Searching for Quantum Gravity with neutrinos, Optical Module Beam Test at Fermilab and Hadronization Model studies for IceCube

IceCube Lab (ICL) Sven Lidstrom, NSF

Find us on Facebook! "Institute of Physics Astroparticle Physics" https://www.facebook.com/IOPAPP

Shivesh Mandalia

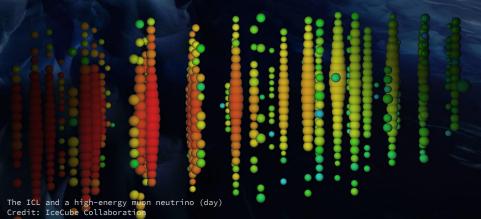
Queen Mary University of London



S. Mandalia QMUL PPRC Seminar - 2019-10-17



Outline



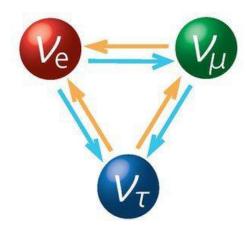
- 1. Neutrino Standard Model
- 2. New Physics with Astrophysical Neutrinos
- 3. IceCube Neutrino Observatory
- 4. Search for New Physics with the Astrophysical Flavour Ratio
- 5. IceCube DOM beam test at the FTBF
- 6. Hadronization modelling in neutrino interactions

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Neutrino Standard Model (vSM)

Neutrinos have nonzero mass!

Discovered via **neutrino oscillations** \rightarrow





The Nobel Prize in Physics 2015 Takaaki Kajita, Arthur B. McDonald

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The Nobel Prize in Physics 2015



Photo © Takaaki Kajita Takaaki Kajita Prize share: 1/2



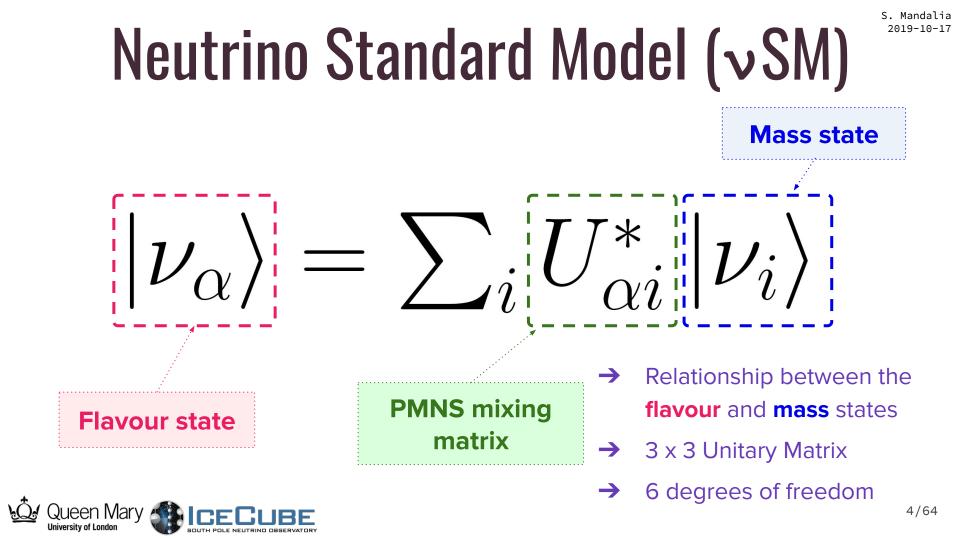
Photo: K. McFarlane. **Oueen's University** SNOLAR

Arthur B. McDonald Prize share: 1/2

The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald "for the discovery of neutrino oscillations, which shows that neutrinos have mass"



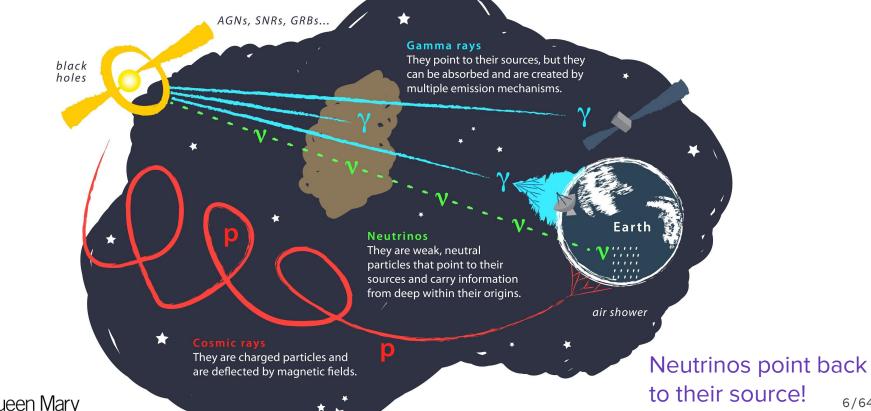
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New Physics with Astrophysical Neutrinos



Astrophysical Neutrinos

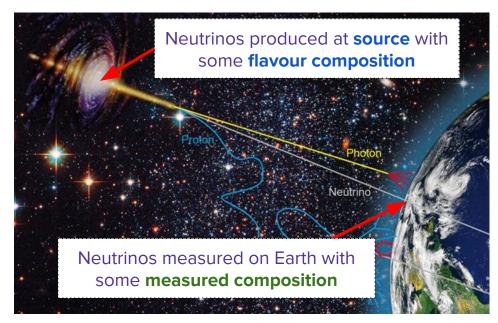


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Astrophysical Neutrino Flavour



$$P_{\nu_{\alpha} \to \nu_{\beta}} = \sum_{i} \mid U_{\alpha i} \mid^{2} \mid U_{\beta i} \mid^{2}$$

Example **source** flavour ratio scenarios:

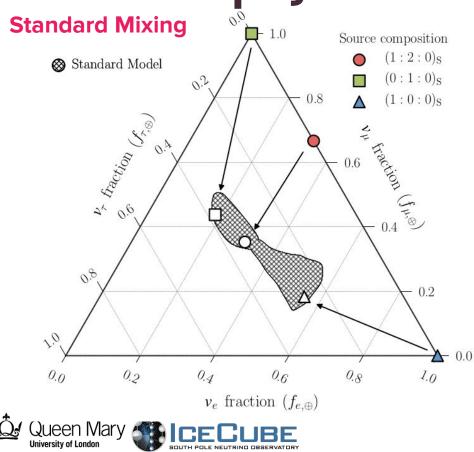
$$(f_e:f_\mu:f_\tau)_{\rm S}$$

- \rightarrow (1:2:0)_s charged pion-decay
- \rightarrow (1:0:0)_s neutron decay dominant
- → $(0:1:0)_{s}$ rapid muon energy loss



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Astrophysical Neutrino Flavour



$$P_{\nu_{\alpha} \to \nu_{\beta}} = \sum_{i} |U_{\alpha i}|^{2} |U_{\beta i}|^{2}$$

Example **source** flavour ratio scenarios:

$$(f_e:f_\mu:f_\tau)_{\rm S}$$

- \rightarrow (1:2:0)_s charged pion-decay
- \rightarrow (1:0:0)_s neutron decay dominant
- \rightarrow (0:1:0)_s rapid muon energy loss
 - $(1:2:0)_{\rm S} \rightarrow (0.31:0.35:0.34)_{\oplus}$
 - $(0:1:0)_{\rm S} \rightarrow (0.18:0.44:0.38)_{\oplus}$
 - $(1:0:0)_{\rm S} \rightarrow (0.55:0.18:0.27)_{\oplus}$

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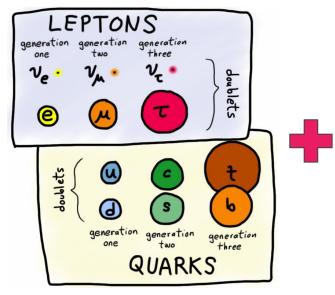
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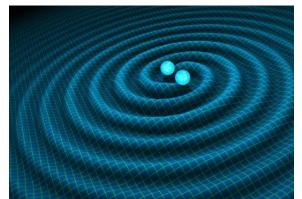
Motivation

Unified theories of the standard model of particle physics and general relativity allow for Lorentz symmetry violation

- → String theory^[1]
- → Quantum gravity^[2]

→ etc.

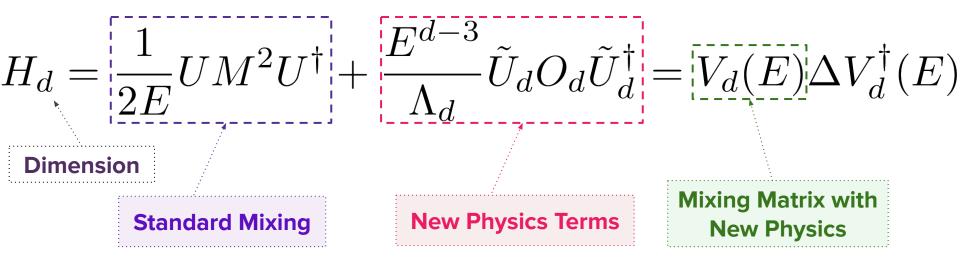






New Physics Effective Hamiltonian

Introduce new physics in the mixing matrix



→ Motivated by Standard Model Extension (SME), which is a general effective field theory framework to look for Lorentz violation

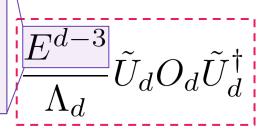


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New Physics Effective Hamiltonian



- → Power of astrophysical v comes from the energy dependence
- → Can provide strong constraints for higher dimensional operators

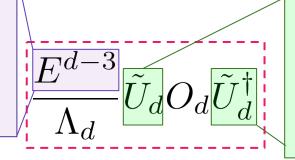




New Physics Effective Hamiltonian



- → Power of astrophysical v comes from the energy dependence
- → Can provide strong constraints for higher dimensional operators



BSM mixing elements

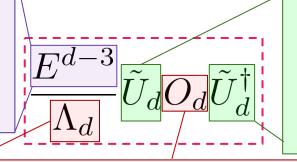
- → 3x3 Unitary Matrix
- → Fix some elements to obtain best case limits
 - Same spirit as LV community



S. Mandalia 2019-10-17 **New Physics Effective Hamiltonian**



- \rightarrow **Power** of astrophysical ν comes from the energy dependence
- Can provide strong constraints for \rightarrow higher dimensional operators

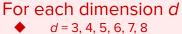


BSM mixing elements

- \rightarrow 3x3 Unitary Matrix
- \rightarrow Fix some elements to obtain best case limits
 - Same spirit as LV community

We will provide limits on **Scale of New Physics** the scale of new physics This is what we are

setting limits on!!





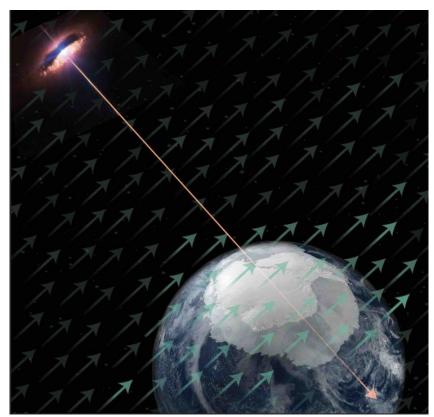
[1] Klop, Ando, Phys. Rev. D 97, 063006 (2018)

New Physics Interpretation

New operators can be interpreted in different new physics contexts

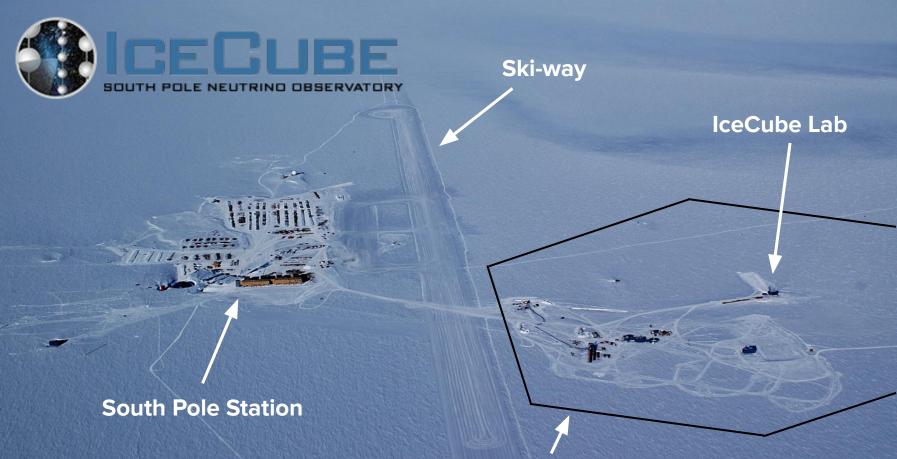
- → Lorentz and CPT Violation
- → Dark Energy Interaction ^[1]

 \rightarrow etc.

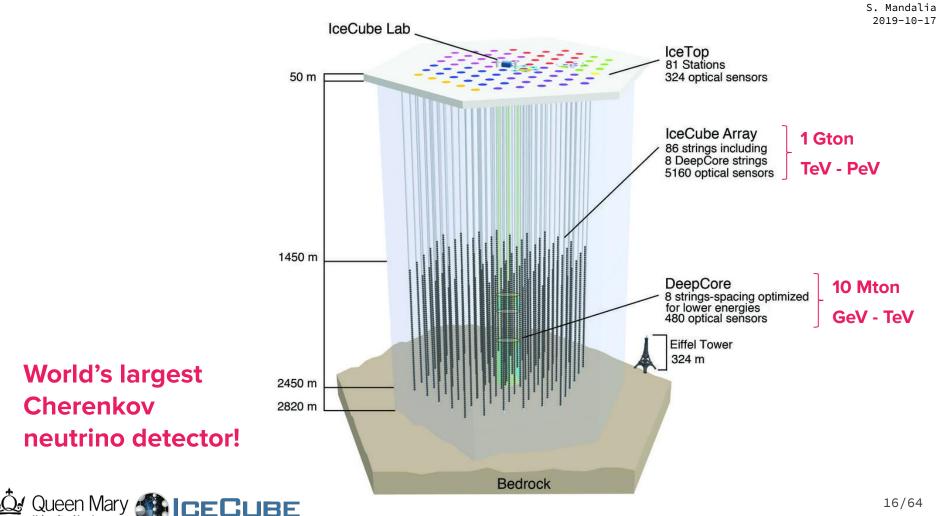




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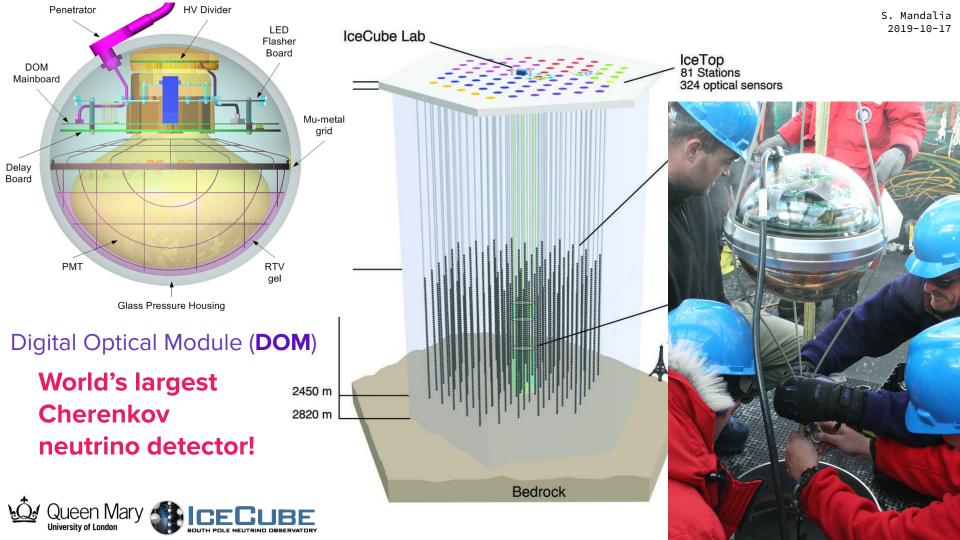


IceCube outline

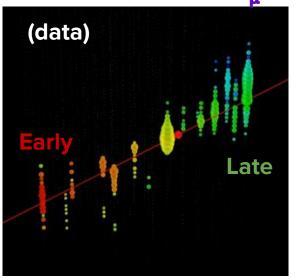


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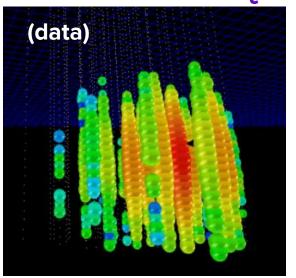
OUTH POLE NEUTRING OBSERVATORY



Charged-current v



Neutral-current / ve



Track

 $\nu_{\mu} + N \rightarrow \mu + X$

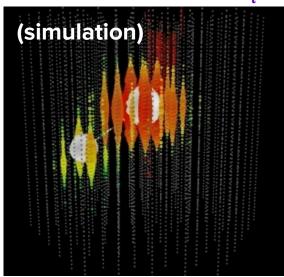
Angular resolution ~ 0.5° Energy resolution ~ factor 2



Cascade $v_e + N \rightarrow e + X$ $v_x + N \rightarrow v_x + X$

Angular resolution ~ 10° Energy resolution ~ 15%

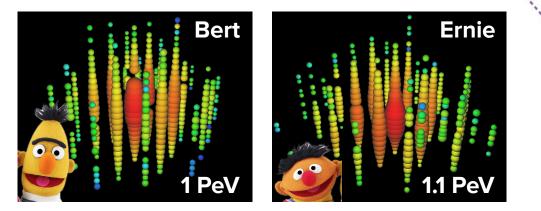
Charged-current v_{τ}



- **Double-Cascade** $\nu_{\tau} + N \rightarrow \tau + X$
- ~ 2 expected in 6 years

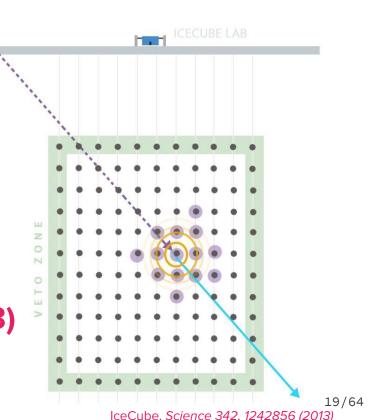
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High Energy Starting Events (HESE)



First Evidence for High-Energy Astrophysical Neutrinos! (Science 2013)



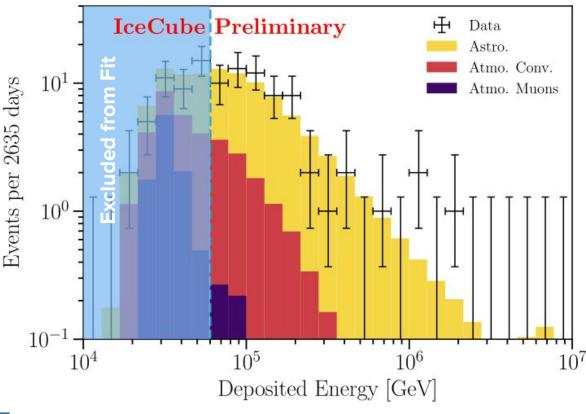


New 7.5 years HESE data!

All energies: 103 events

22 new events in 20169 new events in 2017

> 60 TeV: 60 events 42 cascades 15 tracks 2 double-cascade





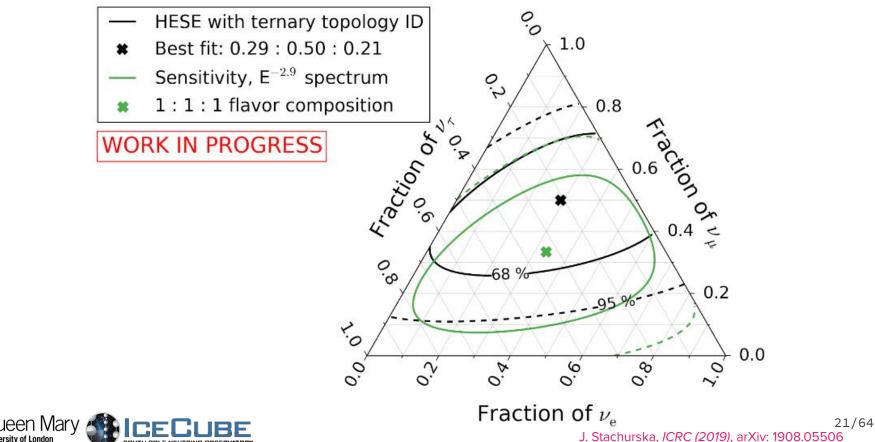
A. Schneider, ICRC (2019), arXiv: 1907.11266

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Flavour Ratio Measurement

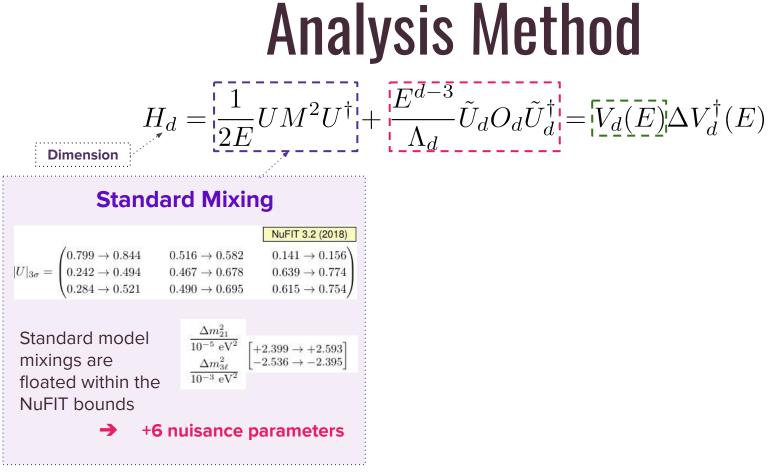
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JHEP 01 (2017) 087 [arXiv:1611.01514] - NuFIT 3.0 (2016), www.nu-fit.org

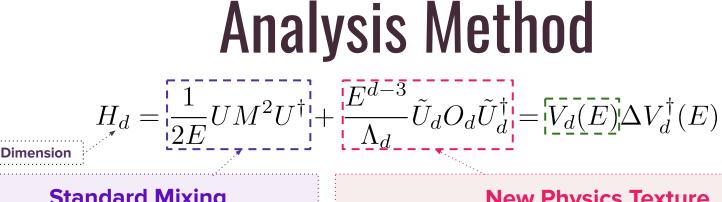
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JHEP 01 (2017) 087 [arXiv:1611.01514] - NuFIT 3.0 (2016), www.nu-fit.org

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Standard Mixing

			NuFIT 3.2 (2018)
	$(0.799 \rightarrow 0.844)$	$0.516 \rightarrow 0.582$	$0.141 \rightarrow 0.156$
$ U _{3\sigma} =$	$0.242 \rightarrow 0.494$	$0.467 \rightarrow 0.678$	$0.639 \rightarrow 0.774$
	$0.284 \rightarrow 0.521$	$0.490 \rightarrow 0.695$	$0.615 \rightarrow 0.754$

Standard model mixings are floated within the NuFIT bounds

$$\begin{array}{c} \frac{\Delta m^2_{21}}{10^{-5}~{\rm eV}^2} \\ \frac{\Delta m^2_{3\ell}}{10^{-3}~{\rm eV}^2} \end{array} \begin{bmatrix} +2.399 \rightarrow +2.593 \\ -2.536 \rightarrow -2.395 \end{bmatrix}$$

+6 nuisance parameters



New Physics Texture

New Physics Matrix is fixed to a **maximally oscillating** texture

Set all but one of the New Physics mixing angles to 0

$$\tilde{\theta}_{ij,d} = \begin{cases} \pi/4 &\in \tilde{\theta}_{13,d} \\ 0 &\text{otherwise} \end{cases} \implies \tilde{U}_d = \mathcal{O}_{e\tau} \equiv \begin{pmatrix} \frac{1}{\sqrt{2}} & 0 & \frac{1}{\sqrt{2}} \\ 0 & 1 & 0 \\ -\frac{1}{\sqrt{2}} & 0 & \frac{1}{\sqrt{2}} \end{pmatrix} + 0 i$$
$$\tilde{\theta}_{ij,d} = \begin{cases} \pi/4 &\in \tilde{\theta}_{23,d} \\ 0 &\text{otherwise} \end{cases} \implies \tilde{U}_d = \mathcal{O}_{\mu\tau} \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\ 0 & -\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix} + 0 i$$

Fit for the New Physics Scale Λ_d^{-1}

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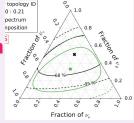
Analysis Method

$\left[\overline{\phi}_{\beta, \oplus}^{(d)} \right] = \frac{1}{\left| \Delta E \right|} \int_{\Delta E} \sum_{\alpha} P_{\nu_{\alpha} \to \nu_{\beta}}^{(d)}(E) \phi_{\alpha, S}(E) \, \mathrm{d}E$

Measured Flavour

$$f^{(d)}_{\beta,\oplus} = \bar{\phi}^{(d)}_{\beta,\oplus} / \sum\nolimits_{\gamma} \bar{\phi}^{(d)}_{\gamma,\oplus}$$

This analysis is an interpretation of this flavour contour





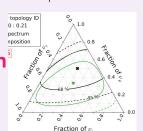
Analysis Method

Measured Flavour

 $P^{(a)}_{\nu_{\alpha} \to \nu_{\beta}}$

$$f^{(d)}_{\beta,\oplus} = \bar{\phi}^{(d)}_{\beta,\oplus} / \sum_{\gamma} \bar{\phi}^{(d)}_{\gamma,\oplus}$$

This analysis is an interpretation of this flavour contour



$$P_{\nu_{\alpha} \to \nu_{\beta}}^{(d)}(E) = \sum_{i} \left\| V_{\alpha i, d}(E) \right\|^{2} \left\| V_{\beta i, d}(E) \right\|^{2}$$

Mixing Matrix with New Physics

BSM introduced in the mixing during propagation



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 $(E) \phi_{\alpha, S}(E) dE$

E) dE

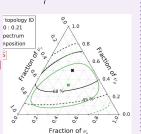
Analysis Method

Measured Flavour

$$f^{(d)}_{\beta,\oplus} = \bar{\phi}^{(d)}_{\beta,\oplus} / \sum\nolimits_{\gamma} \bar{\phi}^{(d)}_{\gamma,\oplus}$$

This analysis is an interpretation of this flavour contour

leen Marv



Mixing

→ Averaged over energy spectrum

$$P_{\nu_{\alpha} \to \nu_{\beta}}^{(d)}(E) = \sum_{i} \left\| V_{\alpha i, d}(E) \right\|^{2} \left\| V_{\beta i, d}(E) \right\|^{2}$$

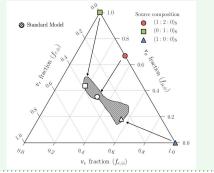
Mixing Matrix with New Physics

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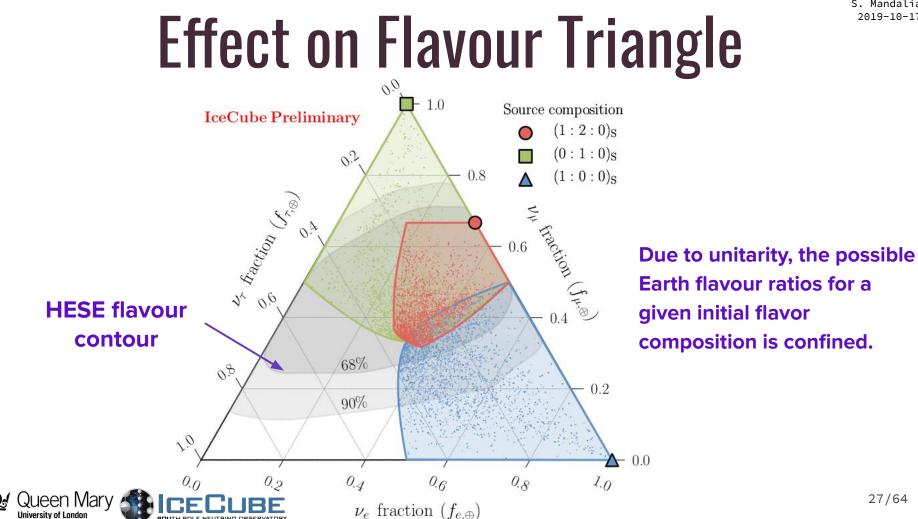
Source Flux

 $)\phi_{\alpha},$

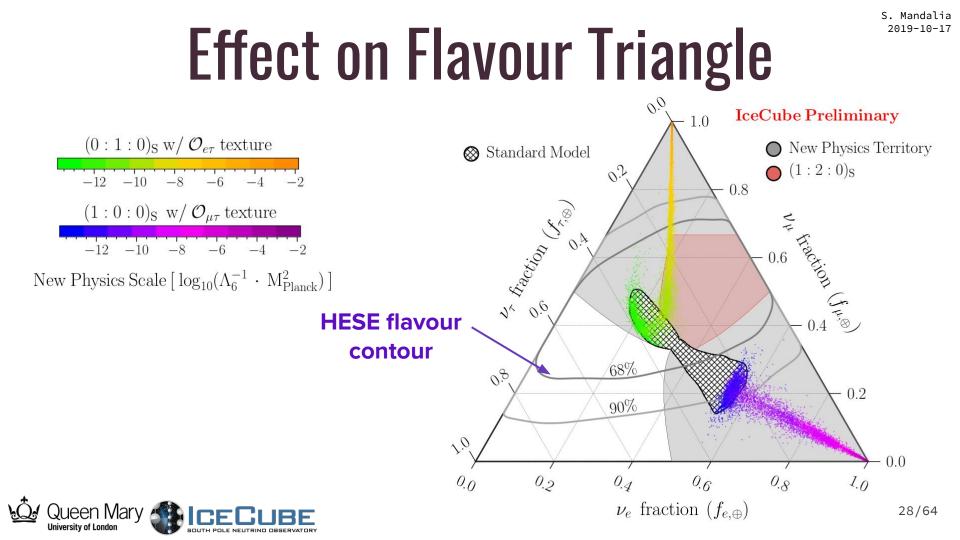
Fix to a certain model

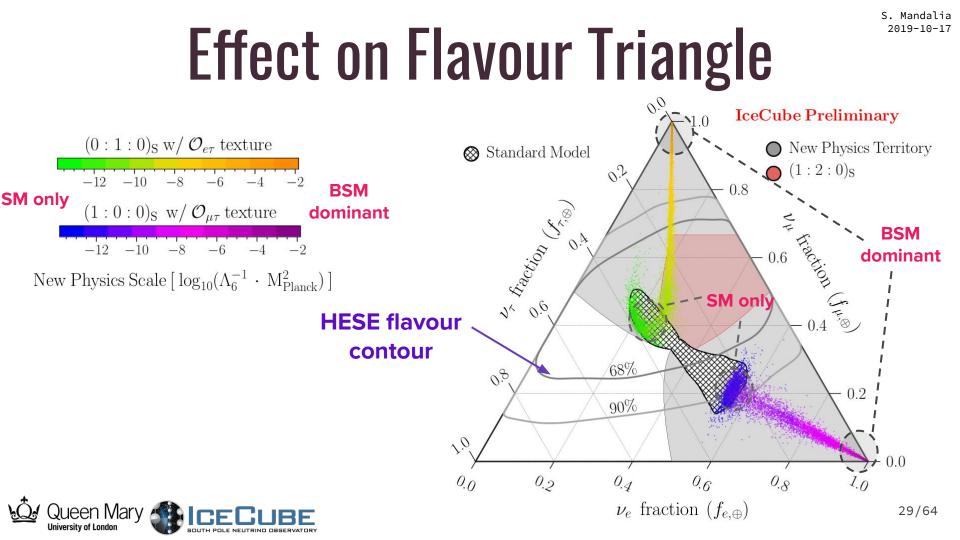






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Analysis Specification

B	Sinning	60 TeV < E < 10 PeV -1 < cos(zenith) < 1			
+ separate binning for double-cascade					
J. Stachurska, <i>JCRC (2019)</i> , arXiv: 1908.05506					

Normalization of flux:

- 1. conventional (40%)
- 2. prompt (±2.4 BERS)
- +5 nuisance

Sytematics

- astrophysical (free) params
- 4. muon background (50%)

Flux components:

3.

1. astrophysical neutrino index (free)

- Atmospheric neutrinos from Honda2006 Honda, Morihiro et al., Phys.Rev. D75 (2007) 043006
- → Prompt neutrinos from BERSS Bhattacharya, Enberg, Reno, Sarcevic, Stasto, JHEP 1506 (2015) 110
- → Simple power law for astrophysical neutrinos
- → DIS cross section from CSMS paper Cooper-Sarkar, Mertsch, Sarkar, JHEP 1108 (2011) 042
- → Analytical oscillation formula
- → Bayesian Markov Chain Monte Carlo (main results)
 - Bayes factor hypothesis test

Fit Methods

Simulation

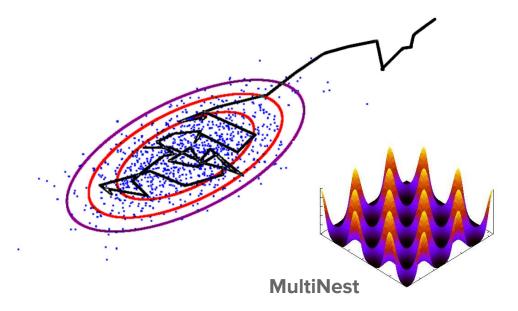
- 11 nuisance parameters in total
- + Λ_d^{-1} parameter of interest

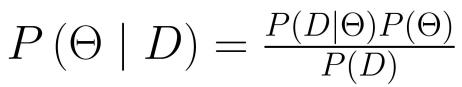
How do we work in such a high-dim space?



Markov Chain Monte Carlo

- Construct a random walk which explores the likelihood space
- Random walk is biased such that the distribution approximates the posterior space
- Great utility when following a **Bayesian** statistical treatment







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Anarchy Sampling

Need to be careful about which parameters we sample over

Could introduce artificial biases if not careful

$$U = e^{i\eta} e^{i\phi_1\lambda_3 + i\phi_2\lambda_8} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} e^{i\chi_1\lambda_3 + i\chi_2\lambda_8}$$

We sample in the space defined by the Haar Measure:

$$dU = d(\sin^2 \theta_{12}) \wedge d(\cos^4 \theta_{13}) \wedge d(\sin^2 \theta_{23}) \wedge d\delta$$
Sample uniformly in s_{12}^2 , c_{13}^4 , s_{23}^2 , δ
Volume Elements



A. de Gouvea and H. Murayama, Phys.Lett.B573, 94 (2003), hep-ph/0301050 L. J. Hall, H. Murayama, and N. Weiner, Phys.Rev.Lett.84, 2572(2000),hep-ph/9911341 32/64 N. Haba and H. Murayama, Phys.Rev. D63, 053010 (2001), hep-ph/0009174

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Constructing a Limit

If no New Physics is found, we can attempt to set a limit

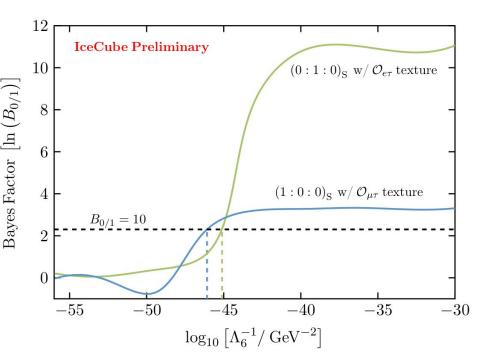
Define Bayes Factor

$$B_{0/1}(\mathbf{x}) = \frac{\mathcal{Z}_0(\mathbf{x})}{\mathcal{Z}_1(\mathbf{x})} \begin{bmatrix} z_0 \\ z_1 \end{bmatrix}$$
$$\mathcal{Z}_j(\mathbf{x}) = f(\mathbf{x}|\mathcal{M}_j) = \int f_j(\mathbf{x}|\boldsymbol{\theta}_j) \pi_j(\boldsymbol{\theta}_j) d\boldsymbol{\theta}_j \end{bmatrix}$$

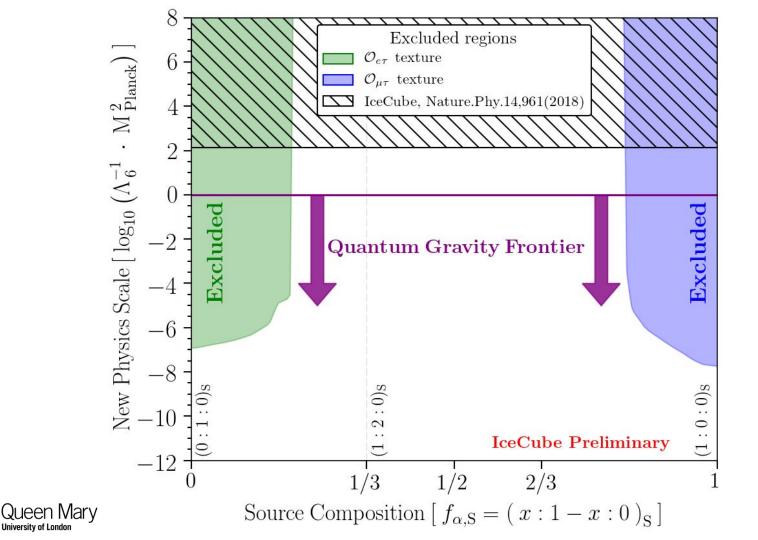
We define our threshold as having a "strong strength of evidence" according to Jeffreys' scale D > 10

$$B_{0/1} > 10$$





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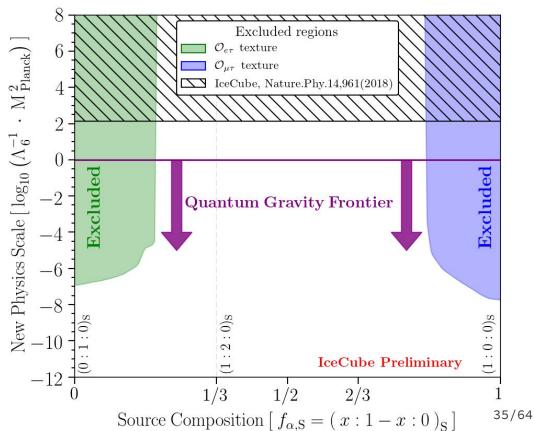


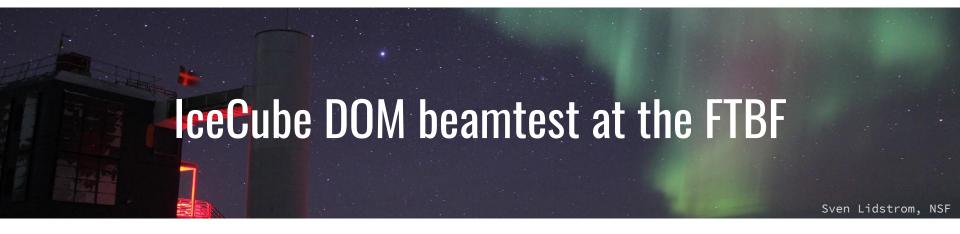
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Preliminary Results

- → Limits given per operator dimensionality
- → Limits given for fixed initial flavor composition (x-axis) and a BSM operator texture (in green and blue)
- → No constraints obtained if the initial flavor composition is pion-based (1:2:0)
- → First BSM physics constraints using astrophysical neutrino flavour information
- Constraints can be reinterpreted in terms of other BSM physics since they have been expressed as effective operators.

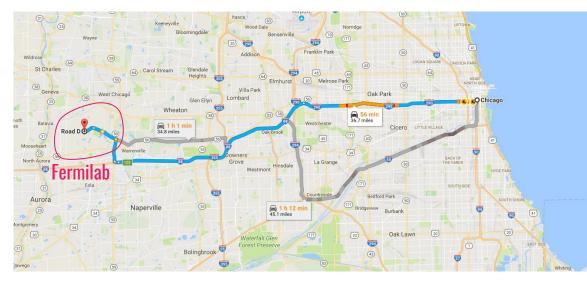






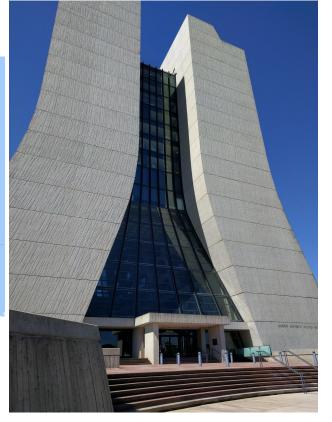


Fermilab



- → Discovery of the bottom/top quark
- → Observation of direct CP violation in kaon decays
- → Observation of tau neutrino (DONUT)
- → Headquarters for DUNE





Wilson Hall

Fermilab







Wilson Hall

Fermilab

One of the beams at Fermilab is available to users.

Fermilab Test Beam Facility (FTBF) From the website: "The FTBF program provides flexible, equal, and open access to test beams for all detector tests, with relatively low bureaucratic overhead and a guarantee of safety, coordination, and oversight."

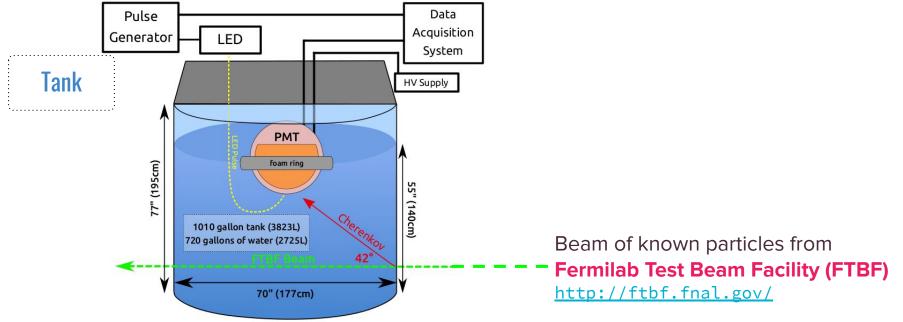


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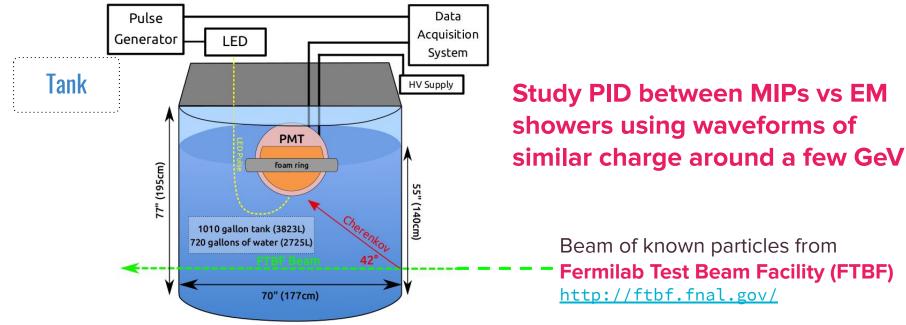
Idea



- → Black tedlar film coats the inside and outside
- → Filled with distilled water



Idea



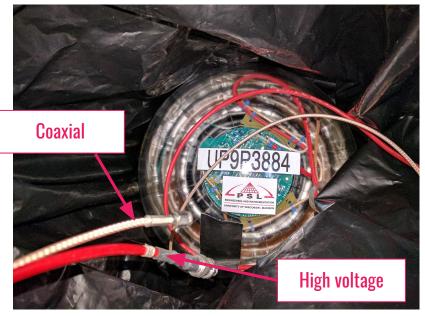
- → Black tedlar film coats the inside and outside
- → Filled with distilled water

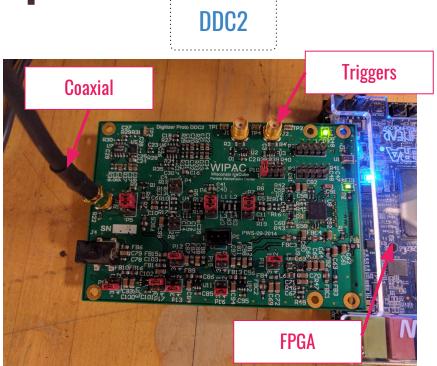






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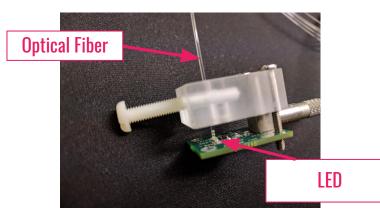




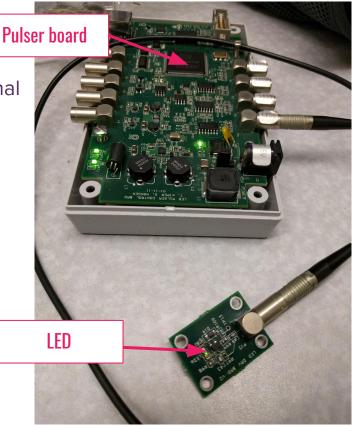
Setup

LED

- Use a well-defined LED pulses to test DOM is operational and produces the expected signal waveform
- → FPGA board allows dynamic control of the LED
- → We couple the LED light pulse to an optical fiber





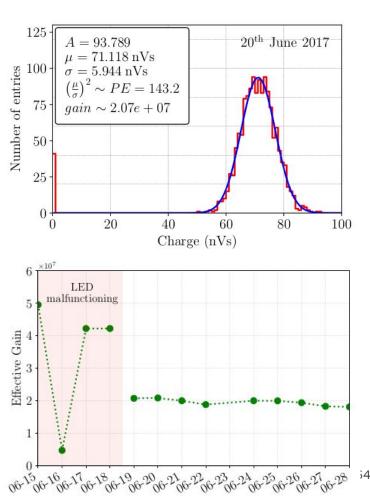


Setup

LED

- We can monitor the gain of the PMT by using the LED
 - Fit the charge distribution to a gaussian
 - Using poisson statistics, the number of PE hits and gain can be calculated
- This is useful to test the degradation of the water over the days when we have beam
 - Degradation of distilled water (i.e. changes in the water transparency) will affect the gain we measure on the PMT





Setup

Tank

- → Delivered to FTBF
- → Wrong colour! We were expecting a black tank
- → We coated the inside and outside layers of the tank in a black Tedlar film





Setup

Tank

- → Delivered to FTBF
- → Wrong colour! We were expecting a black tank
- \rightarrow We coated the inside and outside layers of the tank in a

black Tedlar film

Wires get fed through hole in the lid

Queen Marv 🌶

University of London





Setup



→ Moved the tank into the MT6.2 enclosure



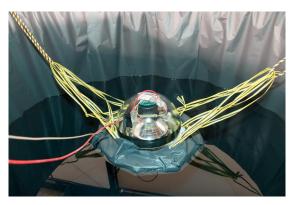


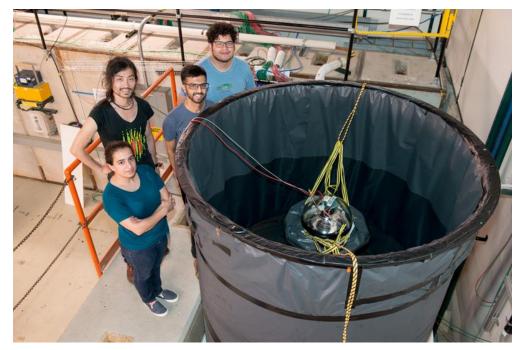
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Setup

Tank

- → Tank was filled with ~ 700 gallons of distilled water
- → DOM was placed at the centre of the tank using ropes





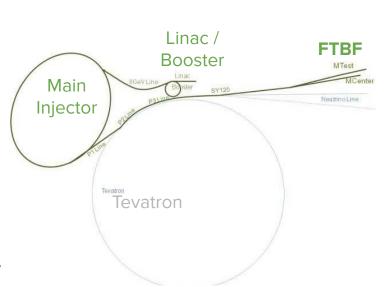


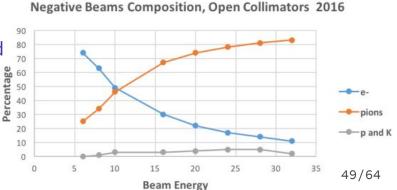
FTBF

Beam details:

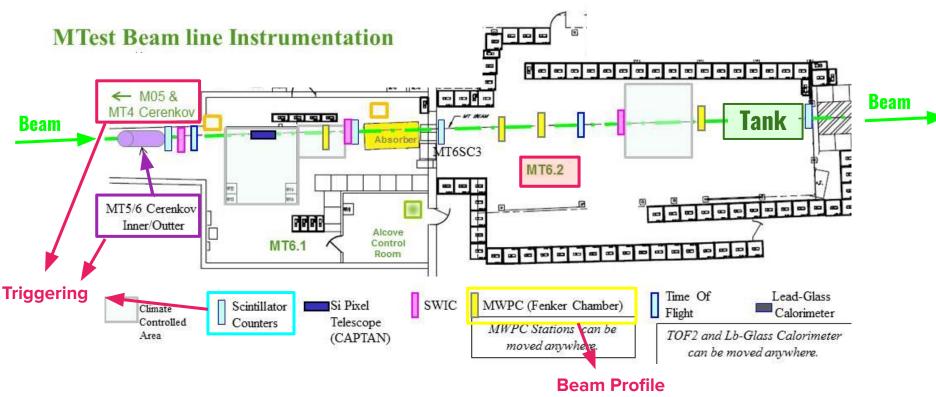
- → Using the secondary beamline at "MTest"
 - (120 GeV: Protons primary beam from Main Injector)
 - 8 60 GeV: Pions, (some protons possible)
 - **1 32 GeV**: Pions, electrons, kaons, or broadband muons
- → 4 second spill every 60 seconds
- → Tunable rate (100 Hz 100,000 Hz), beam available 24/7
- → At 4GeV pions make up ~30% of the beam
 - This fraction gets smaller as energy is decreased
- → Full details can be found:
 - http://ftbf.fnal.gov/beam-overview/





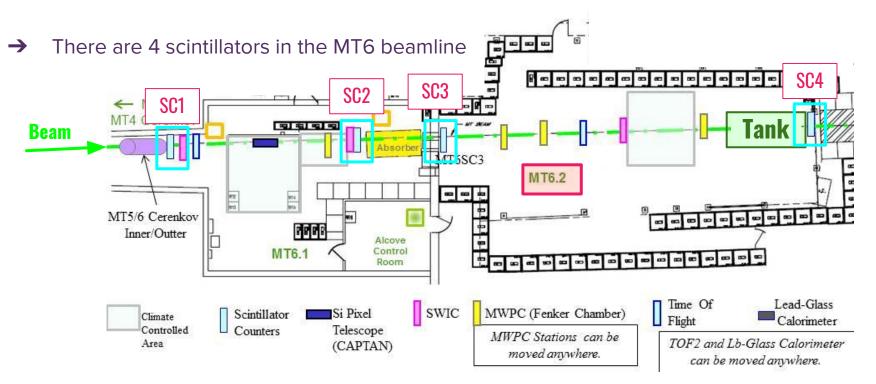


FTBF



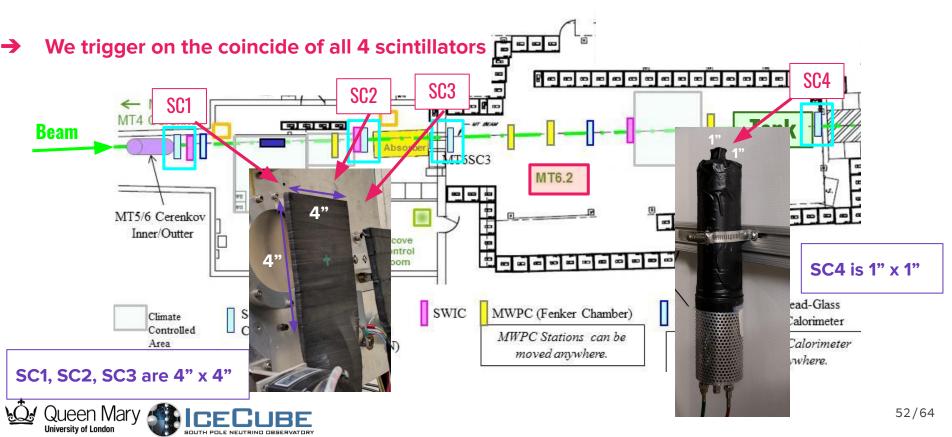


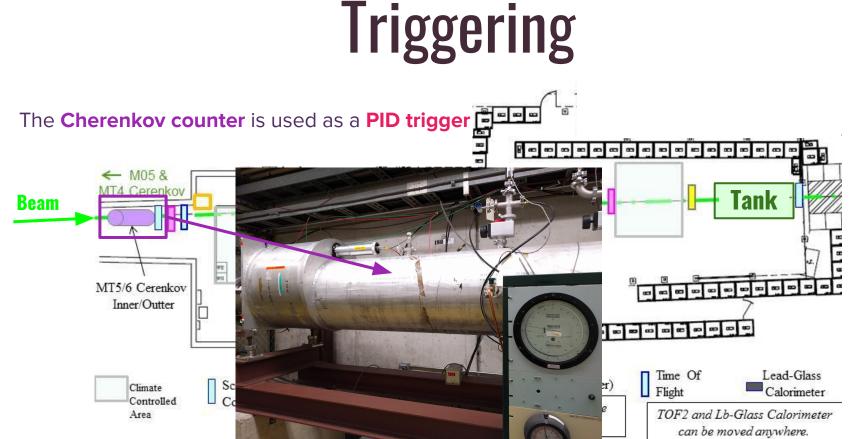
Triggering





Triggering

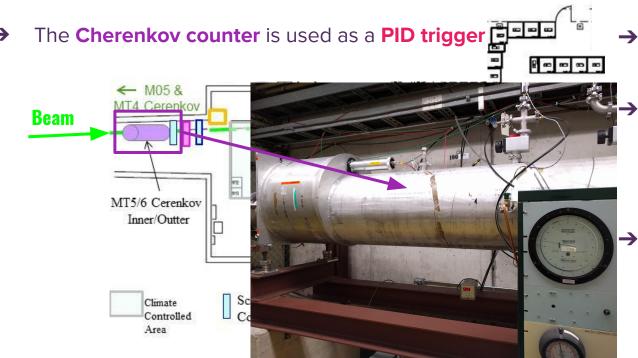






→

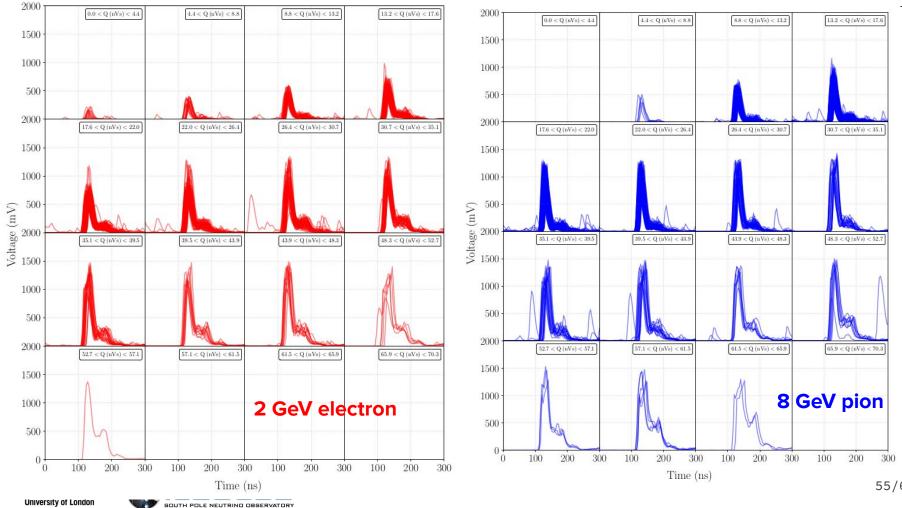
Triggering



- We trigger on the Cherenkov counter to select electrons
 - We can also do an anti-coincidence to trigger on everything but electrons i.e. MIPs
- We use this in (anti-) coincidence with the **4 scintillators** for the final trigger



S. Mandalia .0-17



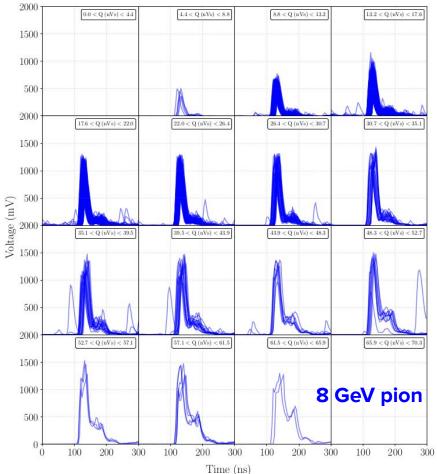
b

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S. Mandalia 10-17

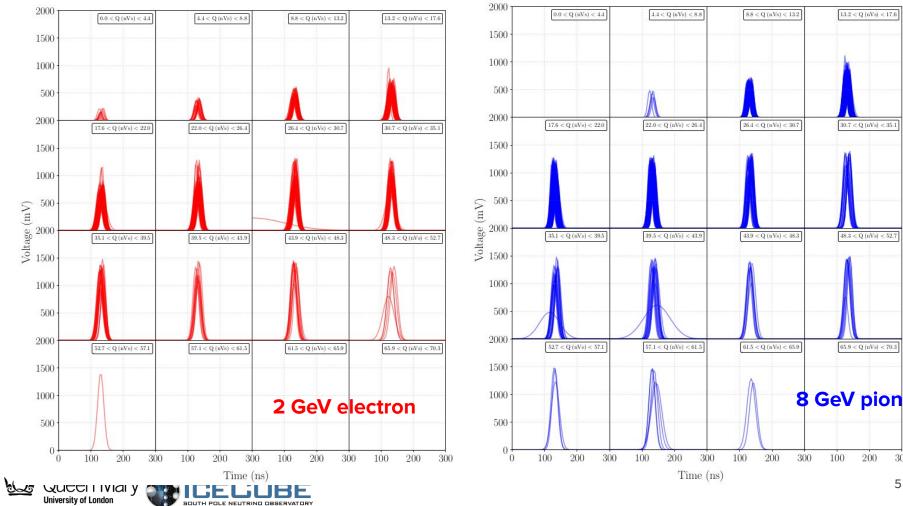
Data

- → Waveforms observe a primary and secondary peak
 - Secondary peak are due to late pulses which are caused by photoelectron backscattering off the dynodes
- → We isolate the primary pulse by fitting the primary peak with a Gaussian





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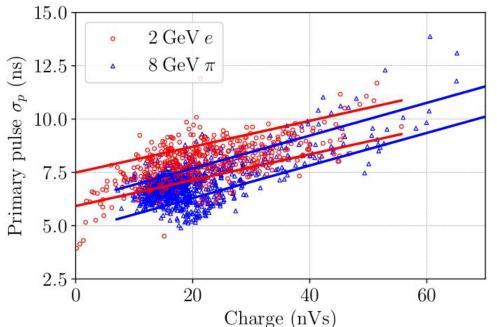
300

Results

→ Comparing waveforms of a similar charge

- Discriminator is simply the width of the primary pulse
- Bands show the one sigma containment region
- → At low charge, there is a possible discrimination that can be made between 2 GeV electron and 8 GeV pions
- → These techniques can be applied to future neutrino telescopes focusing on low energy physics, such as the IceCube-Upgrade









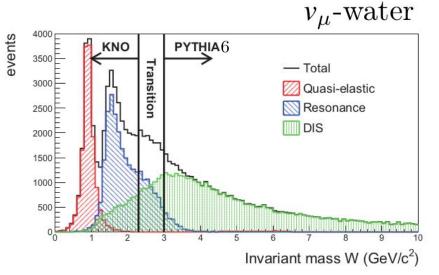
Motivation

Hadronization is important at high energies when DIS dominates

- → High-energy atmospheric neutrinos
 - IceCube (DeepCore) and

IceCube-Upgrade

- few GeV to ~ 100 GeV oscillation physics
- Final state particles from hadronization affects total deposited energy



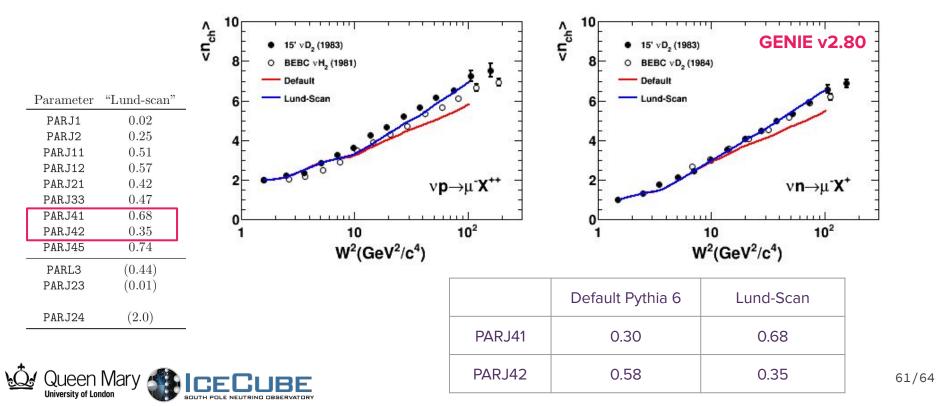
→ ~ TeV energy neutrinos

Hadronization is important if GENIE is extended to the few TeV region (ongoing work)



Pythia 6 Tunings

HERMES (low energy) parameter sets have better agreements with the charged hadron multiplicity data



Pythia 8

Pythia 6 is outdated and no longer supported!

- → Upgrade the GENIE to use Pythia 8
 - Directly affects the Hadronization and Decay routines
 - Pythia 8 can be directly interfaced with in C++ so there is no need for ROOT wrapper classes (e.g. TPythia6, TMCParticle)
- → Incorporate different physics aspects of Pythia 8
 - Physics which were absent in Pythia 6 e.g. beam remnant structure
 - Tuning of Pythia 8 parameters to bubble chamber data

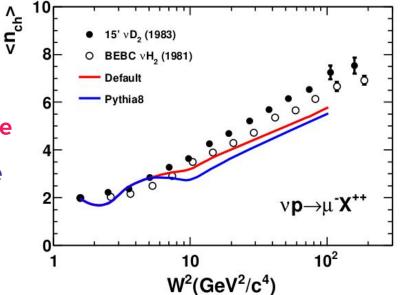


[1] Katori, Mandalia, J.Phys. G42 (2015) no.11, 115004

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Direct Translation to Pythia 8

- Showing averaged charged multiplicity
 - Comparing "Default" Pythia 6 hadronization model with Pythia8
 - Compared also to bubble chamber data
- → Will be implemented in next GENIE release
- Both Pythia 6 and Pythia 8 underestimate bubble chamber data
 - Pythia is tuned to high energy LHC experiments
 - We have the technology needed to tune the the hadronization model [1]



Neutrino interferometry for high-precision tests of Lorentz symmetry with IceCube

The IceCube Collaboration*

IceCube, Nature.Phy.14,961(2018)

T. Katori¹⁴⁵*, S. Mandalia¹⁴⁵*, J. Kiryluk⁴⁶, M. Lesiak-Bzdak⁴⁶, H. Niederhausen⁴⁶, Y. Xu⁴⁶, G. Kohnen⁴⁷,

PYTHIA hadronization process tuning in the GENIE neutrino interaction generator

J. Phys. G42, 115004 (2015)

Teppei Katori and Shivesh Mandalia

Probe Of Sterile Neutrinos Using Astrophysical Neutrino Flavor arXiv:1909.05341

Carlos A. Argüelles,^{1,*} Kareem Farrag,^{2,3,†} Teppei Katori,^{4,‡} Rishabh Khandelwal,^{5,§} Shivesh Mandalia,^{2,¶} and Jordi Salvado^{6,**}

Computational Techniques for the Analysis of Small Signals in High-Statistics Neutrino Oscillation Experiments

K. B. M. Mahn^w, S. Mancina^{ah}, S. Mandalia^{af}, S. Marka^{aq}, Z. Marka^{aq}, R. Maruyama^{ap},

irXiv:1803.05390

Thanks for listening!

Test of Lorentz Violation with Astrophysical Neutrino Flavor in arXiv:1906.09240 IceCube

Teppei Katori¹, Carlos A. Argüelles², Kareem Farrag^{1,3}, and Shivesh Mandalia¹

First look at the PYTHIA8 hadronization program for neutrino interaction generators JPS Conf. Proc. 12, 010033 (2016)

Teppei KATORI, Pierre LASORAK, Shivesh MANDALIA, and Ryan TERRI

Quest for new physics using astrophysical neutrino

flavor in IceCube

arXiv:1908.07602

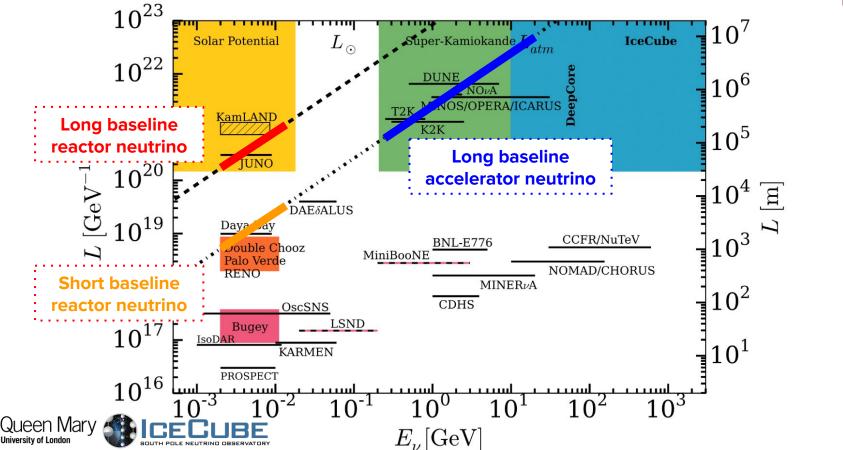
Corresponding authors: Carlos A. Argüelles¹, Kareem Farrag^{†2,3}, Teppei Katori², Shivesh Mandalia²



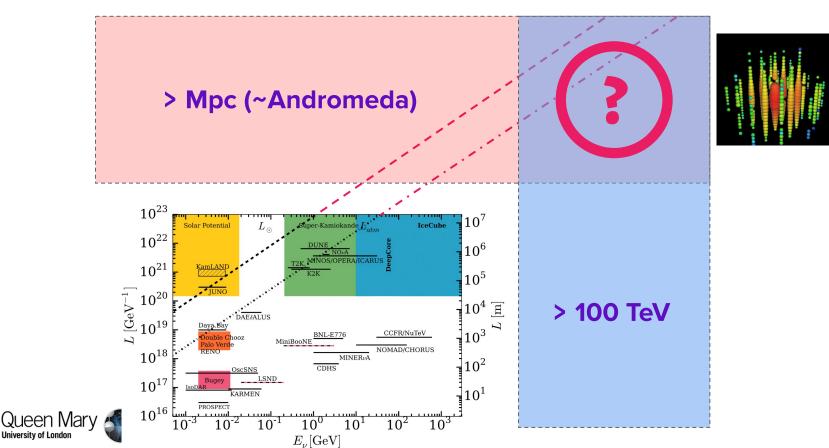
Terrestrial Oscillation Measurement Landscape

S. Mandalia

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Astrophysical Frontier?

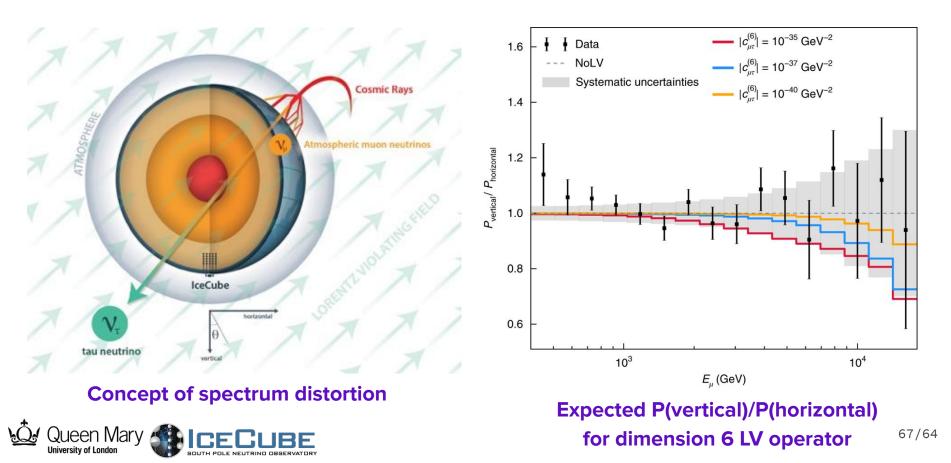


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Neutrino interferometry for high-precision tests of Lorentz symmetry

S. Mandalia



Neutrino interferometry for high-precision tests of Lorentz symmetry

400 GeV < E < 18 TeV ("conventional") -1 < cos(zenith) < 0 ("up-going")

very similar to 2016 sterile neutrino analysis sample
https://icecube.wisc.edu/science/data/HE_NuMu_diffuse

- → Atmospheric neutrinos from MCEq
- → Simple power law for astrophysical neutrinos
- → DIS cross section from **CSS** paper
- → Analytical oscillation formula

Normalization of flux:

Binning

- 1. conventional (40%)
- 2. prompt (free)

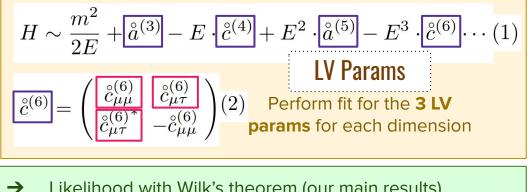
Flux components:

- 3. astrophysical (free)
- 6 nuisance

Sytematics

params

- 1. primary cosmic ray index (2%)
- 2. astrophysical neutrinos index (25%)
- 3. p/K ratio for conventional flux (10%)



- Likelihood with Wilk's theorem (our main results)
 Bayesian Markov Chain Monte Carlo
 Eit M
 - Bayesian Markov Chain Monte Carlo http://dan.iel.fm/emcee/current



S. Mandalia

Simulation

Fit Methods

2019-10-17

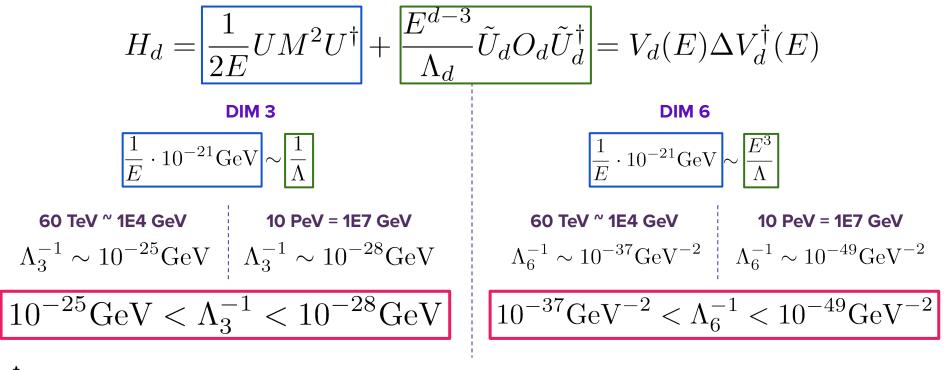
dim.	method	type	sector	limits	ref.	dalia
3	CMB polarization	astrophysical	photon	$\sim 10^{-43} { m GeV}$	[6]	10-17
	He-Xe comagnetometer	tabletop	neutron	$\sim 10^{-34} { m GeV}$	[10]	
	torsion pendulum	tabletop	electron	$\sim 10^{-31}~{ m GeV}$	[12]	
	muon g-2	accelerator	muon	$\sim 10^{-24}~{ m GeV}$	[13]	
	neutrino oscillation	atmospheric	neutrino	$ \operatorname{Re}(\overset{\circ}{a}{}^{(3)}_{\mu\tau}) , \operatorname{Im}(\overset{\circ}{a}{}^{(3)}_{\mu\tau}) < 2.9 \times 10^{-24} \text{ GeV } (99\% \text{ C.L.}) < 2.0 \times 10^{-24} \text{ GeV } (90\% \text{ C.L.})$	this work	
4	GRB vacuum birefringence	astrophysical	photon	$\sim 10^{-38}$	[7]	-
	Laser interferometer	LIGO	photon	$\sim 10^{-22}$	[8]	
	Sapphire cavity oscillator	tabletop	photon	$\sim 10^{-18}$	[5]	
	Ne-Rb-K comagnetometer	tabletop	neutron	$\sim 10^{-29}$	[11]	
	trapped Ca^+ ion	tabletop	electron	$\sim 10^{-19}$	[14]	_
	neutrino oscillation	atmospheric	neutrino	$\frac{ \operatorname{Re}(\mathring{c}_{\mu\tau}^{(4)}) , \operatorname{Im}(\mathring{c}_{\mu\tau}^{(4)}) }{< 2.7 \times 10^{-28} (90\% \text{ C.L.})} $	this work	
5	GRB vacuum birefringence	astrophysical	photon	$\sim 10^{-34} \text{ GeV}^{-1}$	[7]	
	ultra-high-energy cosmic ray	astrophysical	proton	$\sim 10^{-22}$ to $10^{-18} { m GeV^{-1}}$	[9]	_
	neutrino oscillation	atmospheric	neutrino	$\begin{aligned} \text{Re}\left(\overset{\circ}{a}{}^{(5)}_{\mu\tau}\right) , \text{Im}\left(\overset{\circ}{a}{}^{(5)}_{\mu\tau}\right) &< 2.3 \times 10^{-32} \text{ GeV}^{-1} \text{ (99\% C.L.)} \\ &< 1.5 \times 10^{-32} \text{ GeV}^{-1} \text{ (90\% C.L.)} \end{aligned}$	this work	
6	GRB vacuum birefringene	astrophysical	photon	$\sim 10^{-31}~{ m GeV^{-2}}$	[7]	
	ultra-high-energy cosmic ray	astrophysical	proton	$\sim 10^{-42}$ to 10^{-35} GeV ⁻²	[9]	
	gravitational Cherenkov radiation	astrophysical	gravity	$\sim 10^{-31}~{ m GeV^{-2}}$	[15]	_
	neutrino oscillation	atmospheric	neutrino	$ \operatorname{Re}(\mathring{c}_{\mu\tau}^{(6)}) , \operatorname{Im}(\mathring{c}_{\mu\tau}^{(6)}) < 1.5 \times 10^{-36} \text{ GeV}^{-2} (99\% \text{ C.L.}) < 9.1 \times 10^{-37} \text{ GeV}^{-2} (90\% \text{ C.L.})$	this work	
7	GRB vacuum birefringence	astrophysical	photon	$\sim 10^{-28} { m ~GeV^{-3}}$	[7]	
2	neutrino oscillation	atmospheric	neutrino	$\begin{aligned} \text{Re}\left(\mathring{a}_{\mu\tau}^{(7)}\right) , \text{Im}\left(\mathring{a}_{\mu\tau}^{(7)}\right) &< 8.3 \times 10^{-41} \text{ GeV}^{-3} \text{ (99\% C.L.)} \\ &< 3.6 \times 10^{-41} \text{ GeV}^{-3} \text{ (90\% C.L.)} \end{aligned}$	this work	
8	gravitational Cherenkov radiation	astrophysical	gravity	$\sim 10^{-46} { m GeV^{-4}}$	15	_
	neutrino oscillation	atmospheric	neutrino	$ \operatorname{Re}(\mathring{c}_{\mu\tau}^{(8)}) , \operatorname{Im}(\mathring{c}_{\mu\tau}^{(8)}) < 5.2 \times 10^{-45} \text{ GeV}^{-4} (99\% \text{ C.L.}) < 1.4 \times 10^{-45} \text{ GeV}^{-4} (90\% \text{ C.L.})$	this work	
8						101

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TABLE I: Comparison of attainable best limits of SME coefficients in various fields. IceCube, Nature.Phy.14,961(2018)

Back of the Envelope Sensitivity

Binning = [60 TeV - 10 PeV] 20 bins



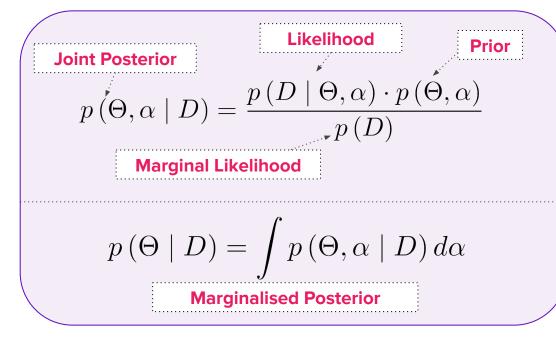
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2019-10-17

Markov Chain Monte Carlo

 $p(D) = \int p(D \mid \Theta, \alpha) \cdot p(\Theta, \alpha) \ d\Theta \ d\alpha$

The goal is to obtain the posterior distribution marginalised over your nuisańce parameters



Problem: In a multidimensional space, the integral over the nuisance parameters is difficult to perform directly

Instead: Use an MCMC to sample over the joint posterior, after which the sampled points can be integrated over in a more straightforward way to obtain the marginalised posterior

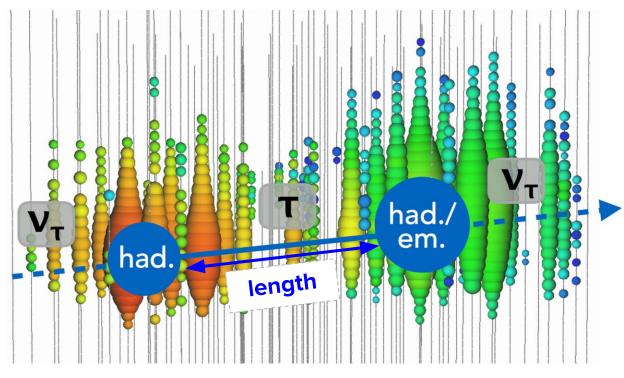


Flavour Identification

Updated search for astrophysical tau neutrinos with starting events with double cascades!

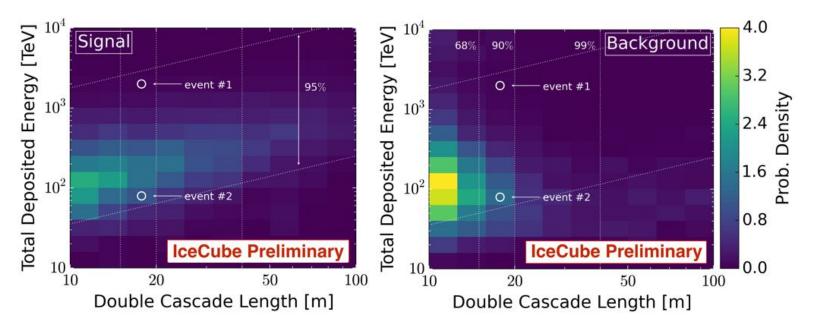
Require:

- → Both cascades with E > 1 TeV
- → Separation distance > 10 m





Two Candidate Events



- Two double cascade events have been identified!
- -> Could have arisen from v_{τ} or background (astro. or atm.)
- Dedicated studies are on-going

Jueen Marv

University of London

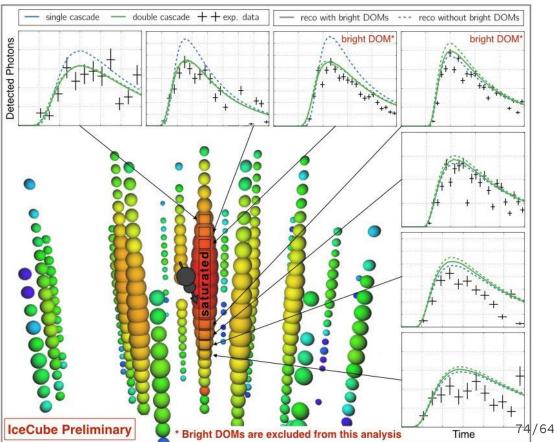
Two Candidate Events

Candidate Event 1

Some preference for double cascade over single cascade based on DOM-to-DOM charge distributions

Dedicated studies ongoing





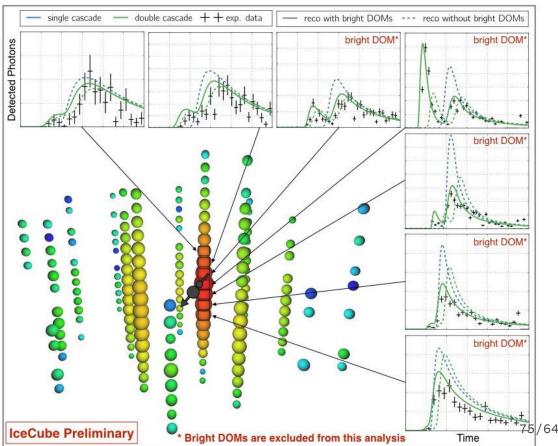
Two Candidate Events

Candidate Event 2

Strong preference for double cascade over single cascade based on DOM-to-DOM charge distributions

Dedicated studies ongoing





Setup

DOM

- → DOM must be made water-tight
- → Penetrator was closed up by using grey RTV glue





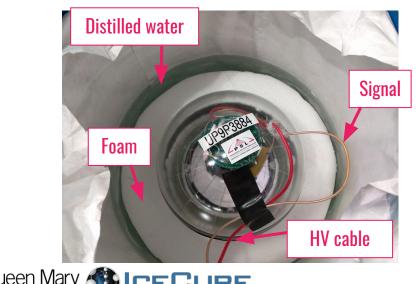
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Setup

DOM

University of London

- → DOM must be made water-tight
- → Penetrator was closed up by using grey RTV glue
- → Used a 55-gallon drum filled with distilled water as test



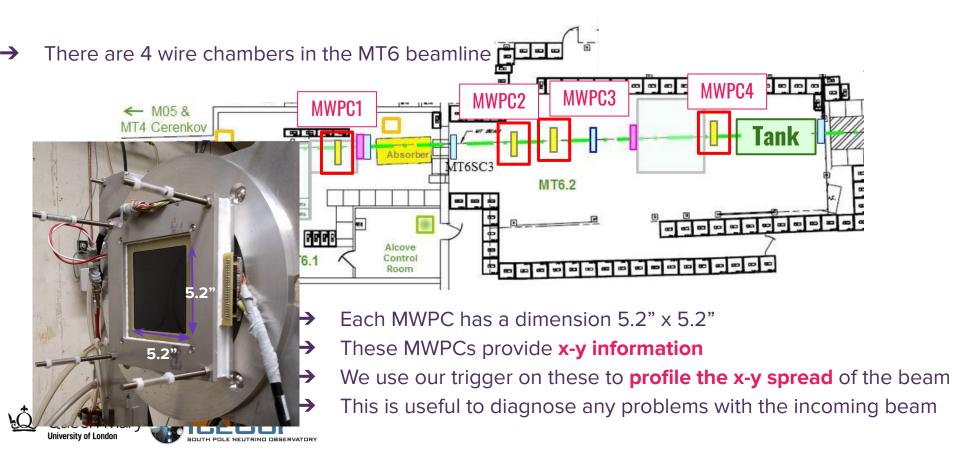


- → DOM is very buoyant!
- → For stability we made a foam

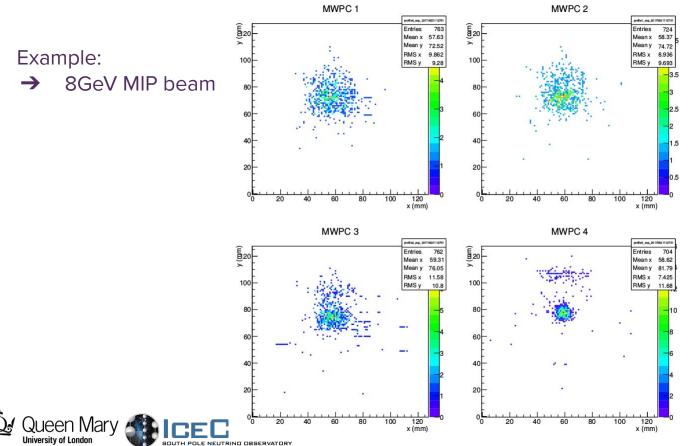
structure which houses the DOM -

"floating island"

Beam Profile



Beam Profile



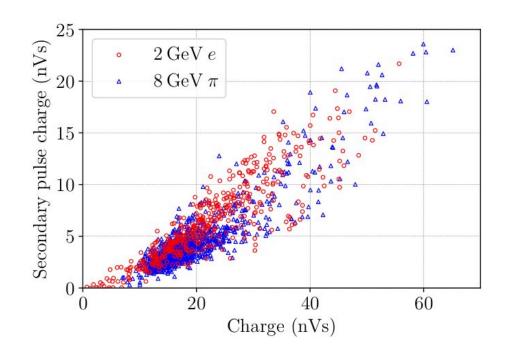


Figure 5.15: Charge of the secondary pulse plotted as a function of the charge of the waveform. In red circles show data points for a 2 GeV electron beam and in blue triangles show data points for an 8 GeV pion beam. For both configurations, it can be seen that the second pulse grows linearly with the total waveform.



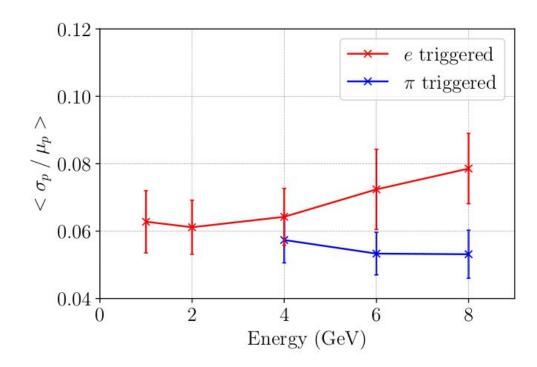


Figure 5.16: Average of the mean normalised primary peak sigma as a function of the beam energy for an electron trigger in red and a pion trigger in blue.



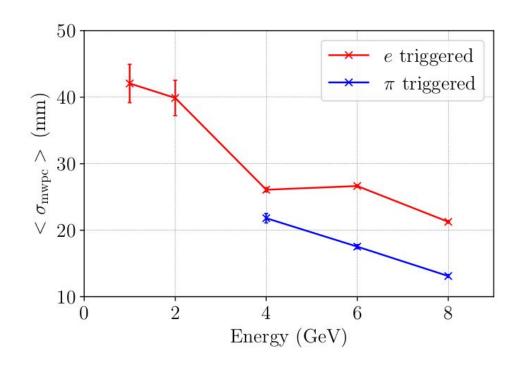


Figure 5.17: Average spread of the beam as a function of the beam energy for an electron trigger in red and a pion trigger in blue. Spread was computed from beam profile data which was taken using 3 multi-wire proportional chambers.



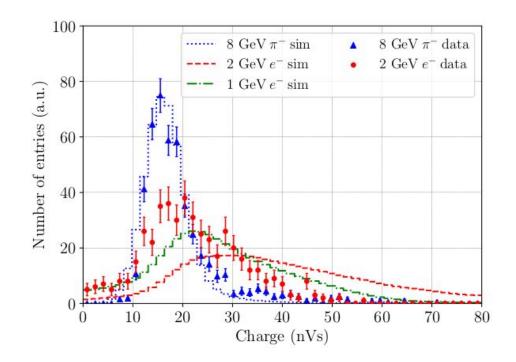


Figure 5.20: Hit distribution of the number of photons hitting the PMT for both electron and pion events. Data points taken from this beam test are compared to Geant4 simulations. The presence of saturation can be seen to have a large affect for very high charge data, such as for the 2 GeV data with electron triggering.



Driving PYTHIA6 from GENIE

Some amount of monkey business in making **quark** + **diquark** assignments most certainly due to our own unfamiliarity with PYTHIA. Luckily, overall generation outcomes not sensitive to choices made.

Init state	Hit	Leading	Remnant	PYTHIA6 assignment		Weirdness
	quark	quark	system			level
$\nu + p CC$	d valence	$(d \rightarrow) u$	uu	u	uu	
$\nu + p CC$	d sea	$(d \rightarrow) u$	\bar{d} + uud	u	uu	*
$\nu + p CC$	s sea	$(s \rightarrow) u$	s + uud	u	uu	**
$\nu + p CC$	ū sea	$(\bar{u} \rightarrow) \bar{d}$	u + uud	u	uu	***
$\nu + n CC$	d valence	$(d \rightarrow) u$	ud	u	ud	
$\nu + n CC$	d sea	$(d \rightarrow) u$	\bar{d} + udd	u	ud	*
$\nu + n CC$	s sea	$(s \rightarrow) u$	s + udd	u	ud	**
ν + n CC	ū sea	$(\bar{u} \rightarrow) \bar{d}$	<mark>u + uud</mark>	u	ud	***



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Improvements

