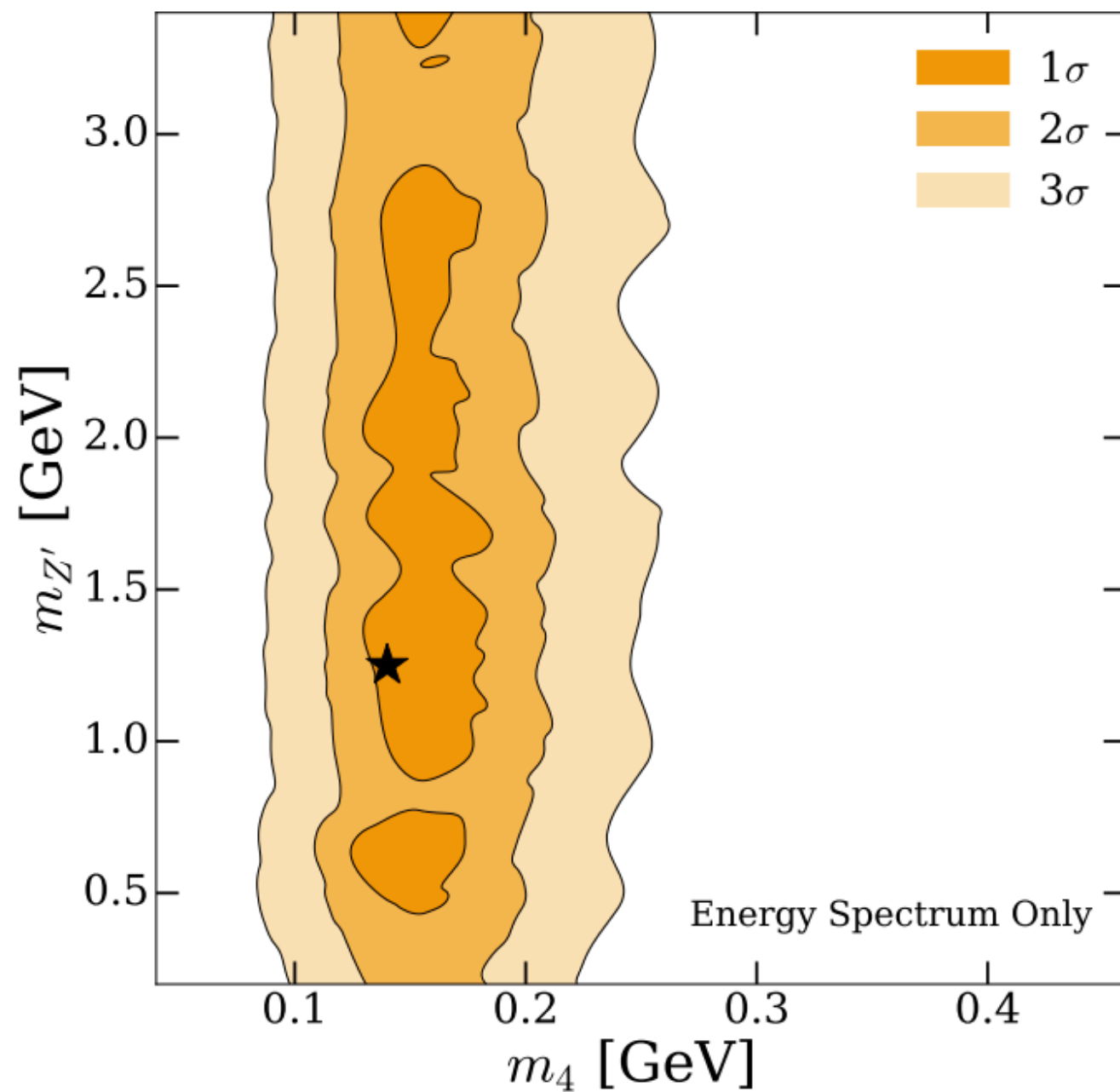
Incoherent signal

— single EM shower + some hadronic activity —

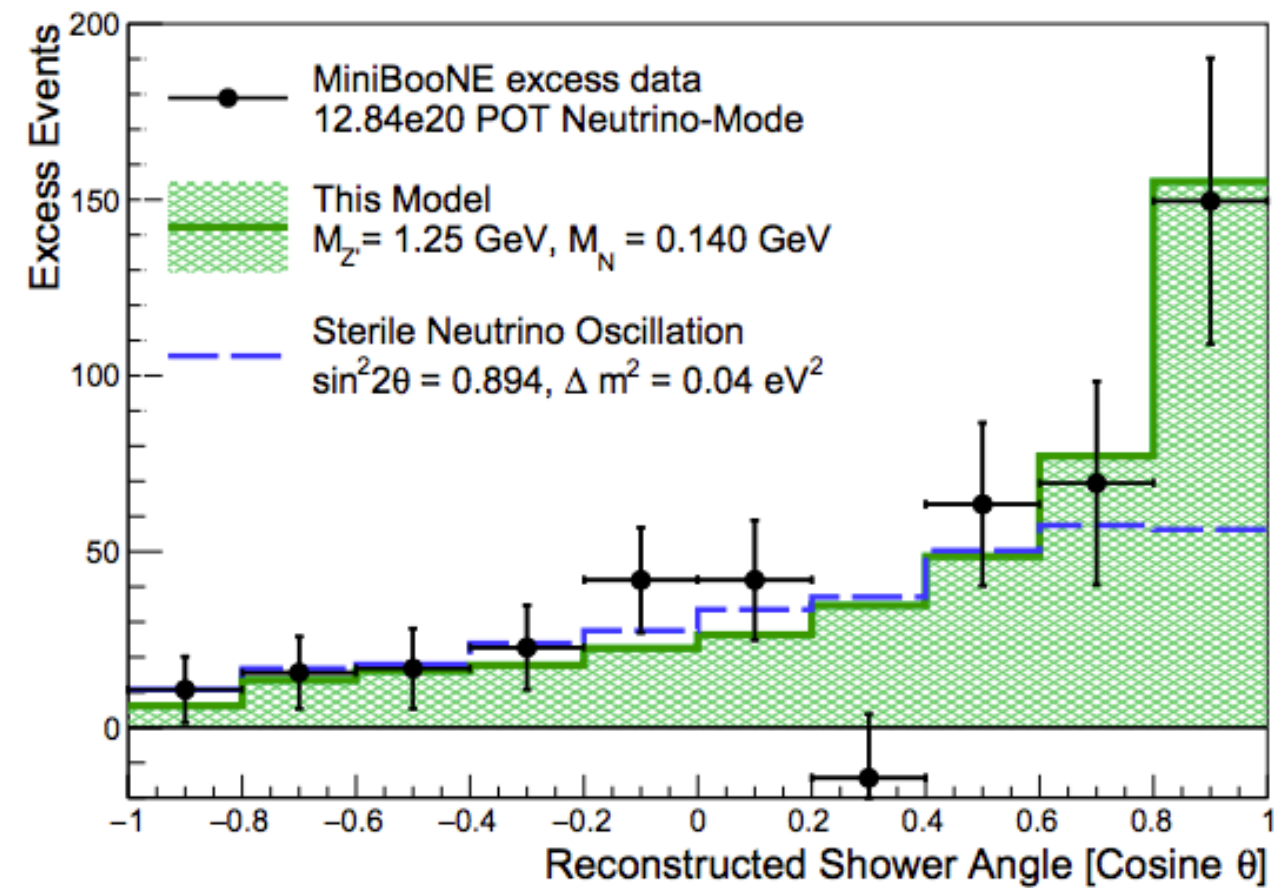
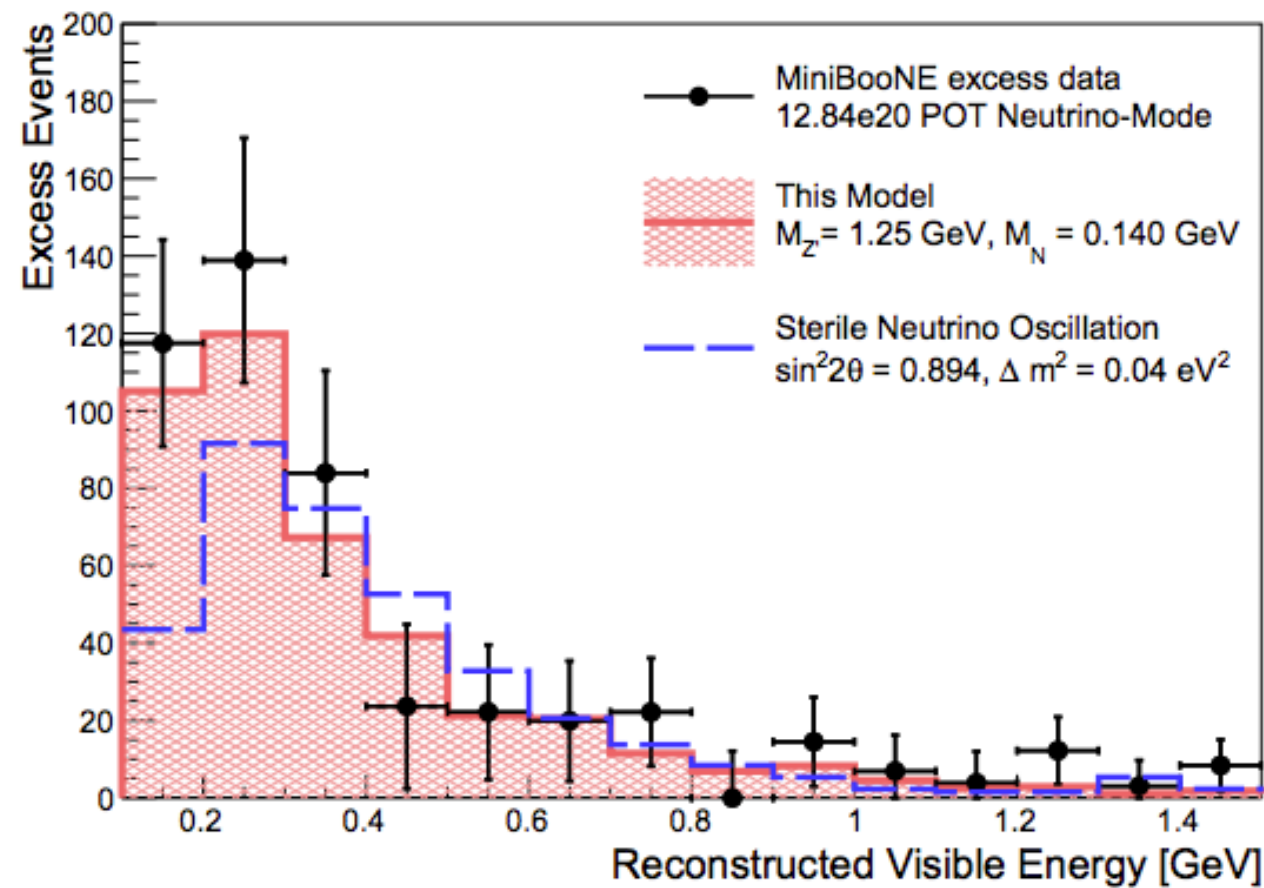


Heavy dark photon case

Shape-only fit



Heavy dark photon case



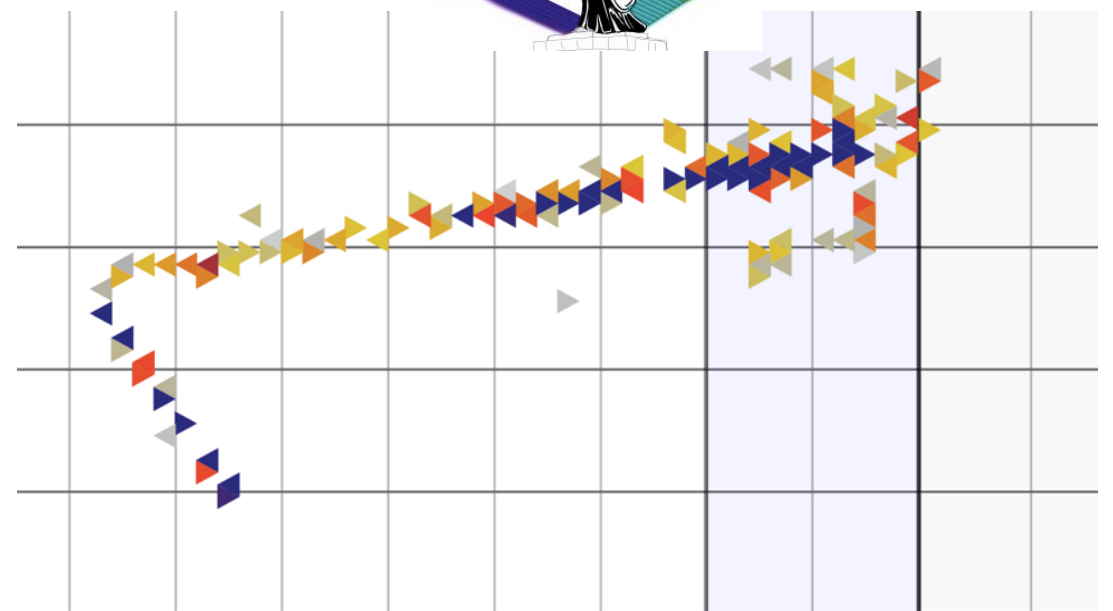
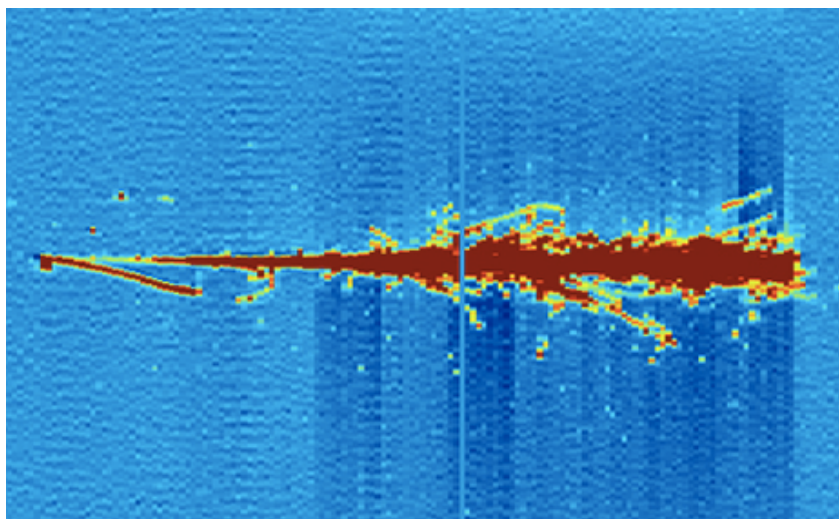
Much better angular fit to MiniBooNE data throughout parameter space.

Prospects for this case?

Harder to constrain at neutrino experiments, but could try!

proton + EM shower (+ gap?)

μ BooNE 



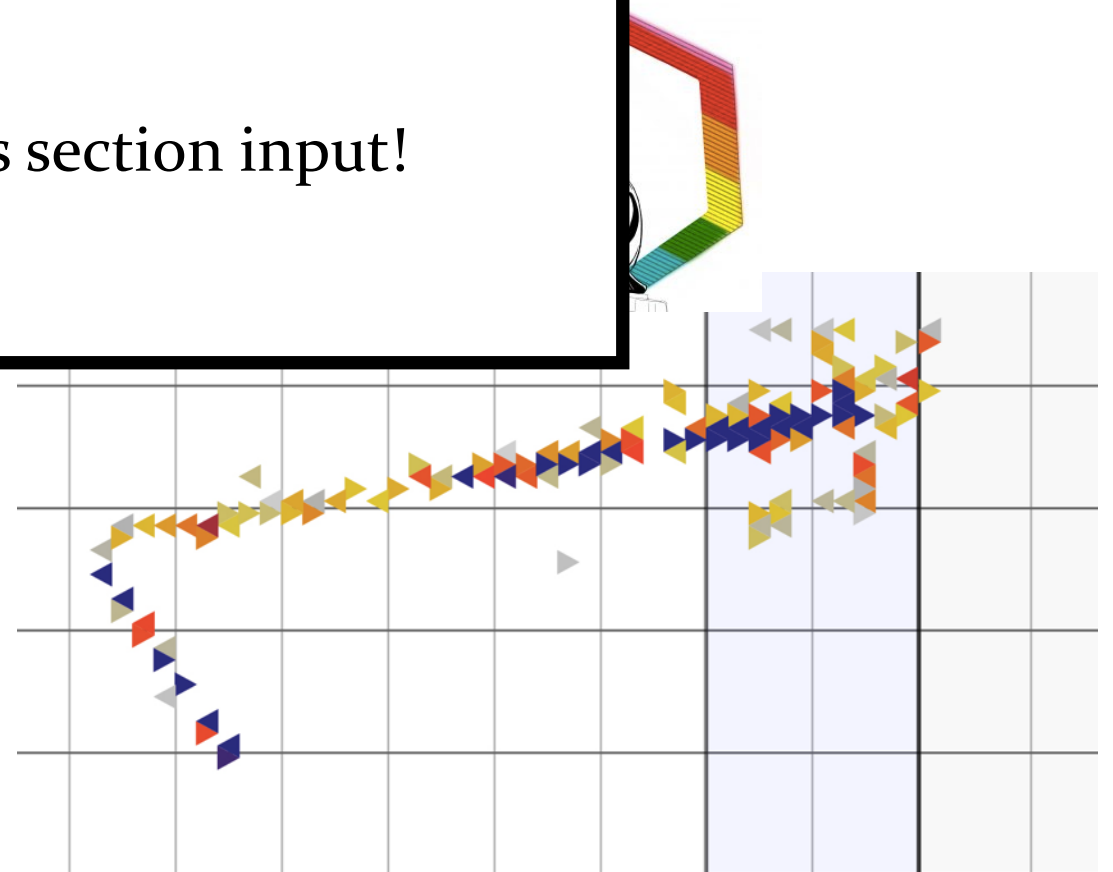
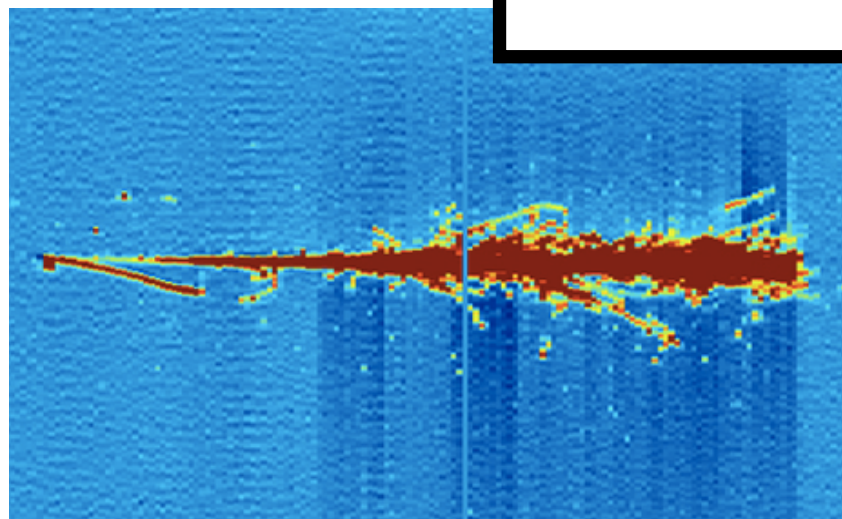
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μ BooN

Need neutrino cross section input!



Conclusions

Dark neutrinos are a class of models with **novel interactions** connected to the neutrino sector.

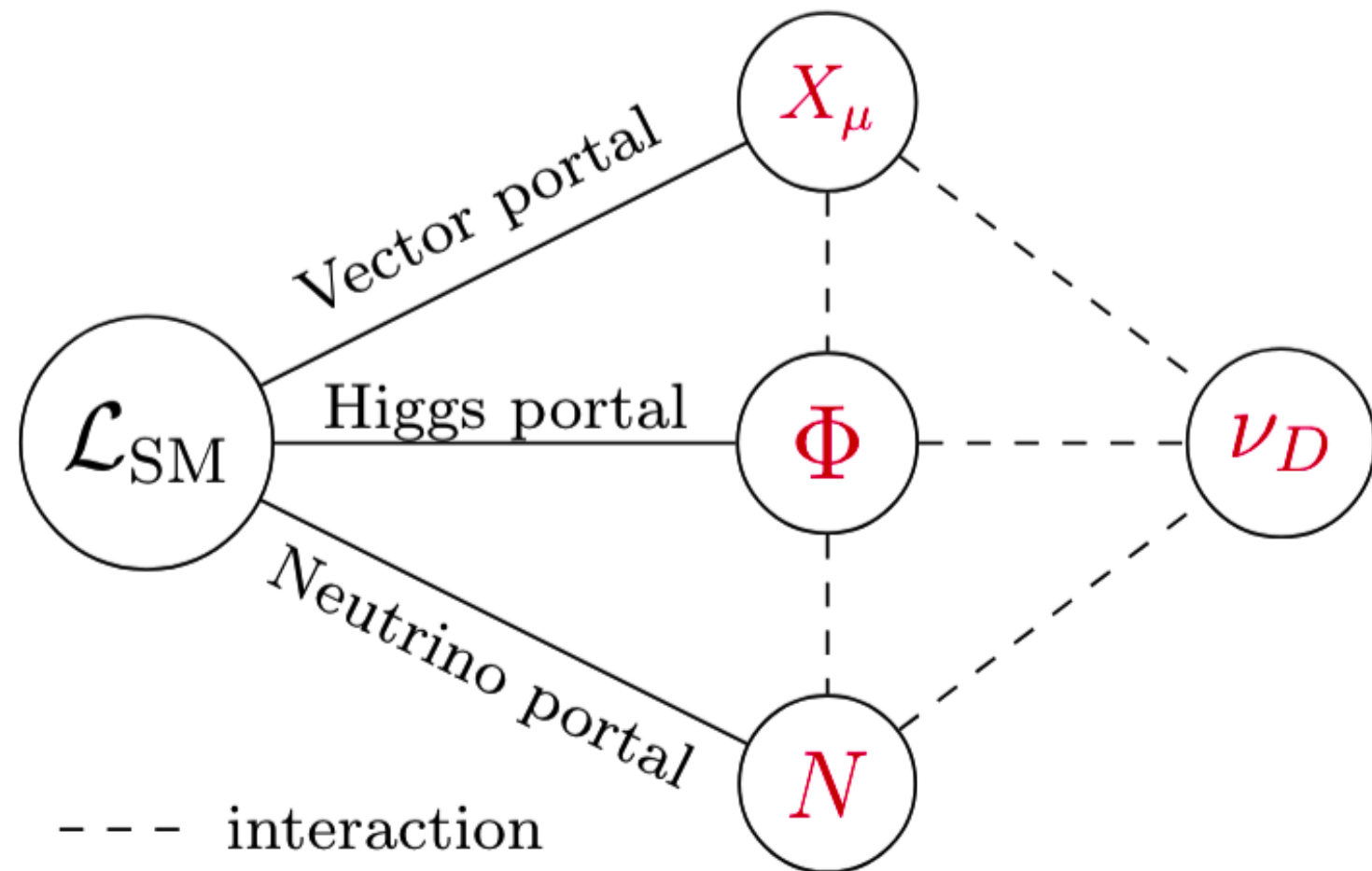
Can also generate **neutrino masses** at one-loop level — novel mechanism

Dark photon and scalar searches do not apply at face value
— more work needed to understand how **open** the **parameter space** is—

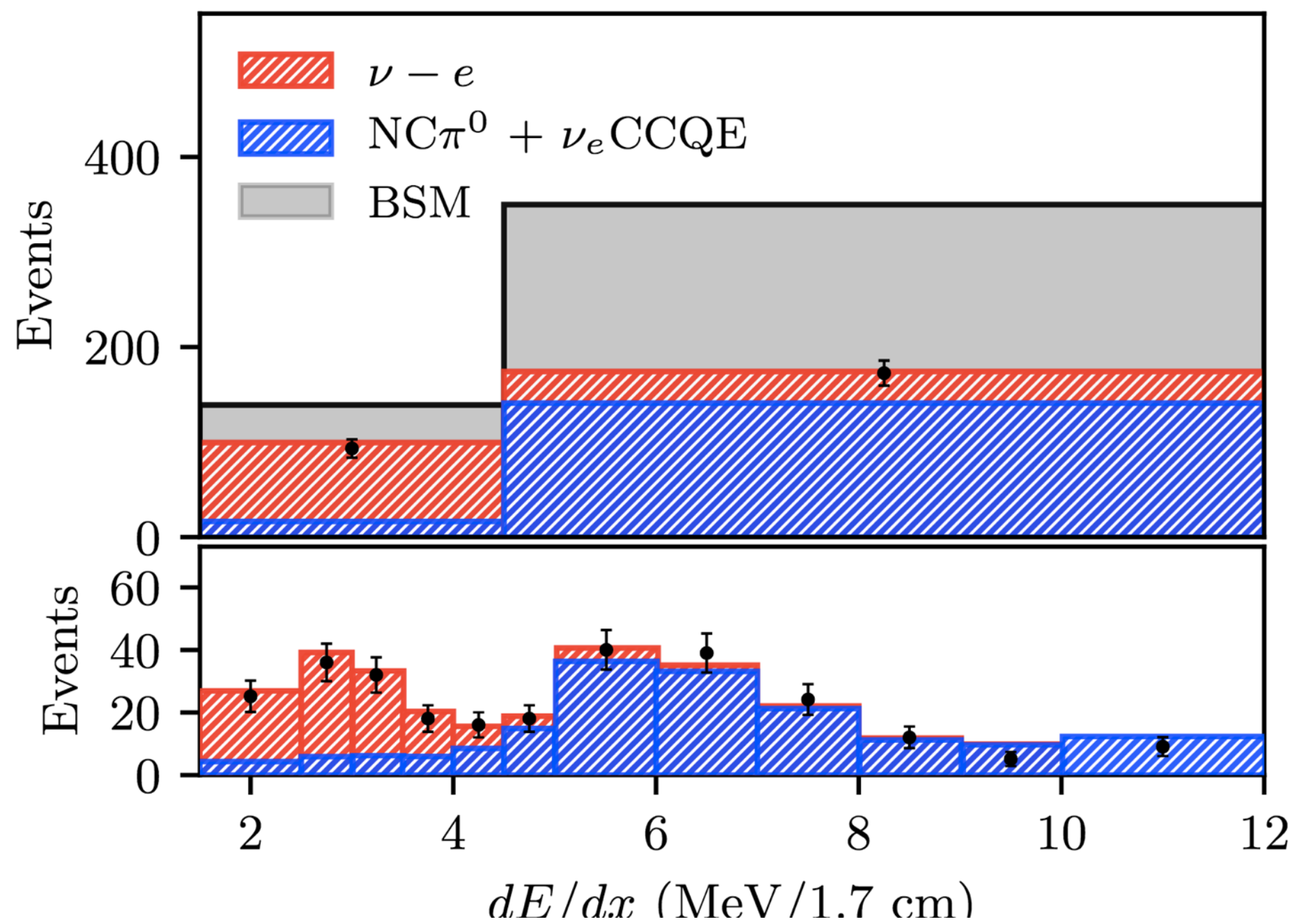
Predicts the **MiniBooNE** result in two different realisations.

Angular information and nu-e scattering data points us
to **heavy dark photon masses (~ 1 GeV)** boson mass.

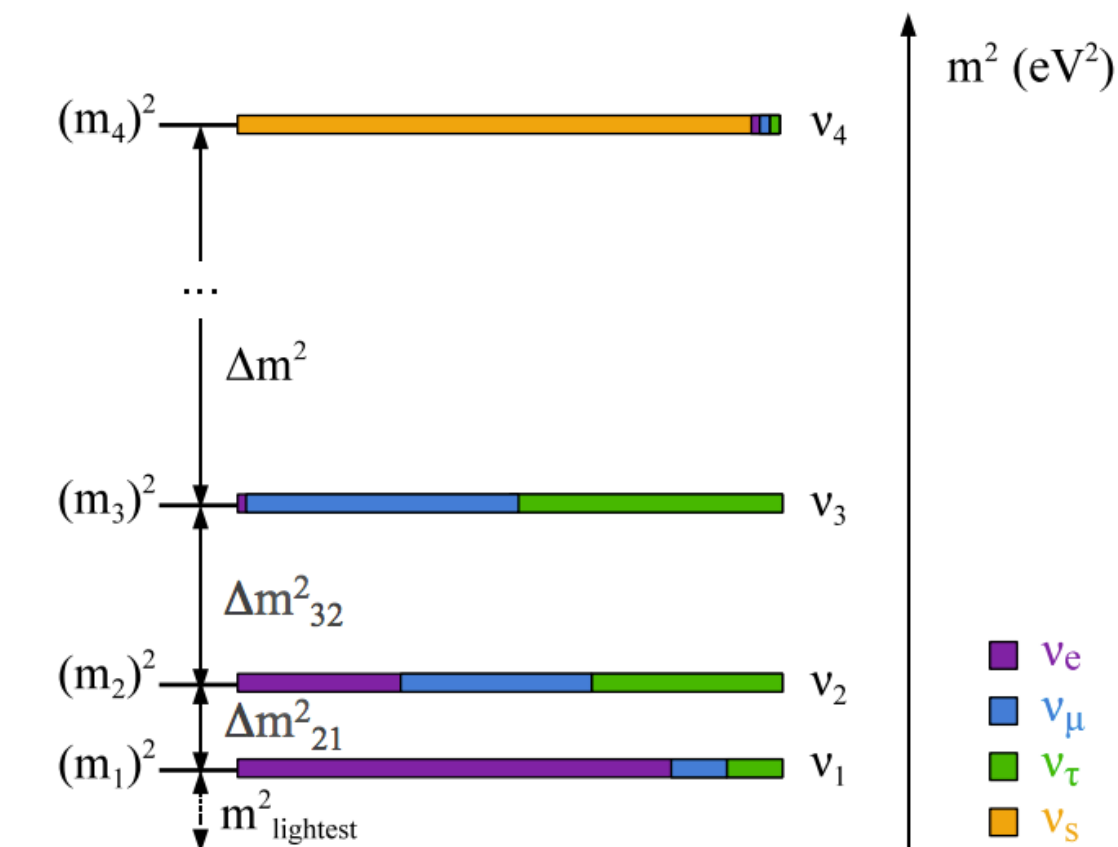
THANK YOU



APPENDIX



eV sterile neutrinos



Sterile neutrino with eV mass

$$P_{\nu_e \rightarrow \nu_\mu}^{3+1} = 4|U_{e4}|^2|U_{\mu4}|^2 \sin^2 \left(\frac{\Delta m^2}{4E} \right)$$

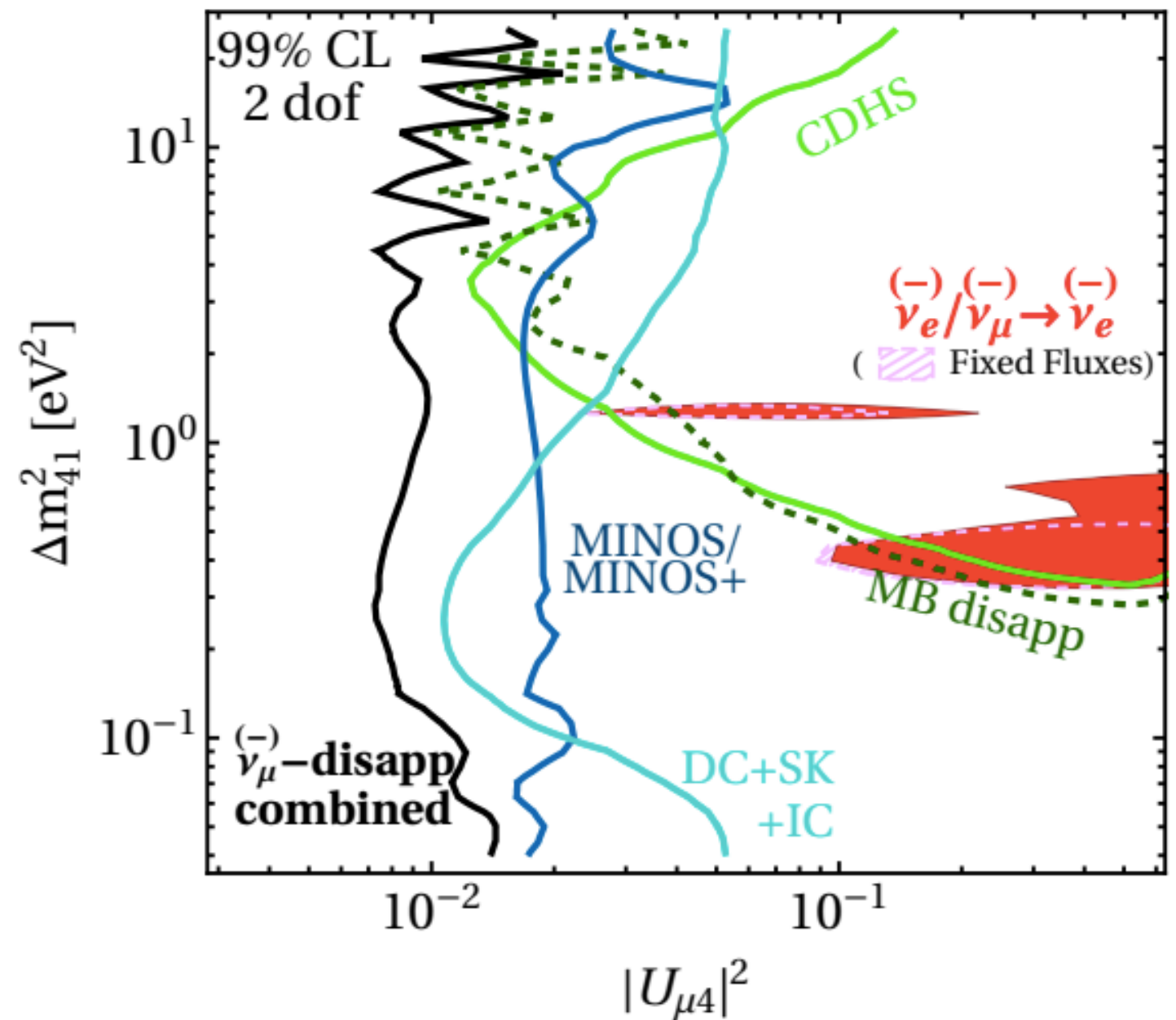
$$P_{\nu_\mu \rightarrow \nu_\mu}^{3+1} = 1 - 4|U_{\mu4}|^2(1 - |U_{\mu4}|^2) \sin^2 \left(\frac{\Delta m^2}{4E} \right)$$

Need electron and muon disappearance

for an appearance signal.

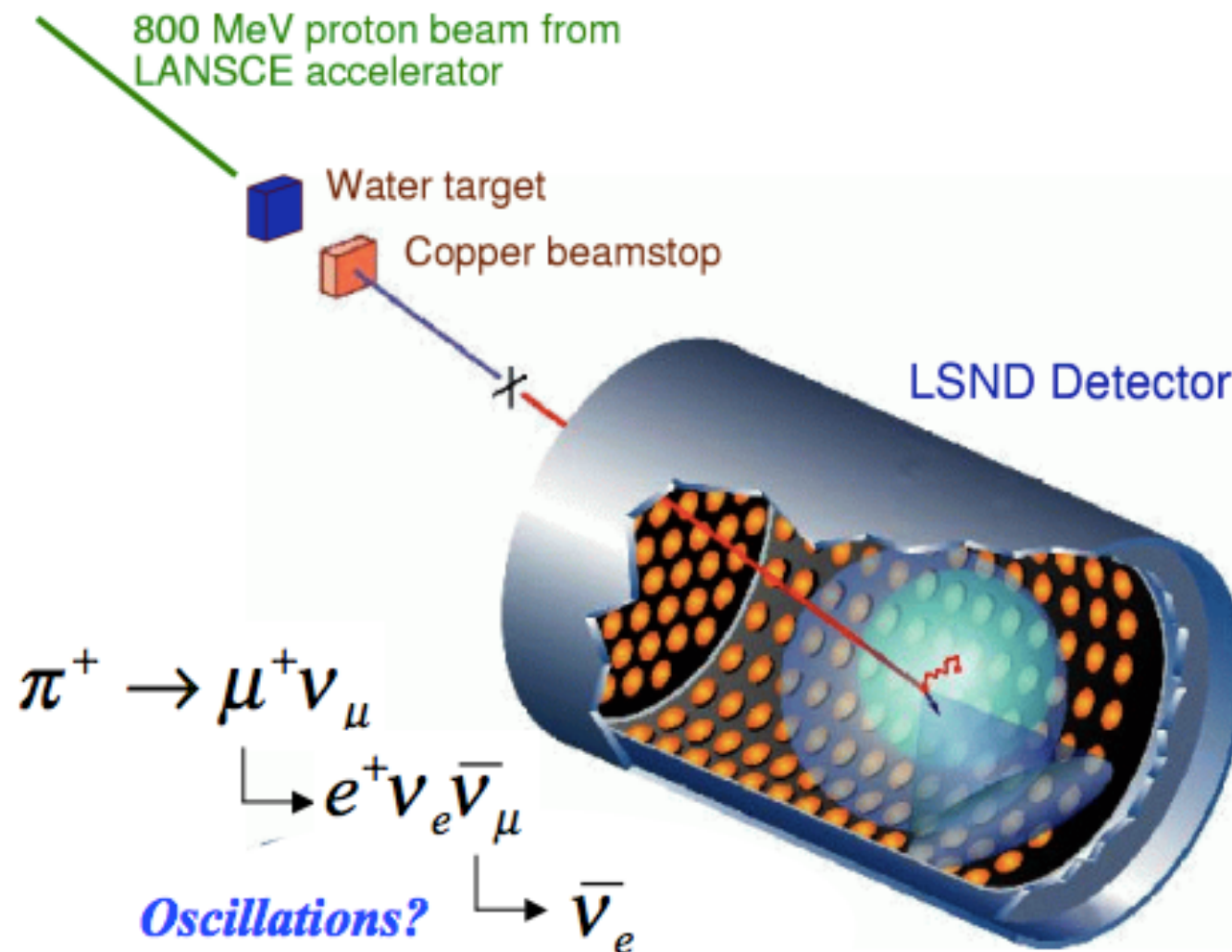
eV sterile neutrinos

Standard eV sterile
explanation in large tension
with disappearance
experiments.

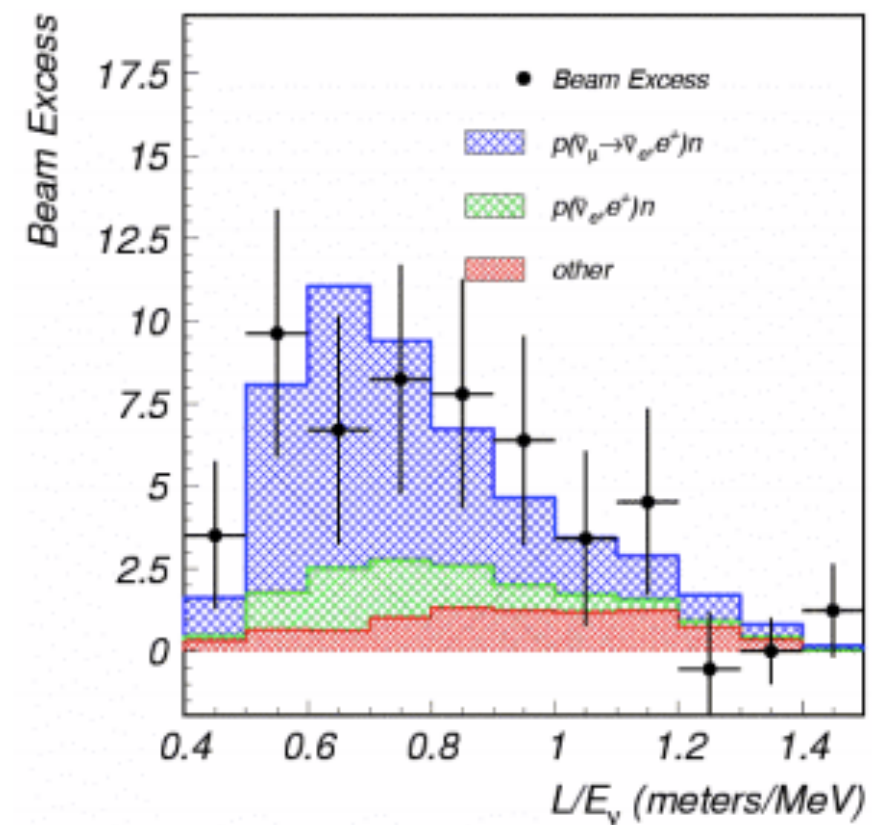


See Dentler et al, 1803.10661

LSND



LSND in conjunction with the atmospheric and solar oscillation results needs more than 3 ν 's



Saw an excess of:
 $87.9 \pm 22.4 \pm 6.0$ events.

With an oscillation probability of
 $(0.264 \pm 0.067 \pm 0.045)\%$.

3.8 σ evidence for oscillation.

A closer look

Series of published papers by HARP~CDP collaboration.

Anomaly goes from 3.8σ to 2.3σ .

Revisiting the 'LSND anomaly' I: impact of new data

The HARP~CDP Group: A. Bolshakova, I. Boyko, G. Chelkov, D. Dedovitch, A.

Revisiting the 'LSND anomaly' II: critique of the data analysis

Corrections to the HARP~CDP Analysis of the LSND Neutrino Oscillation Backgrounds

G. T. Garvey, W. C. Louis, G. B. Mills, D. H. White

Unpublished 1 page paper as reply.

Conclusion claimed to hold, but discussions end here because...

Reply to 'Corrections to the HARP~CDP Analysis of the LSND Neutrino Oscillation Backgrounds'

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(Submitted on 16 Dec 2011)